

# 6d SCFTs compactified on a circle/torus

Kazuya Yonekura, IAS

Mainly based on

- **To appear** and arXiv:1503.06217  
with K. Ohmori, H. Shimizu, Y. Tachikawa,

also

- arXiv:1410.6868 with H. Hayashi Y. Tachikawa,
- arXiv:1505.04743

# Introduction

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In six dimensions, we have many superconformal field theories!

- 6d  $N=(2,0)$  SCFTs: classified by ADE Lie algebra
- 6d  $N=(1,0)$  SCFTs: so many of them, classified recently in F-theory

[Heckman, Morrison, Vafa, 2013]

[Del Zotto, Heckman, Tomasiello, Vafa, 2014]

[Heckman, Morrison, Rudelius, Vafa, 2015]

# Introduction

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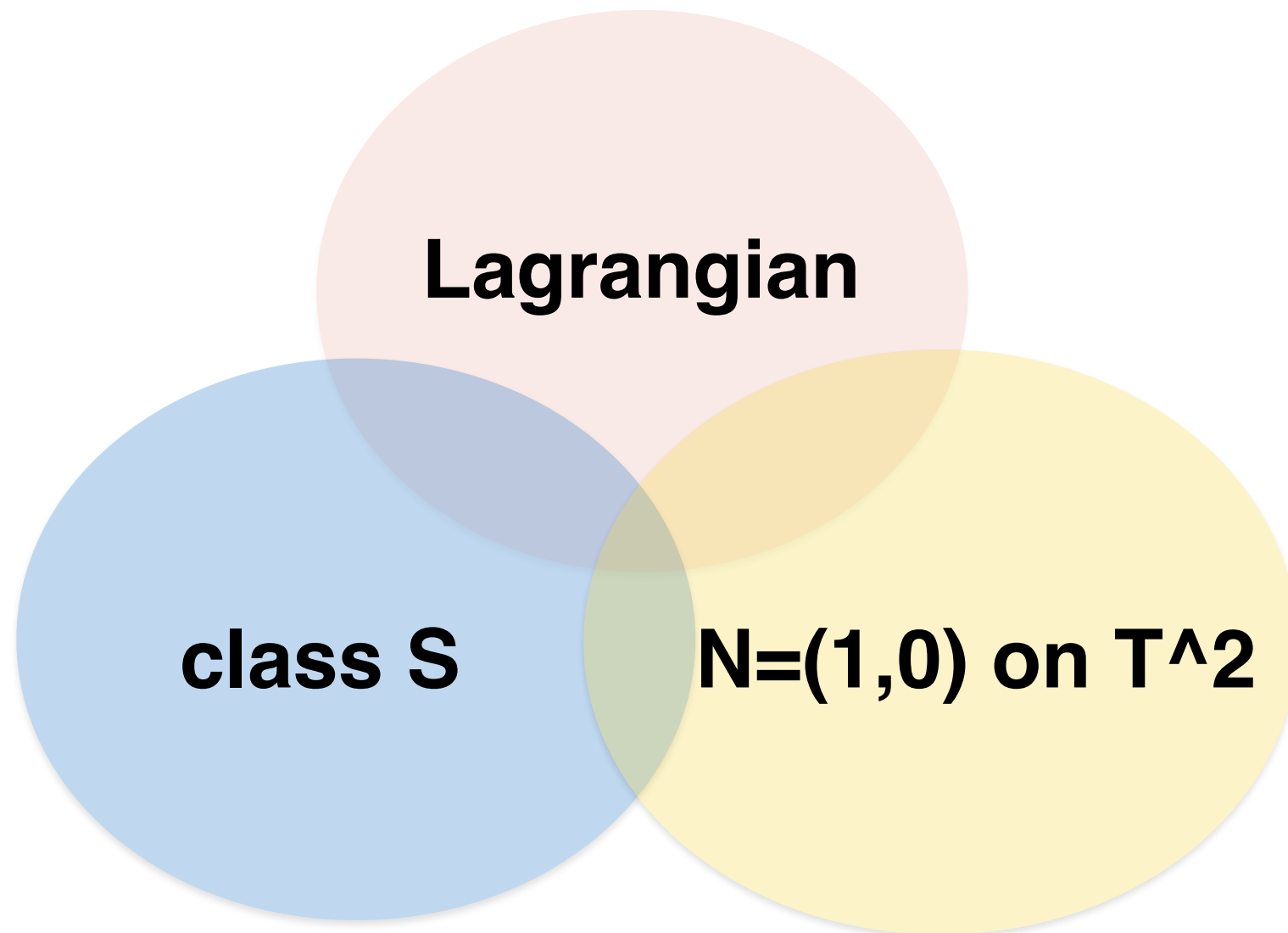
There are at least three ways to get 4d  $N=2$  theories.

- Lagrangian field theories.
- 6d  $N=(2,0)$  theories compactified on Riemann surfaces.
- 6d  $N=(1,0)$  theories compactified on a torus.

Compactification of  $N=(2,0)$  theories, called **class S**, have given us many new theories which might not have Lagrangian description.

# Introduction

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- What theories do we get from **6d N=(1,0)** SCFTs on  $T^2$ ?
- How much overlap do they have?

# Contents

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1. Introduction

2. A class of  $N=(1,0)$  generalizing  $N=(2,0)$  theories

3. Compactification on a circle

4. Compactification a torus

5. Further generalization

6. Summary

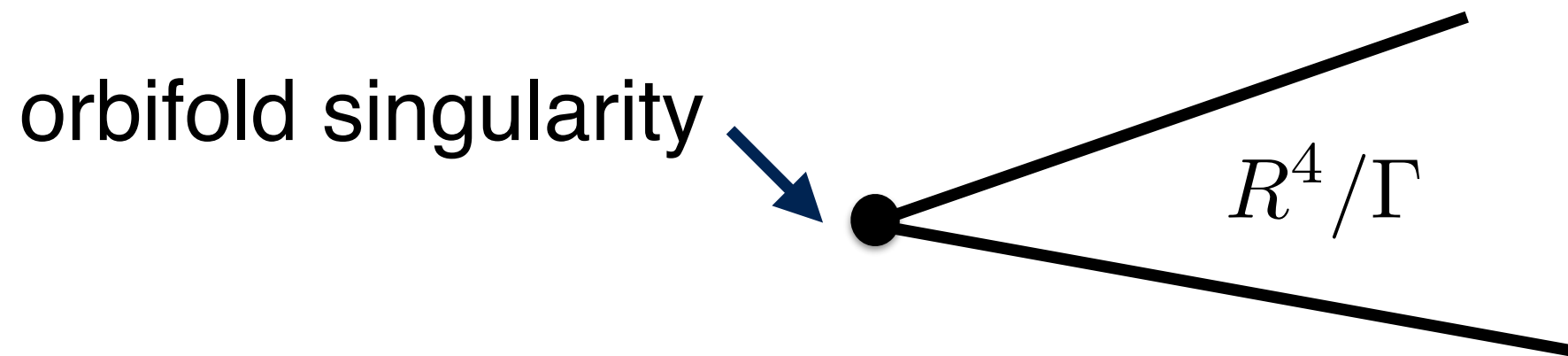
# N=(2,0) in type IIB string

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N=(2,0) theories can be constructed by type IIB string on

$$R^{1,5} \times (R^4/\Gamma)$$

$\Gamma$  is an orbifold group classified by ADE Dynkin diagram.  
(E.g., for A-type,  $\Gamma = Z_N$  )



The N=(2,0) theory is realized on the singularity.

[Witten, 1995]

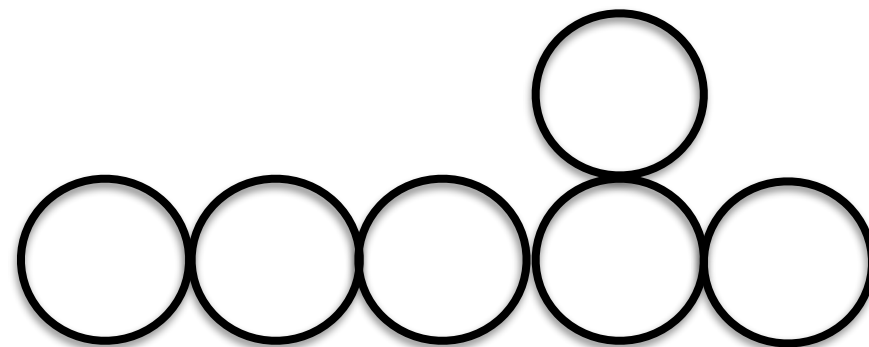
# $N=(2,0)$ in type IIB string

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singular point ●



resolution of singularity  
(blowup)



(2-cycles)

By the resolution of the singularity,  
the singular point becomes several intersecting 2-cycles.

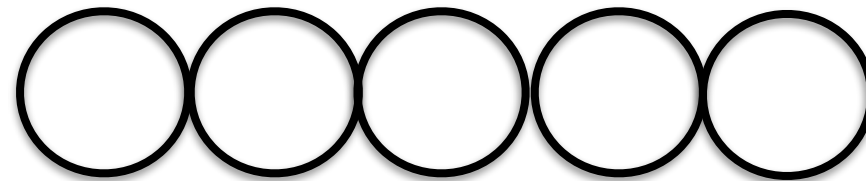
Field theoretically, this corresponds to giving vevs to some  
scalar fields of the  $N=(2,0)$  theory.

# N=(2,0) in type IIB string

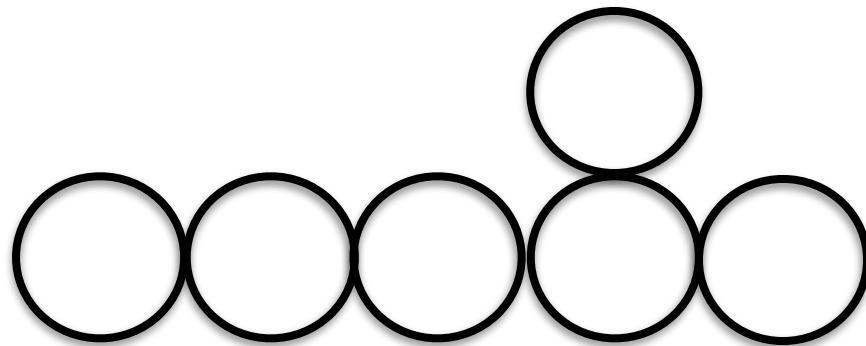
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The intersection of the 2-cycles is determined by the Dynkin diagram corresponding to the orbifold group  $\Gamma$ .

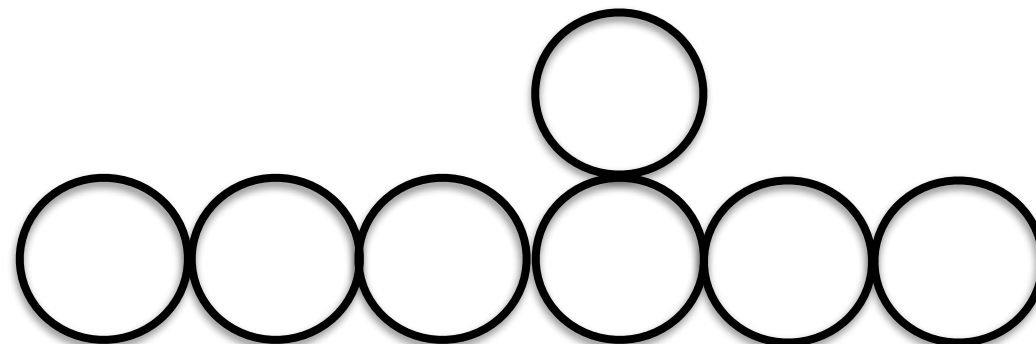
$\Gamma : A_N$  type



$\Gamma : D_N$  type



$\Gamma : E_N$  type

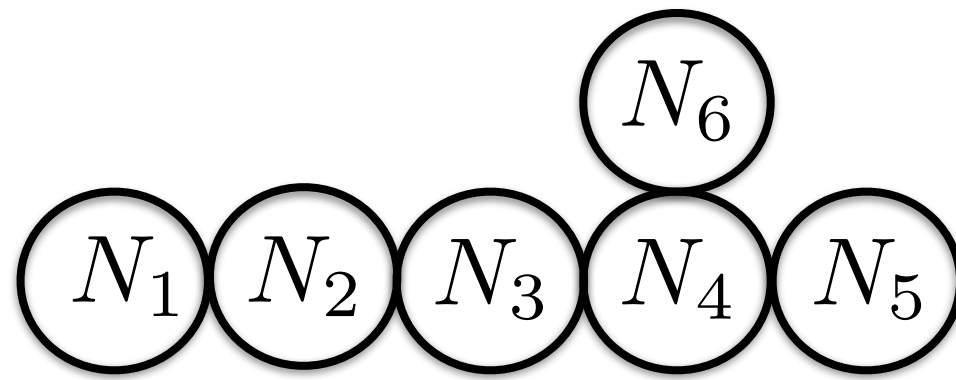




# Generalization to $N=(1,0)$

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We can **wrap D7-branes** on each of the 2-cycles.



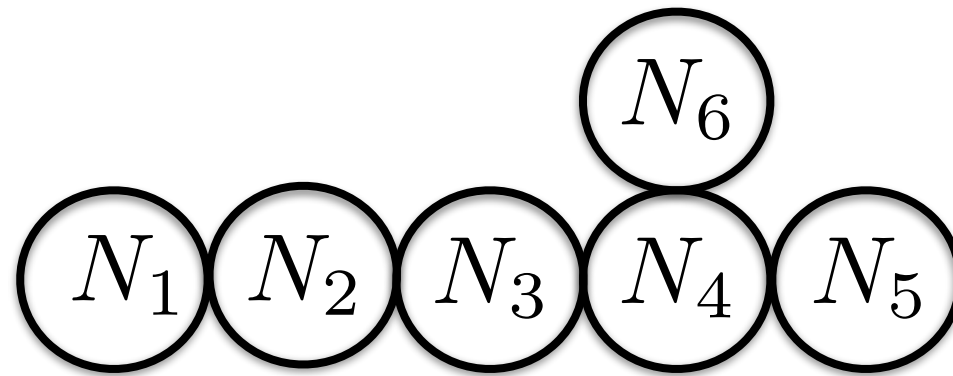
Each 2-cycle now supports  $SU(N_i)$  gauge group.

By shrinking all the 2-cycles with D7-branes,  
we get  $N=(1,0)$  SCFT at the singularity.

$$\mathcal{L} \sim \phi \operatorname{tr}(F_{\mu\nu} F^{\mu\nu}) \quad (\phi : \text{size of the 2-cycles.})$$

# Generalization to $N=(1,0)$

---



- Each 2-cycle has  $SU(N_i)$  gauge group.
- There are bifundamental between adjacent 2-cycles
- Anomaly cancellation requires additional fundamentals such that each gauge group has

$$[\text{Total flavor number}] = 2 [\text{Color number}]$$

We get a quiver gauge theory with the above properties.

# Contents

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1. Introduction

2. A class of  $N=(1,0)$  generalizing  $N=(2,0)$  theories

3. Compactification on a circle

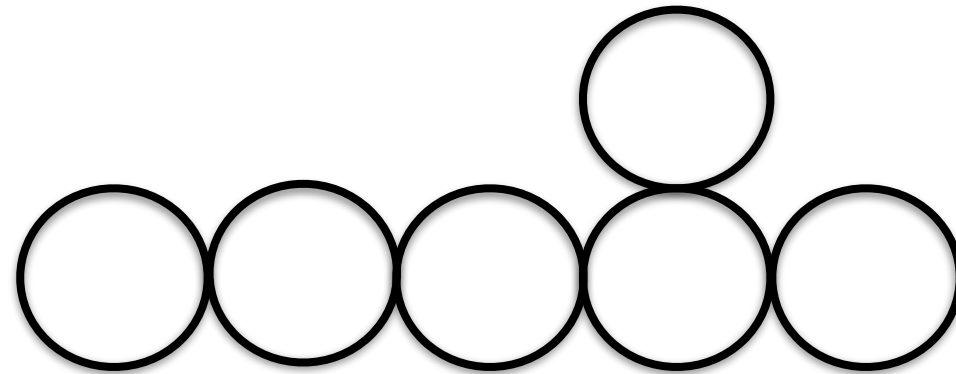
4. Compactification on a torus

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6. Summary

# $N=(2,0)$ on a circle

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The size of each 2-cycle is given by a vev of a scalar field in the so-called tensor multiplet in 6d.

6d tensor multiplets



on a circle

5d vector multiplets



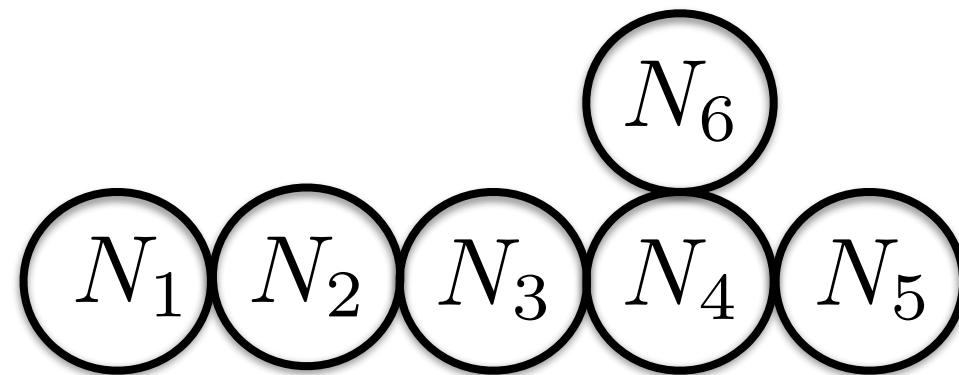
vevs  $\rightarrow 0$

5d  $N=2$  Super-Yang-Mills of type  $G=ADE$

# $N=(1,0)$ on a circle

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Now we consider  $N=(1,0)$  theories on a circle.



- This gives a 5d  $SU(N)$  quiver after compactification.
- The gauge couplings are determined by the size of the 2-cycles, which are vevs of free vector multiplets.

$$\mathcal{L} \sim \phi \operatorname{tr}(F_{\mu\nu} F^{\mu\nu}) \quad (\phi : \text{scalar in vector mult.})$$

What happens when vevs  $\rightarrow 0$  ?

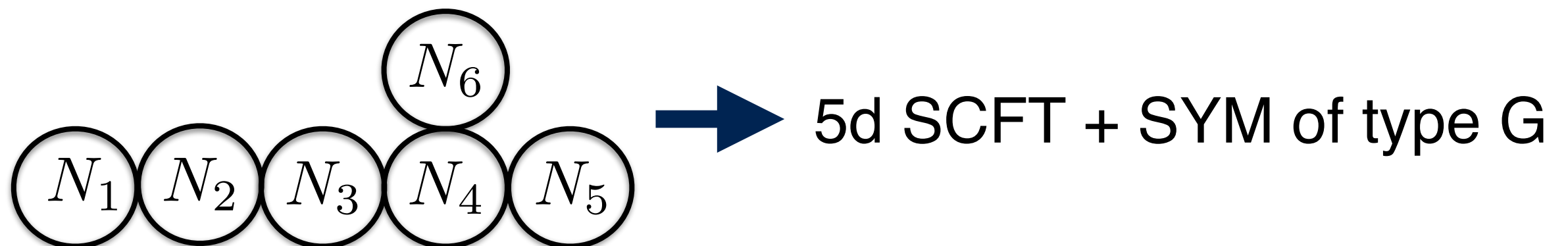
# $N=(1,0)$ on a circle

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## The two effects of $\text{vevs} \rightarrow 0$ ;

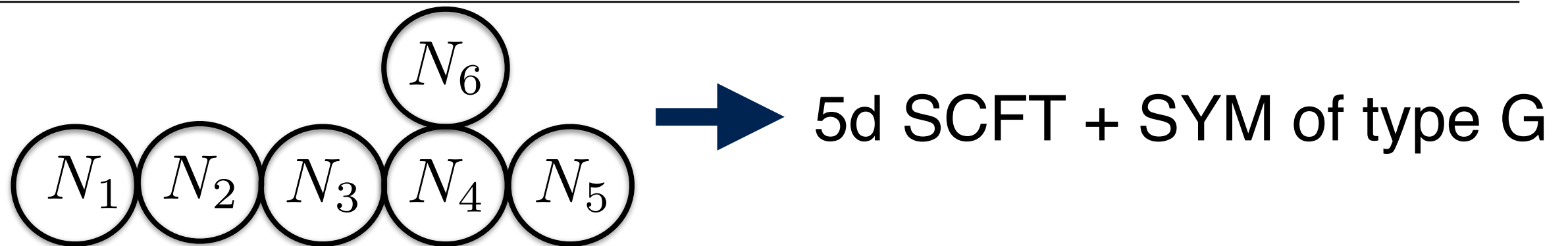
- 5d  $SU(N)$  gauge couplings become infinitely strong.  
→ 5d  $SU(N)$  quiver becomes a **5d SCFT**
- The gauge symmetry of the underlying  $N=(2,0)$  restored.  
→ we get **5d SYM of type G**

(G is determined by the ADE Dynkin diagram.)



# $N=(1,0)$ on a circle

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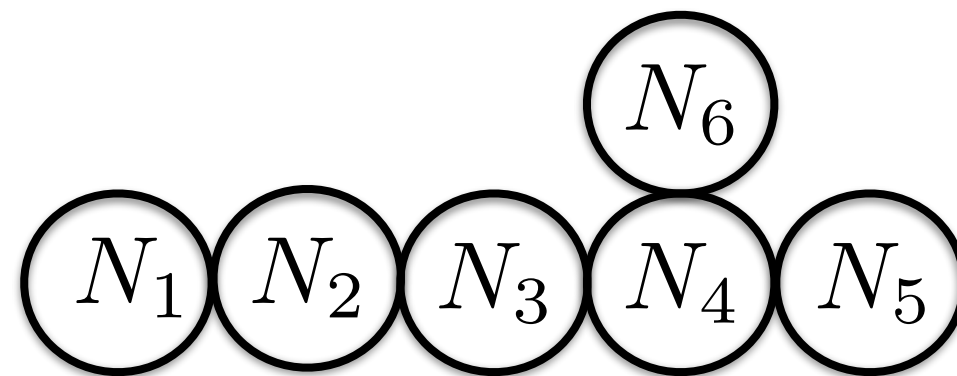


The 5d SCFT and SYM of type G must be coupled, so the 5d SCFT must have flavor G symmetry which is gauged by SYM of type G.

This is indeed the case by **5d symmetry enhancement!**

# N=(1,0) on a circle

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In 5d, each gauge group has instanton symmetry and baryon symmetry  $j_I \sim *(\text{tr} F \wedge F)$   $j_B$

They give  $U(1)^{2\text{rank}G}$  flavor symmetries.

If each  $SU(N)$  gauge group has  $[\text{flavor}]=2[\text{color}]$  and zero Chern-Simons level,  $G \times G$  flavor symmetry is enhanced.  
[Tachikawa,2015], [KY,2015]

Diagonal subgroup is gauged by SYM.



# $N=(1,0)$ on a circle

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## Summary of circle compactification

6d SCFT of the class we are considering



[5d SCFT with flavor  $G$ -symm.] + [5d SYM of type  $G$ ]

The simplest example: in the case of  $N=(2,0)$ ,

[5d SCFT with  $G$ -symm.] = free adjoint hyper of  $G$ .

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---

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# N=(1,0) on a torus

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$$T^2 \sim S^1 \times S^1$$

6d SCFT



[5d SCFT with flavor G-symm.] + [5d SYM of type G]



???



[4d SYM of type G]

coupling  $\tau = \frac{4\pi i}{g^2} + \frac{\theta}{2\pi}$   
=complex structure of  $T^2$


$SL(2, Z)$  S-duality

# $N=(1,0)$ on a torus

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$T^2$  compactification of 6d SCFT is now reduced to  $S^1$  compactification of 5d SCFT.

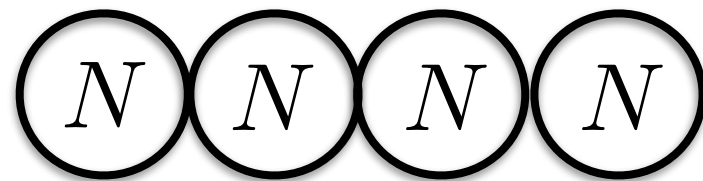
[5d SCFT with flavor G-symm.]

 on  $S^1$   
**???**

# Example: Linear $SU(N)$ quiver

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As interesting example, consider the theory



$$G = A_{M-1}$$

- All the gauge groups have the same color  $N$ .
- The underlying  $N=(2,0)$  theory is of type  $G = A_{M-1}$
- The corresponding 5d quiver is

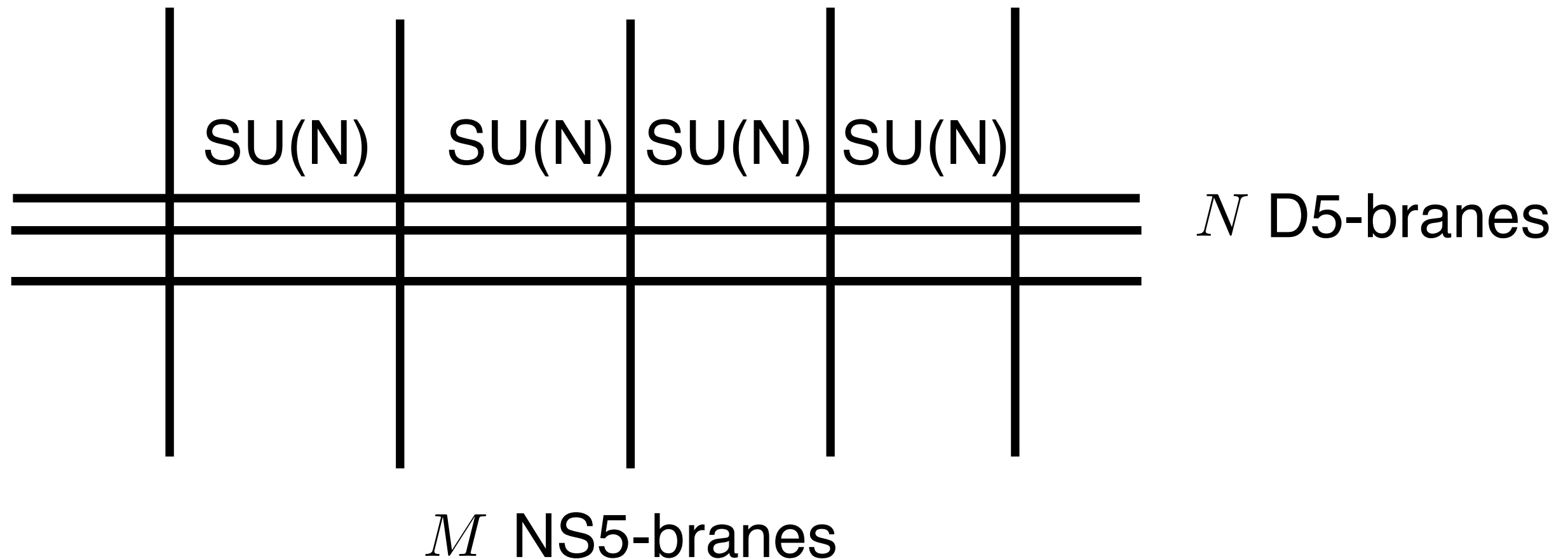
$$[N] - SU(N) - SU(N) - SU(N) - SU(N) - [N]$$

This is  $SU(N)^{M-1}$  quiver.

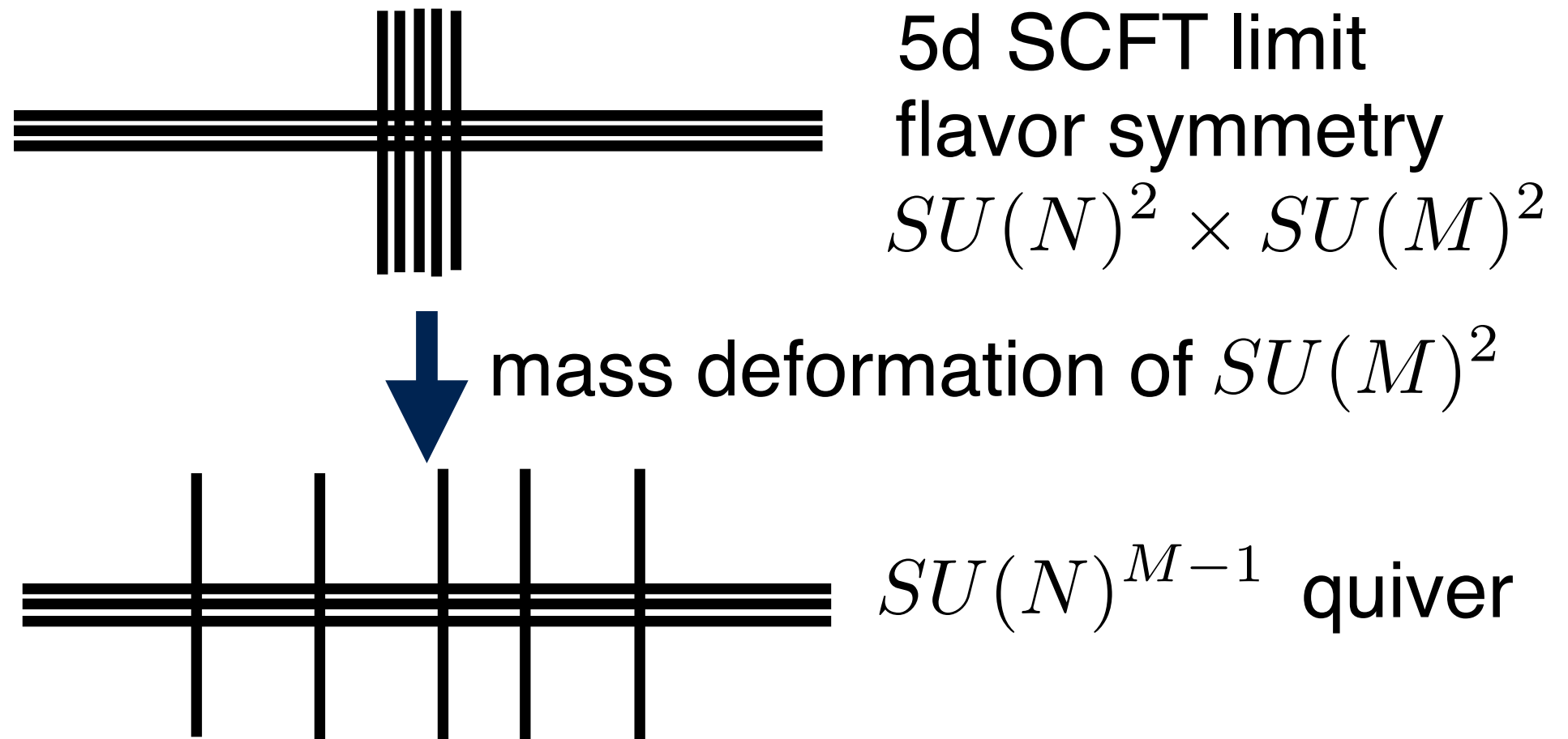
# Brane web realization

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The 5d theory can be realized by a NS5-D5 system



# Brane web realization



# 4d situation

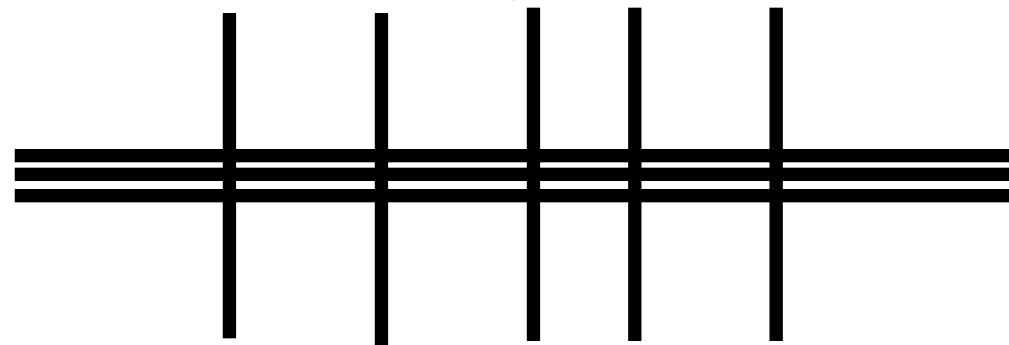
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???

(obtained from the 5d SCFT  
on circle compactification)



mass deformation of  $SU(M)^2$



$SU(N)^{M-1}$  quiver in 4d



# 4d situation

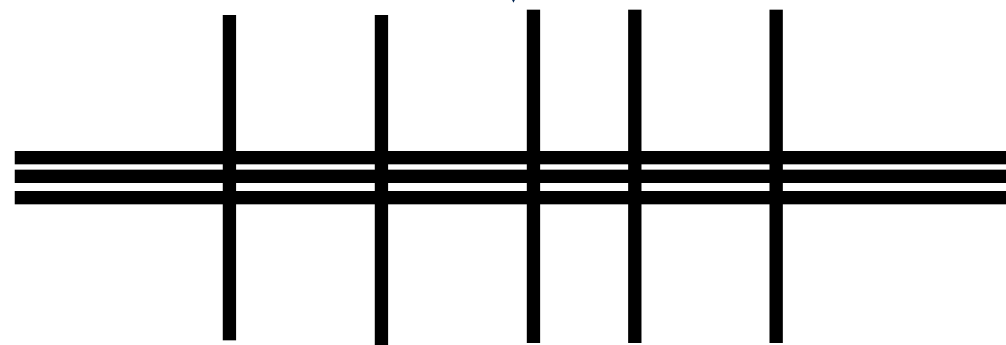
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???

(obtained from the 5d SCFT  
on circle compactification)

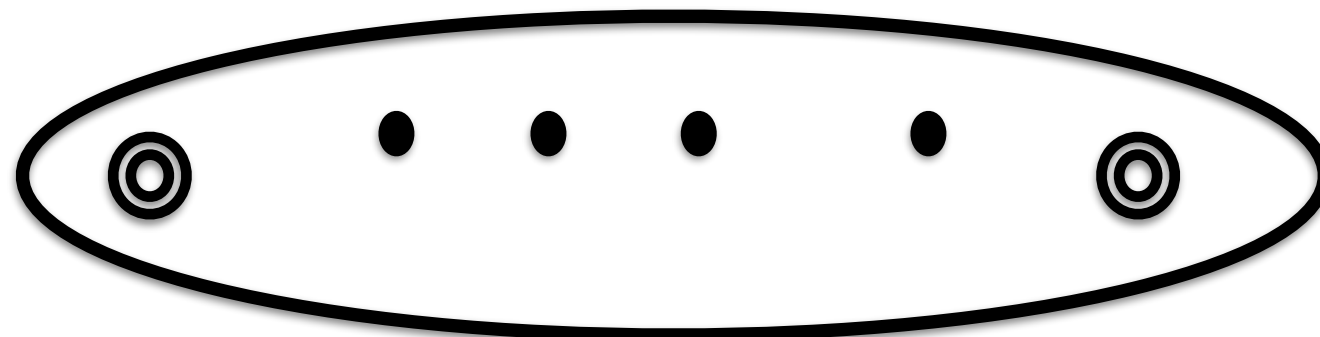


mass deformation of  $SU(M)^2$



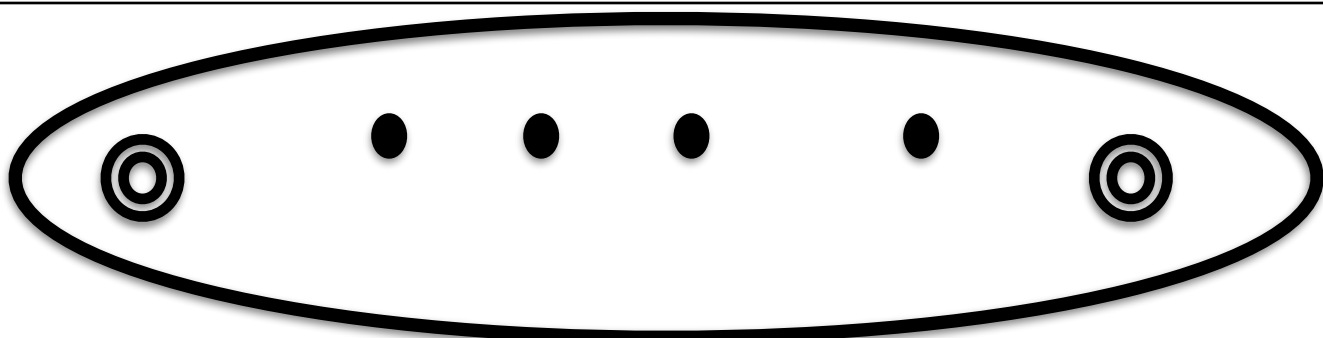
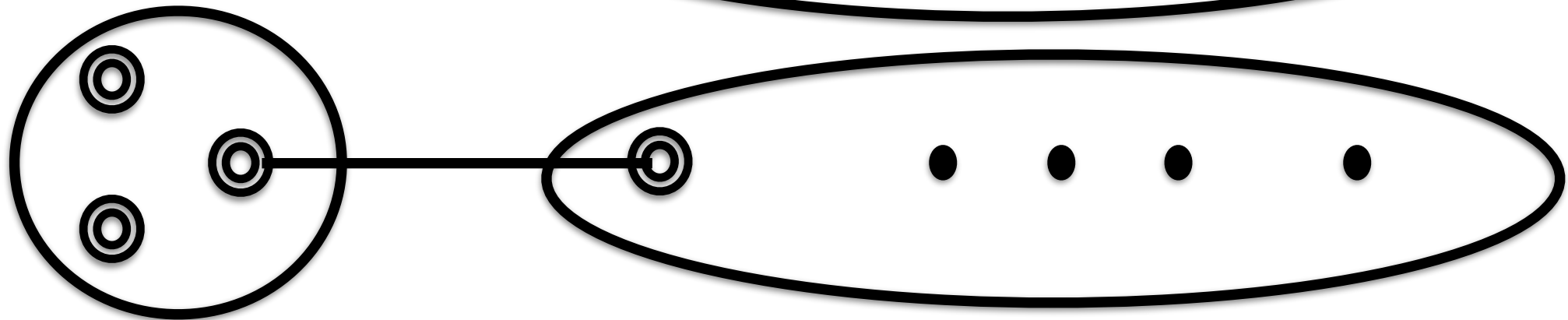
$SU(N)^{M-1}$  quiver in 4d

=



A sphere with two full punctures and  $M$  simple punctures  
in **class S theory**.

# Gaiotto duality

$$\begin{aligned}
 SU(N)^{M-1} &= \text{Diagram 1} \\
 &= \text{Diagram 2} \\
 &= T_N(\text{full}, \text{full}, Y) - SU(M) - SU(M-1) - \dots - SU(1)
 \end{aligned}$$



A relative of the so-called  $T_N$  theory appears.  
 (Just for simplicity I have taken  $N > M$  .)

$T_N(\text{full}, \text{full}, Y)$  is the theory with two full punctures and one puncture named  $Y$  . (Explicitly  $Y = [N - M, 1^M]$  )

# Mass deformation

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??? (obtained from the 5d SCFT  
on circle compactification)

↓ mass deformation of  $SU(M)^2$

$$T_N(\text{full}, \text{full}, Y) - SU(M) - SU(M-1) - \cdots - SU(1)$$

**A fact:**

$$T_M(\text{full}, \text{full}, \text{full})$$

↓ mass deformation of  $SU(M)^2$

$$[M] - SU(M-1) - \cdots - SU(1)$$

[Bergman, Zafrir, 2014]

[Hayashi, Tachikawa, KY, 2014]

# The results

---

??? (obtained from the 5d SCFT  
on circle compactification)

$$= T_N(\text{full}, \text{full}, Y) - SU(M) - T_M(\text{full}, \text{full}, \text{full})$$

It turns out that the gauge group  $SU(M)$  is infrared free by computing the beta function.

The two isolated SCFTs are connected by an IR free gauge group.

# The results

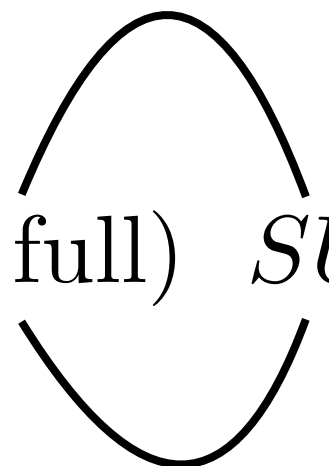
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6d SCFT  $\bigcirc N \bigcirc N \bigcirc N \bigcirc N$   $G = A_{M-1}$



[5d SCFT] + [5d SYM of type  $G = \text{SU}(M)$ ]



$$T_N(\text{full, full, } Y) - \text{SU}(M) - T_M(\text{full, full, full}) \quad \text{SU}(M)$$


The  $\text{SU}(M)$  is coupled to the diagonal subgroup of the  $\text{SU}(M)^2$  of two full punctures.

# The results

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The result can also be generalized to

$$\bigcirc_{N_1} \bigcirc_{N_2} \bigcirc_{N_3} \bigcirc_{N_4} \bigcirc_{N_5} \quad G = A_{M-1}$$

[Gaiotto, Tomasiello 2014]



on  $T^2$

$$T_N(Y_L, Y_R, Y) - SU(M) - T_M(\text{full}, \text{full}, \text{full}) \quad \begin{array}{c} \text{---} \text{---} \text{---} \\ \text{---} \text{---} \end{array} SU(M)$$

- The punctures become more general.
- The case of  $N < M$  is also obtained in the same way.

## Conclusion:

The 4d theory is not exactly class S, but is described by using class S theories as building blocks.

# Contents

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1. Introduction

2. A class of  $N=(1,0)$  generalizing  $N=(2,0)$  theories

3. Compactification on a circle

4. Compactification on a torus

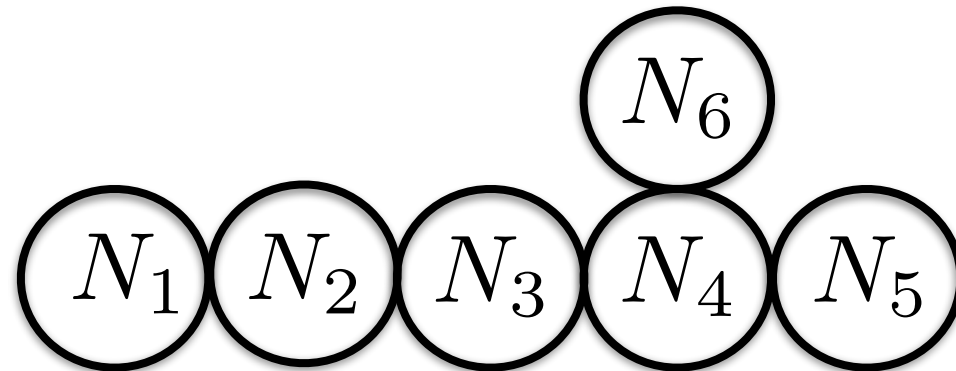
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6. Summary

# 6d SCFTs in F-theory

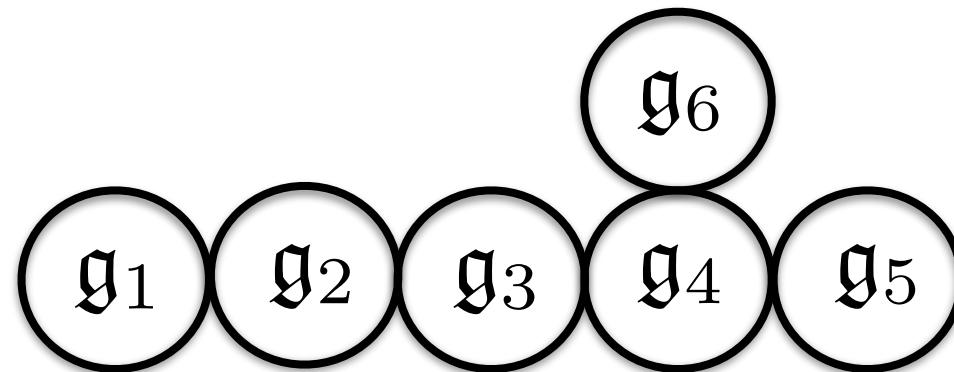
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We can wrap D7-branes on 2-cycles in Type IIB.



Type IIB  $\rightarrow$  F-theory

D7-branes  $\rightarrow$  singular Kodaira fibers with gauge algebra  $\mathfrak{g}$

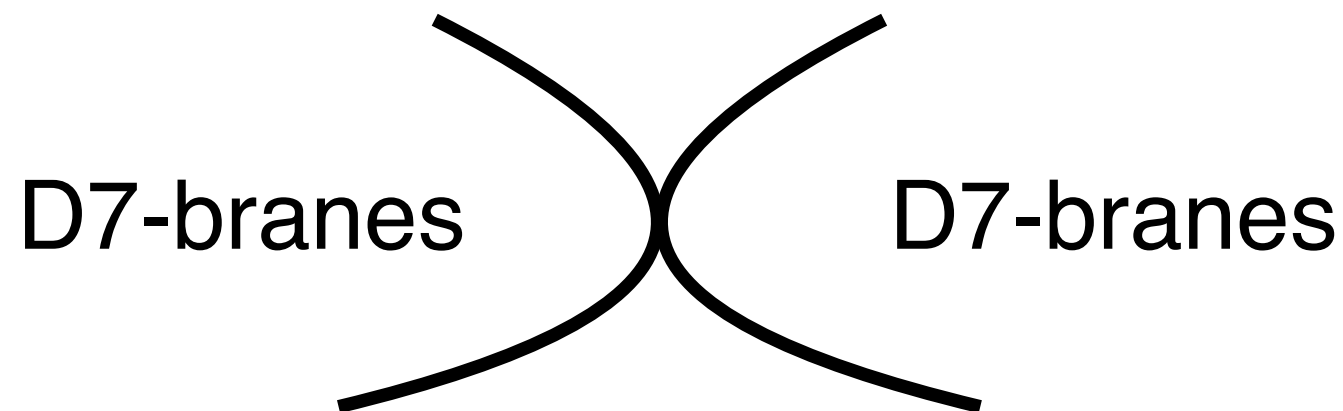


We focus on the case in which the intersection matrix of the 2-cycles is the same as  $N=(2,0)$  of type  $G=ADE$ .

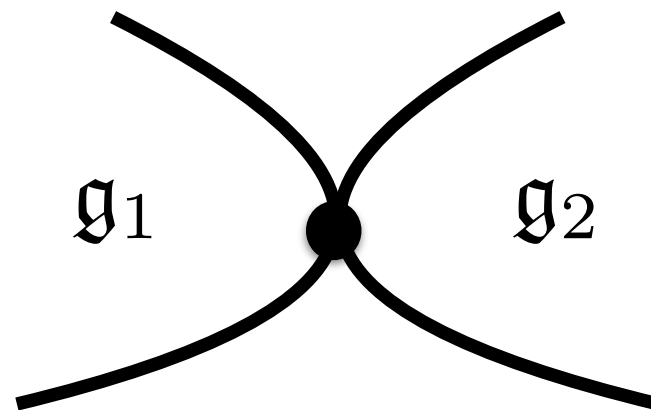


# 6d SCFTs in F-theory

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Bifundamental matters appear at the intersection of the D7-branes.

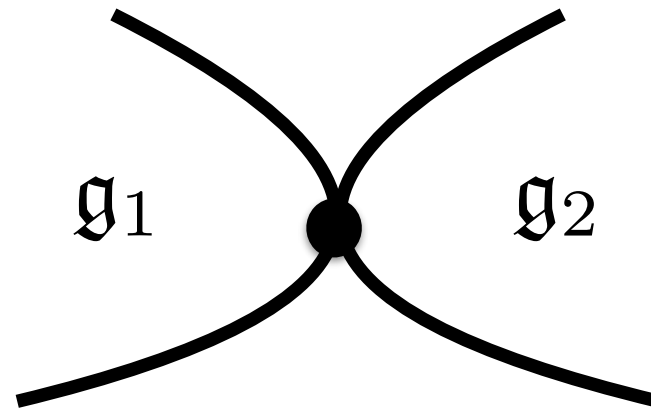


In F-theory, when two singular fibers intersect, we get strongly interacting 6d SCFTs, called **conformal matter**.

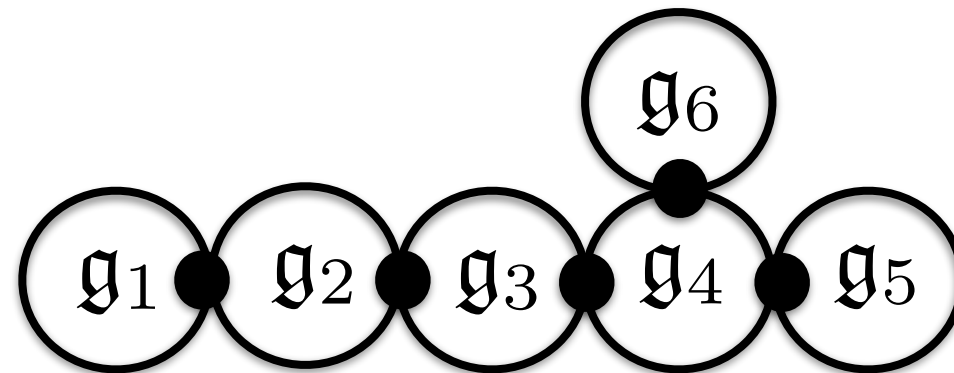
[Del Zotto, Heckman, Tomasiello, Vafa, 2014]

# Conformal matter

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Conformal matter has at least  $\mathfrak{g}_1 \times \mathfrak{g}_2$  flavor symmetries and it is regarded as a **generalized bifundamental** in quivers



More generally, they are **very-Higgsable theories**, but I will not talk about them. See our paper.

[Ohmori, Shimizu, Tachikawa, KY]

# Conformal matter on circle

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$D_N$  conformal matter may be tractable  
for 5d theorists.

$(D_N, D_N)$  conformal matter



on  $S^1$  with **flavor holonomies**

SU(N-2) gauge theory with 2N flavors

[Hayashi, Kim, Lee, Taki, Yagi, 2015]

[KY, 2015]

## Remark:

Without holonomies, the 5d version of the conformal matter is a strongly coupled theory.

# Conformal matter on torus

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$(\mathfrak{g}, \mathfrak{g})$  conformal matter



on  $T^2$

Class S theory of type  $\mathfrak{g}$  with two full punctures and one simple puncture.

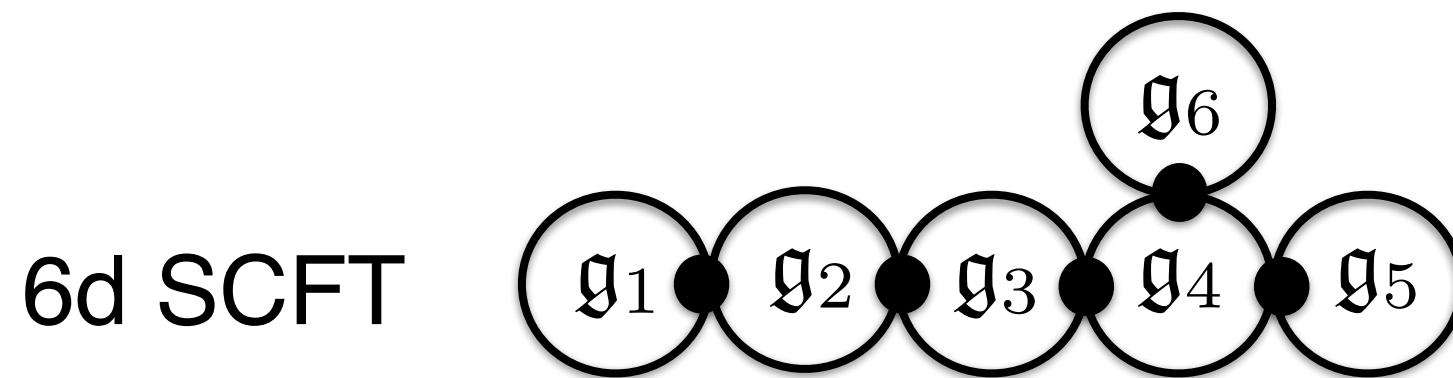
- This is a strongly interacting isolated theory.
- When  $\mathfrak{g} = A$ , it is just an ordinary bifundamental.

[Ohmori, Shimizu, Tachikawa, KY, 2015]

[Del Zotto, Vafa, Xie, 2015]

# General case

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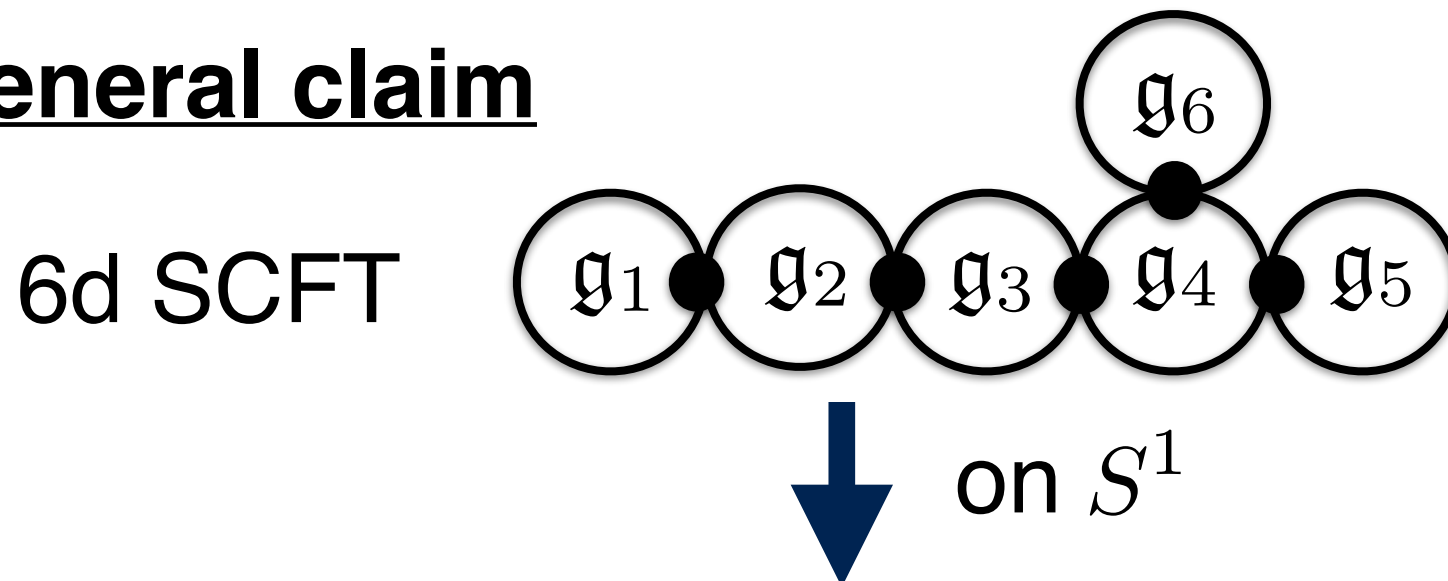


We want to find the compactification of the 6d SCFTs of the class whose base is the  $N=(2,0)$  theory of type  $G$ .

# General proposal

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## Our general claim



[5d SCFT with flavor  $G$ -symm.] + [5d SYM of type  $G$ ]

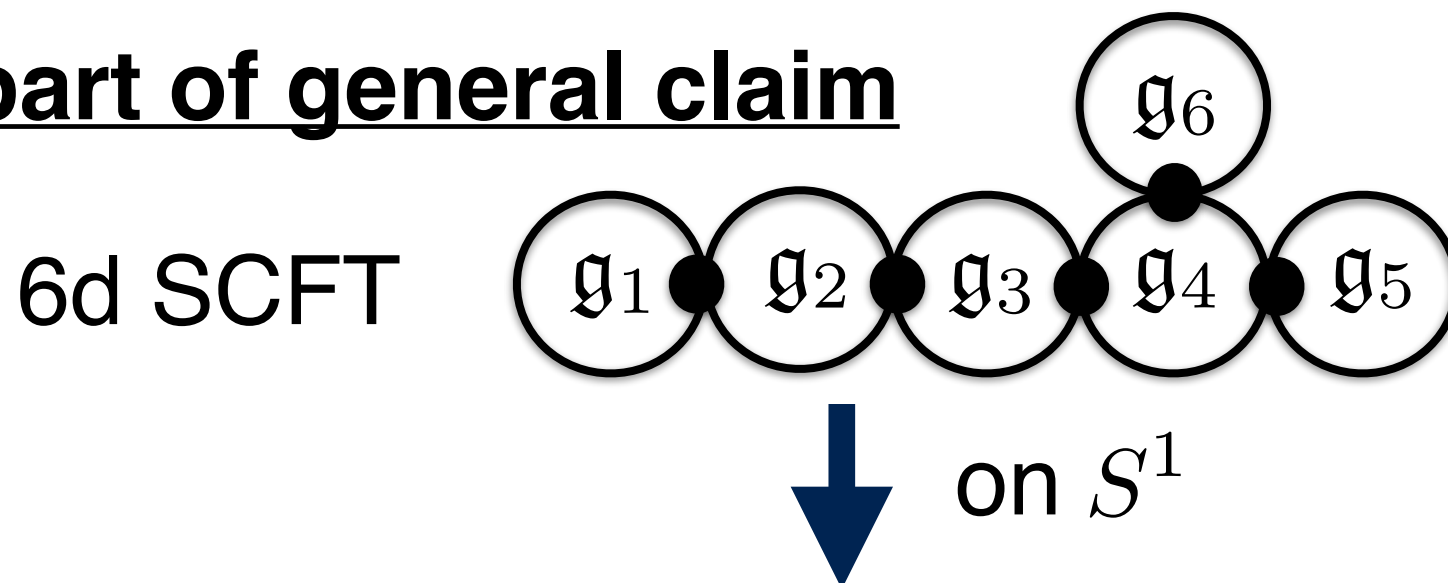


$[4d \text{ SCFT}(1)] - [\text{IR free gauge group}] - [4d \text{ SCFT}(2)]$   
|
[4d SYM of type  $G$ ]

# Justification of general claim (1)

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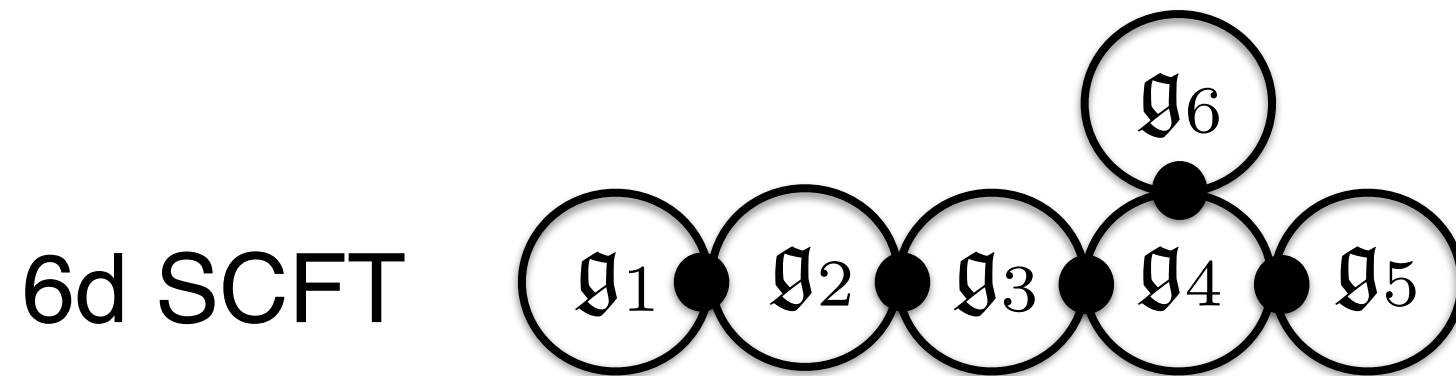
## First part of general claim



[5d SCFT with flavor  $G$ -symm.] + [5d SYM of type  $G$ ]

# Justification of general claim (1)

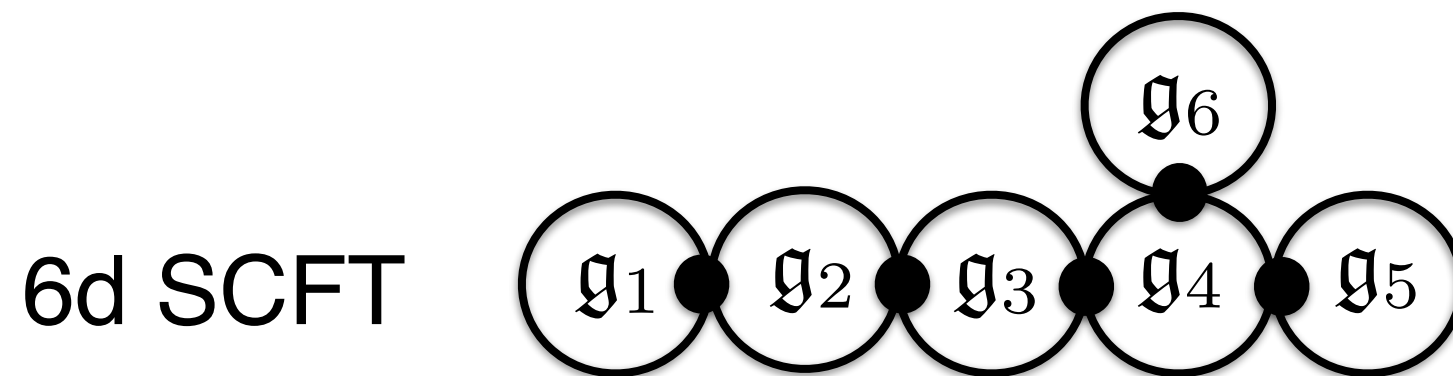
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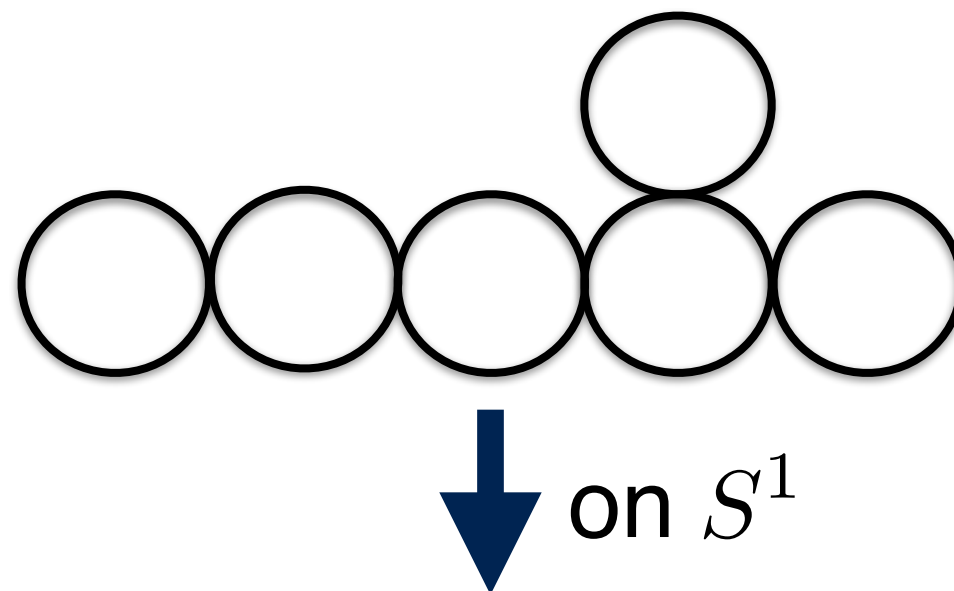


# Justification of general claim (1)

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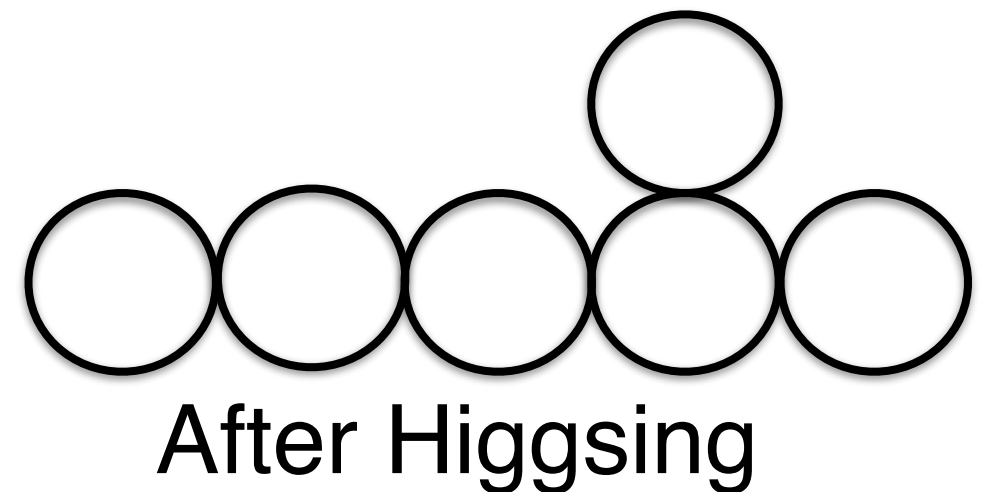
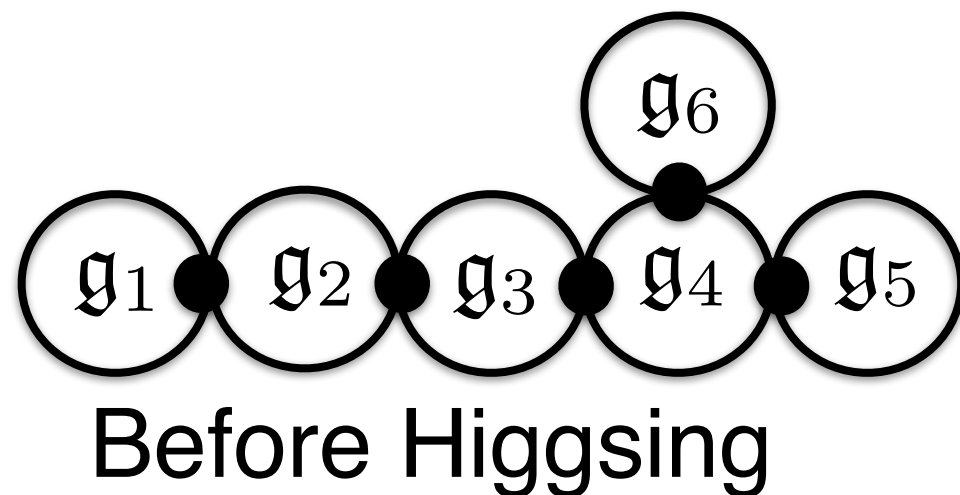
We can always Higgs this theory to the  $N=(2,0)$  theory



[5d  $N=2$  SYM of type G]

# Justification of general claim (1)

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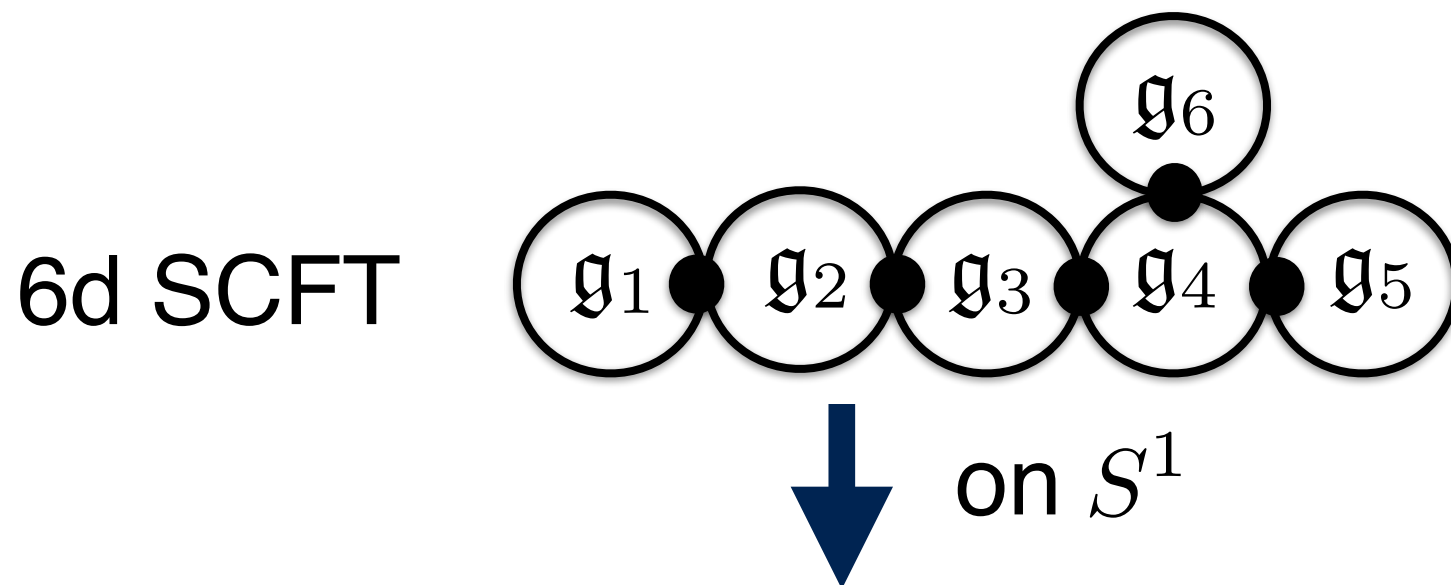
General fact about theories with 8 supercharges:

Higgs branch fields and tensor/vector fields **do not mix**.

If the theory contains 5d SYM after Higgsing, it contains 5d SYM before Higgsing.

# Justification of general claim (1)

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[Something] + [5d SYM of type G]

We conjecture that something is always 5d SCFT.

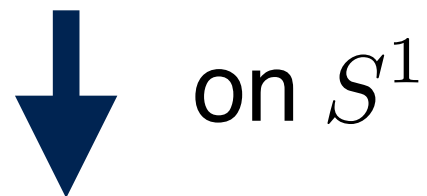
Gauge couplings of  $\mathfrak{g}$ 's  $\rightarrow$  infinity when vevs  $\rightarrow$  zero

# Justification of general claim (2)

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## Second part of general claim

[5d SCFT with flavor G-symm.] + [5d SYM of type G]



[4d SCFT(1)] — [IR free gauge group] — [4d SCFT(2)]  
|  
[4d SYM of type G]

Justification of this claim is more difficult...  
See our paper.

# Implication of General proposal

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For example, 5d version of a generalized quiver like

$$CM_{\mathfrak{g}} - \mathfrak{g} - CM_{\mathfrak{g}} - \dots - \mathfrak{g} - CM_{\mathfrak{g}}$$

where  $CM_{\mathfrak{g}}$  is the conformal matter of type  $(\mathfrak{g}, \mathfrak{g})$

This quiver must have a UV fixed point as a 5d SCFT!  
The UV fixed point has enhanced  $G=\mathrm{SU}(N)$  symmetry.

More studies need to be done.....

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---

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# Summary

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- There is a class of 6d SCFTs which is a certain generalization of the  $N=(2,0)$  theories of type  $G=ADE$ .
- The circle compactification always gives  
[5d SCFT with flavor  $G$ -symm.] + [5d SYM of type  $G$ ]
- The problem of the torus compactification of 6d SCFTs is reduced to the circle compactification of the corresponding 5d SCFTs.

Thank you very much!