

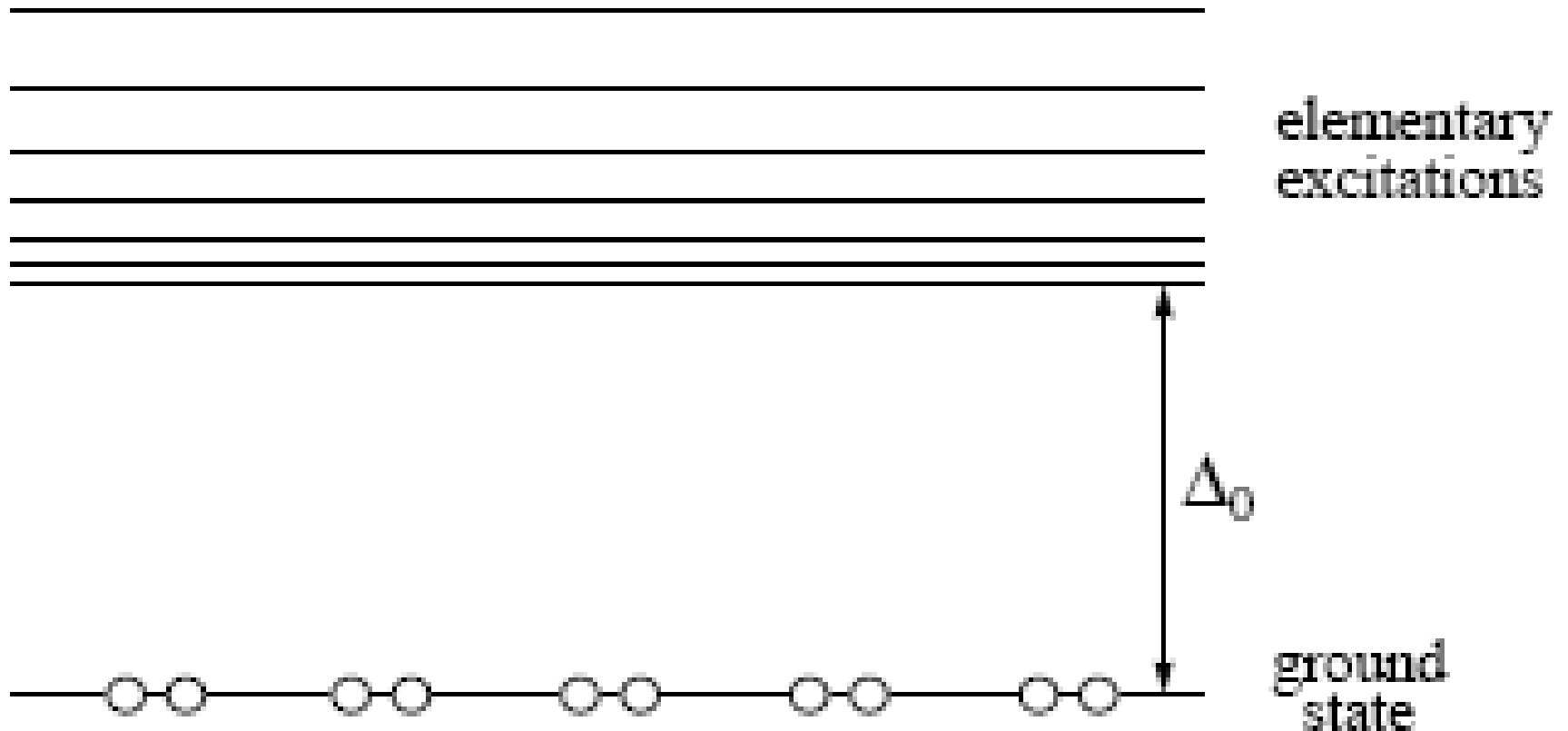
The origin of pseudogap in HTSC

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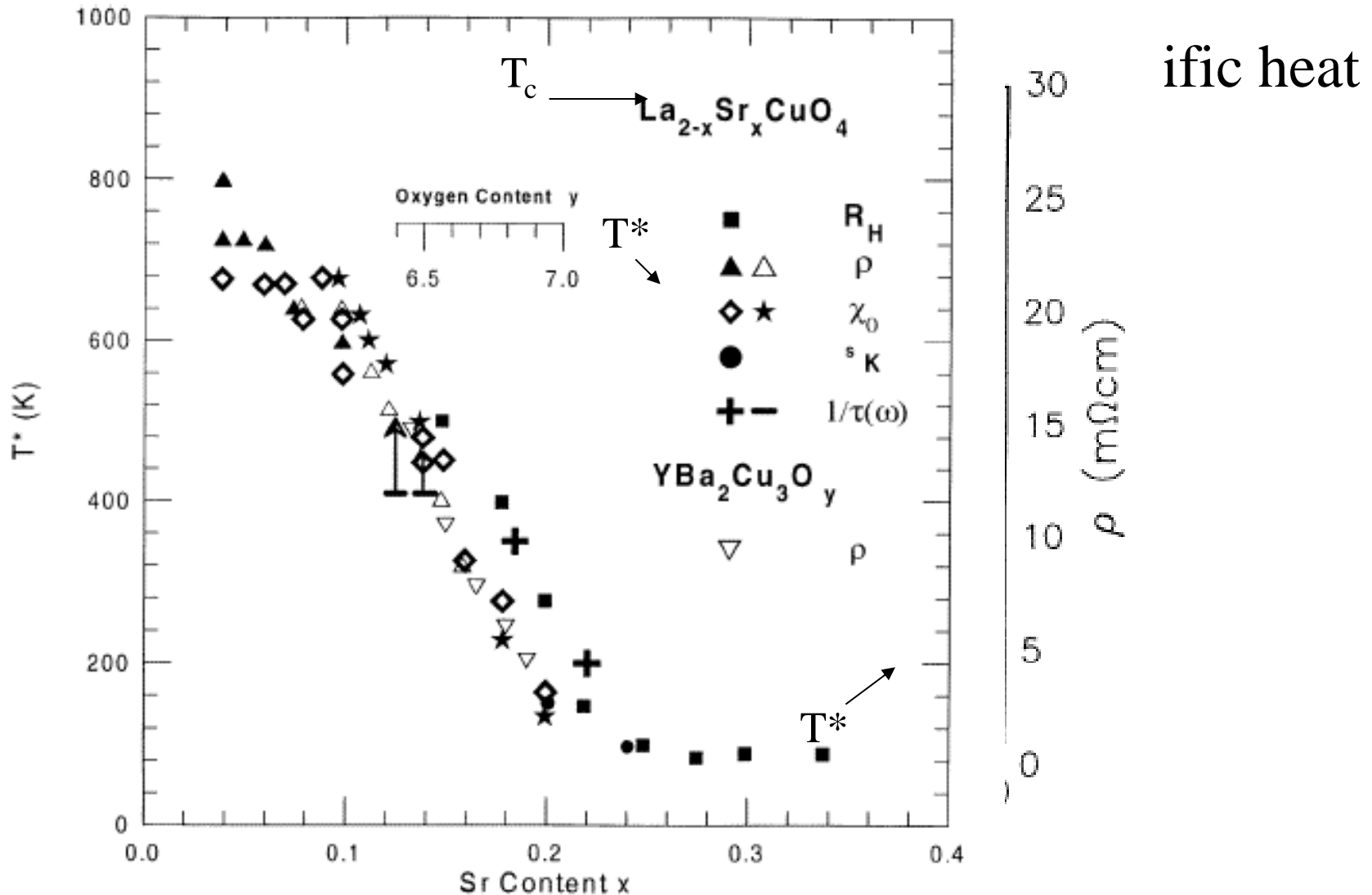
The superconductor energy gap

The BCS superconductor



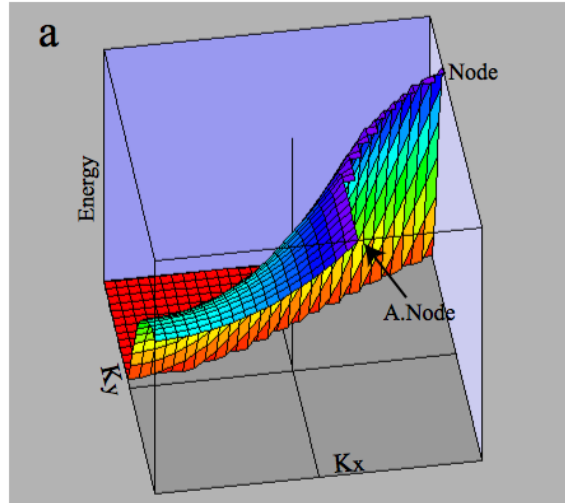
The pseudogap temperature

Timusk Rep. Phys. 62 61-122 1999

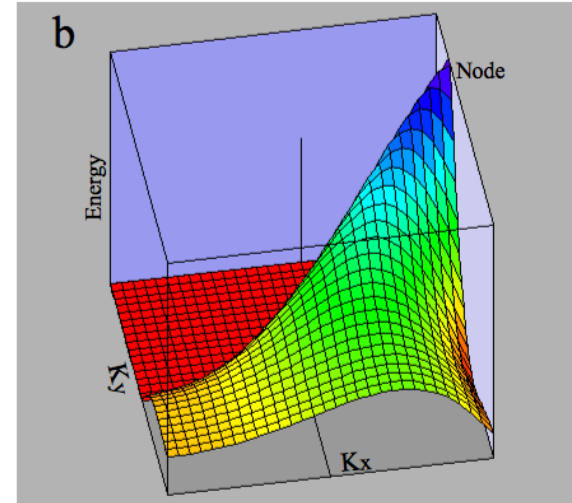


ARPES measurements

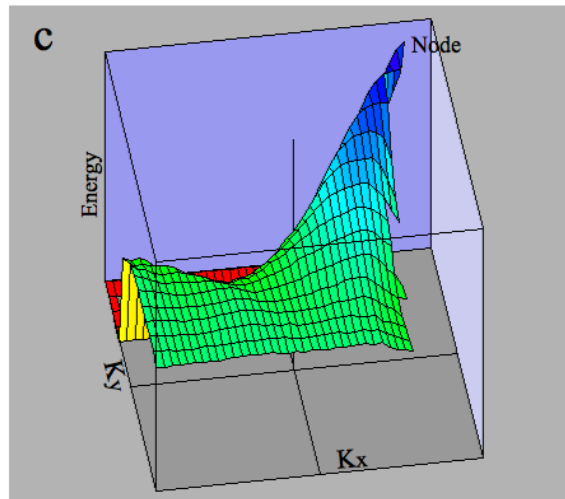
Theory
Normal state



Theory
dSC state

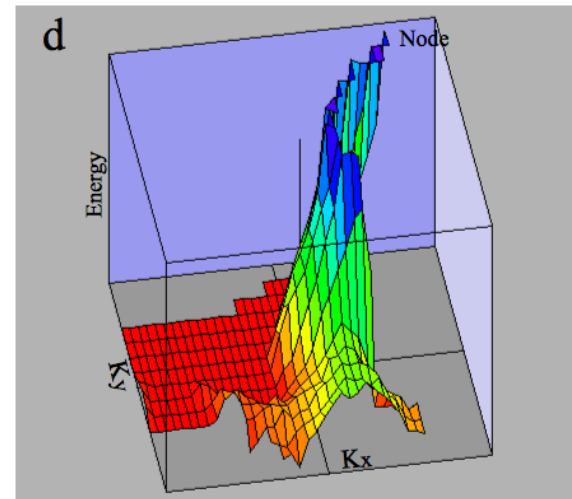


Experiment
dSC



Kanigel

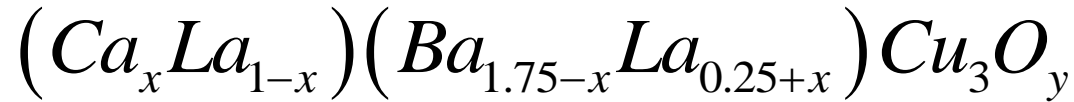
Experiment
PG state



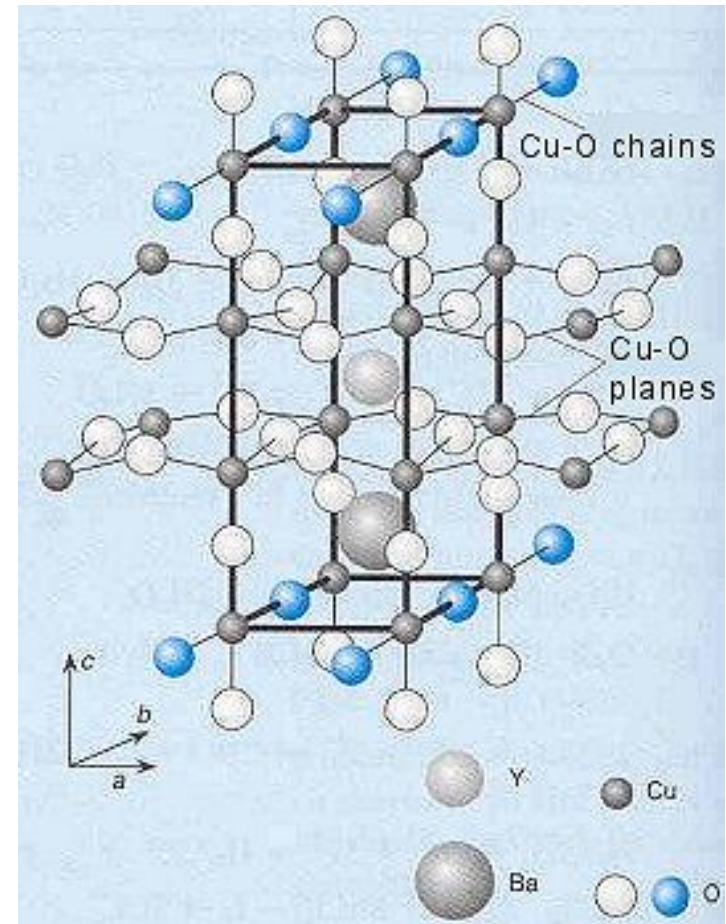
Main question

What are the interactions that affect the T^* ?

The CLBLCO system

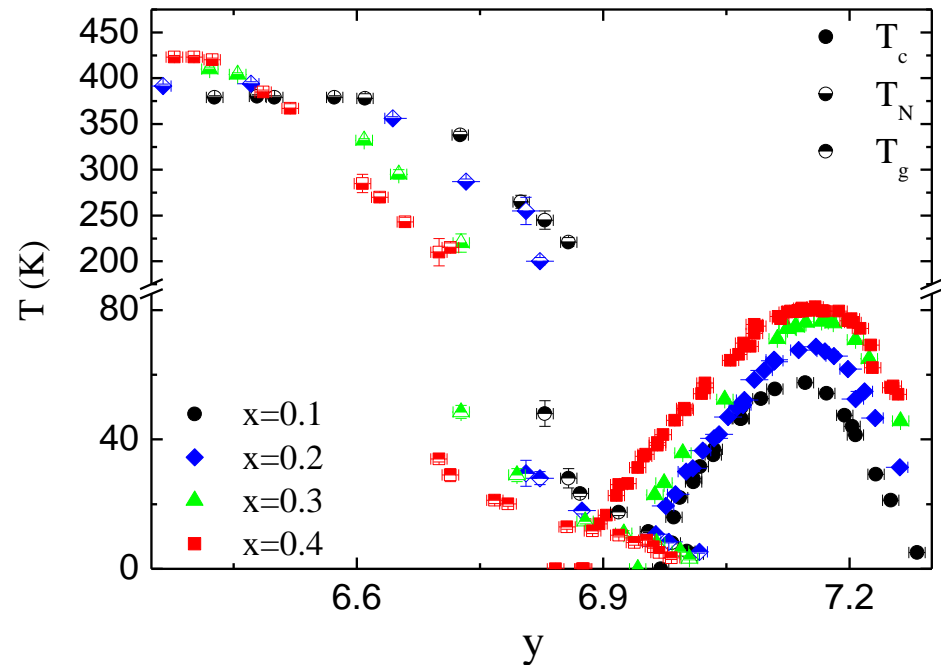


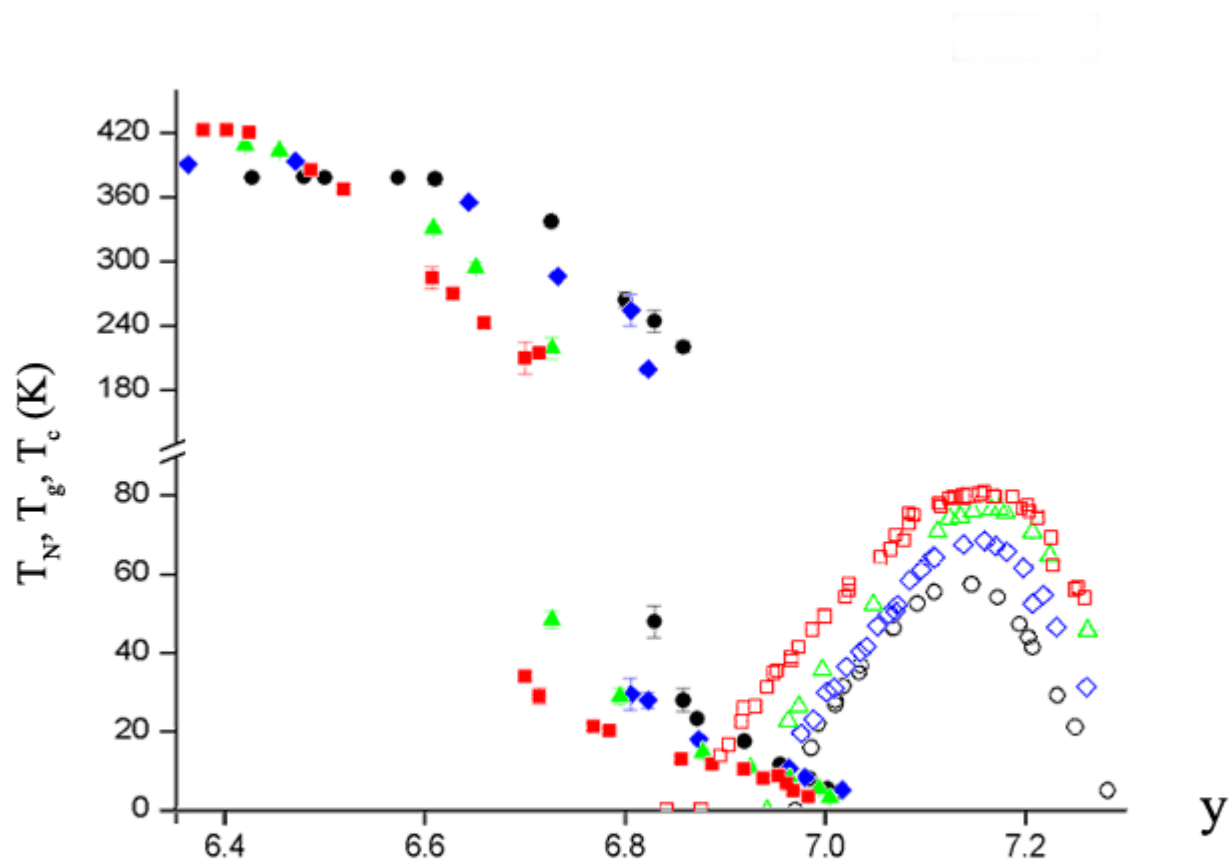
- Similar structure as the well known YBCO
- 1:2:3 atomic ratio
- The main structure doesn't change with the families
- Controllable doping level (y parameter)
- **Controllable magnetic coupling (x parameter)**

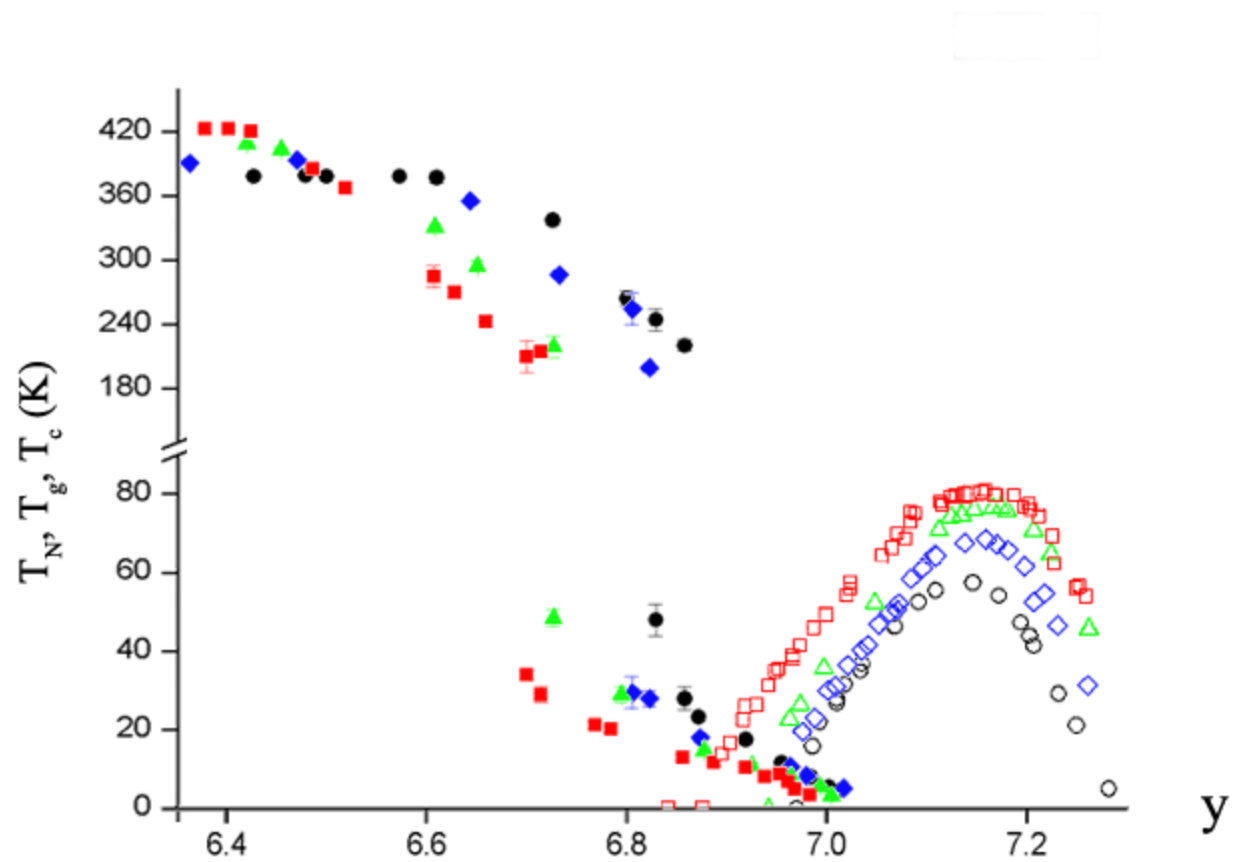


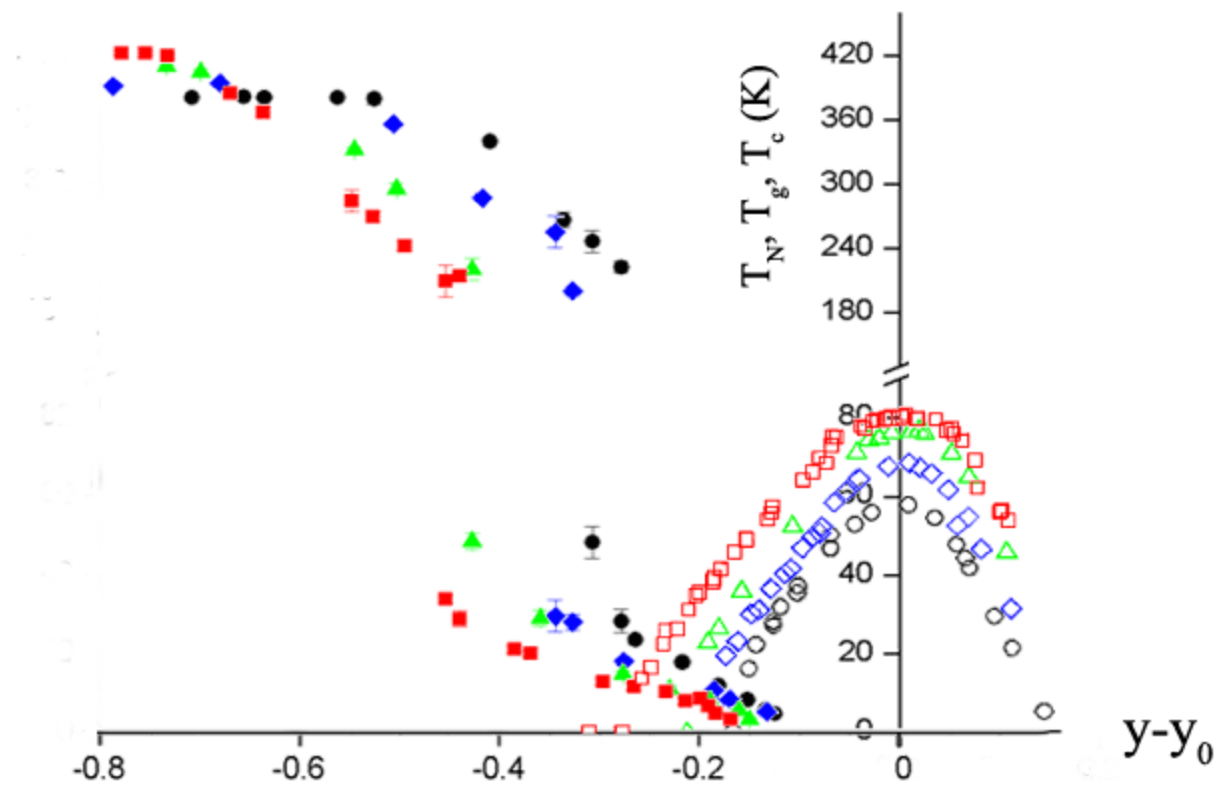
CLBLCO phase diagram

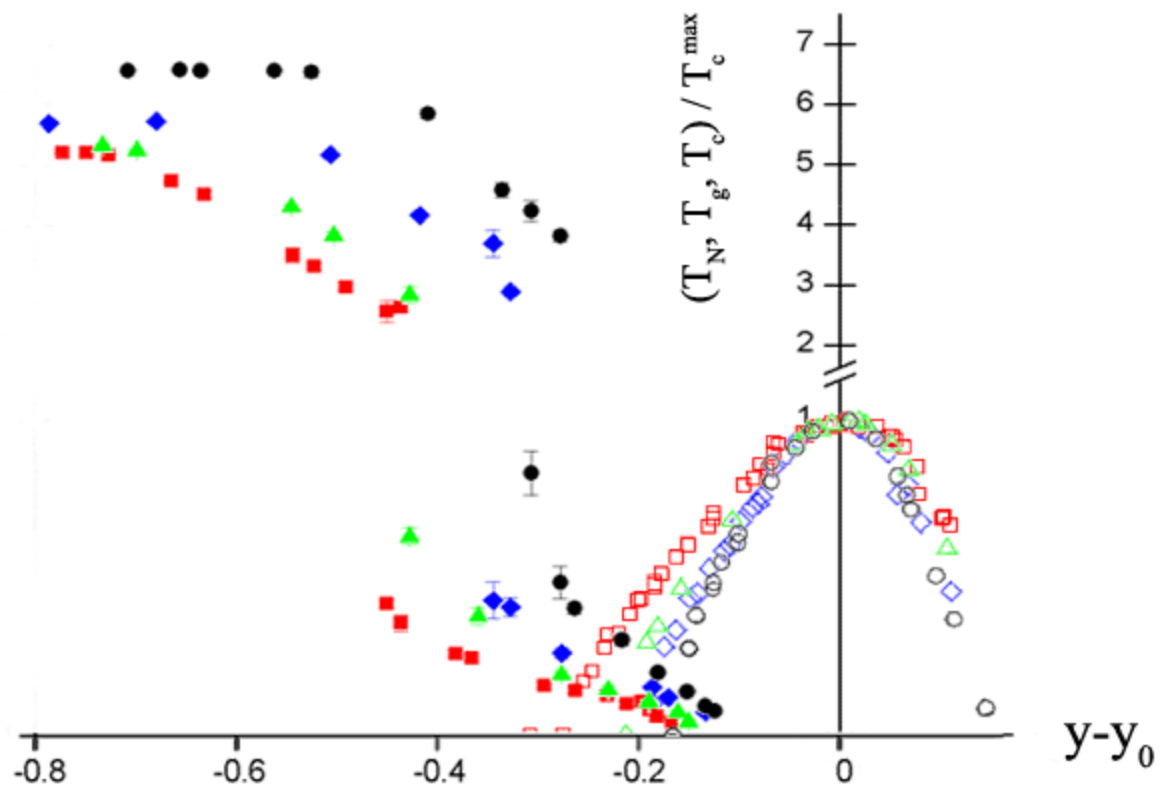
- Similar phase diagrams
- The family with the highest T_c have the highest T_N on the lowest doping.
- Big difference at T_c^{\max} between the families

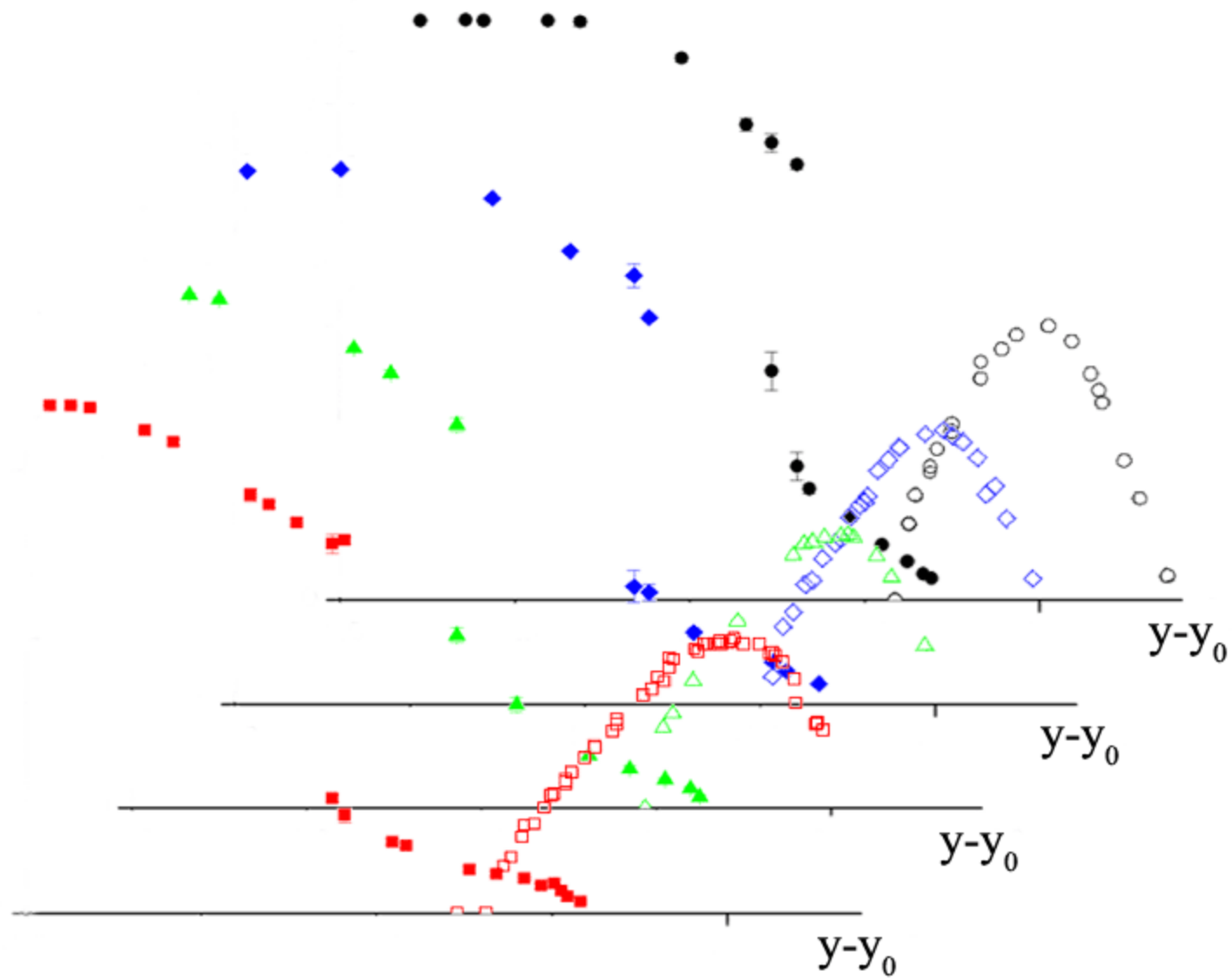


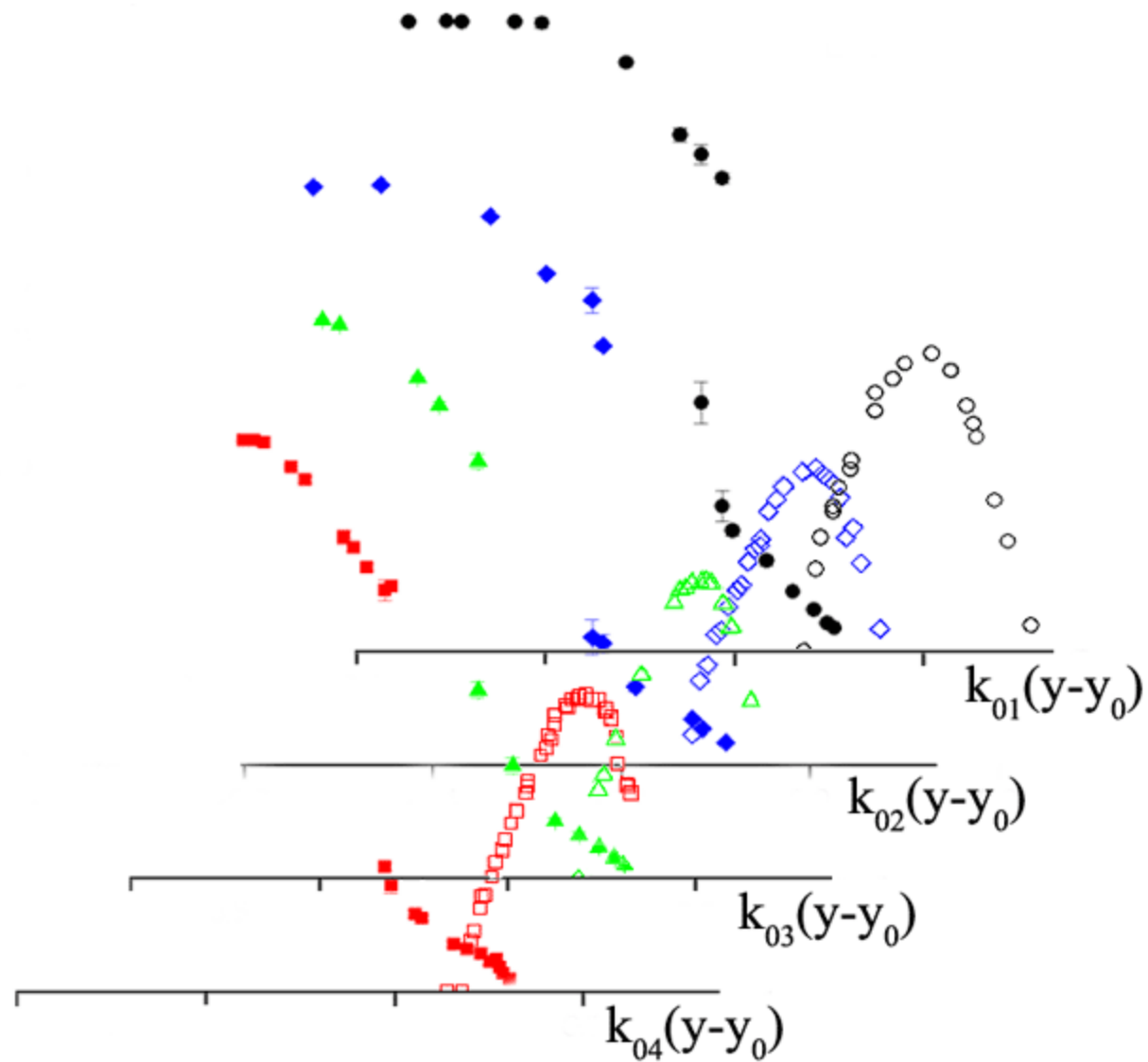


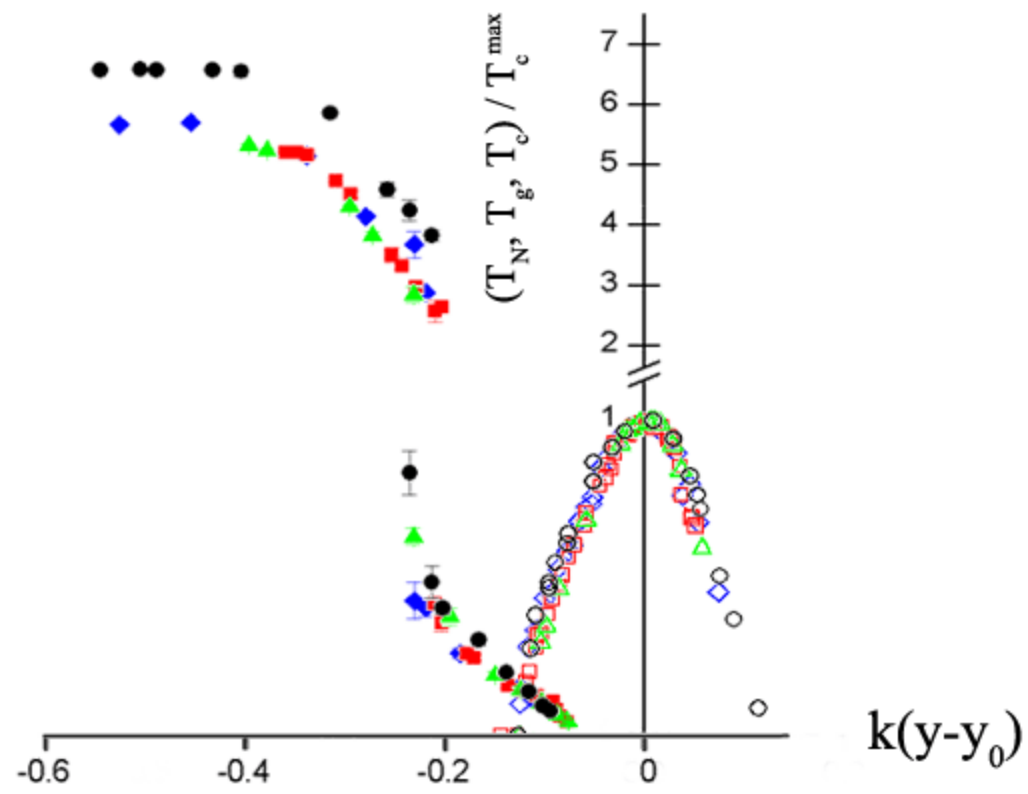






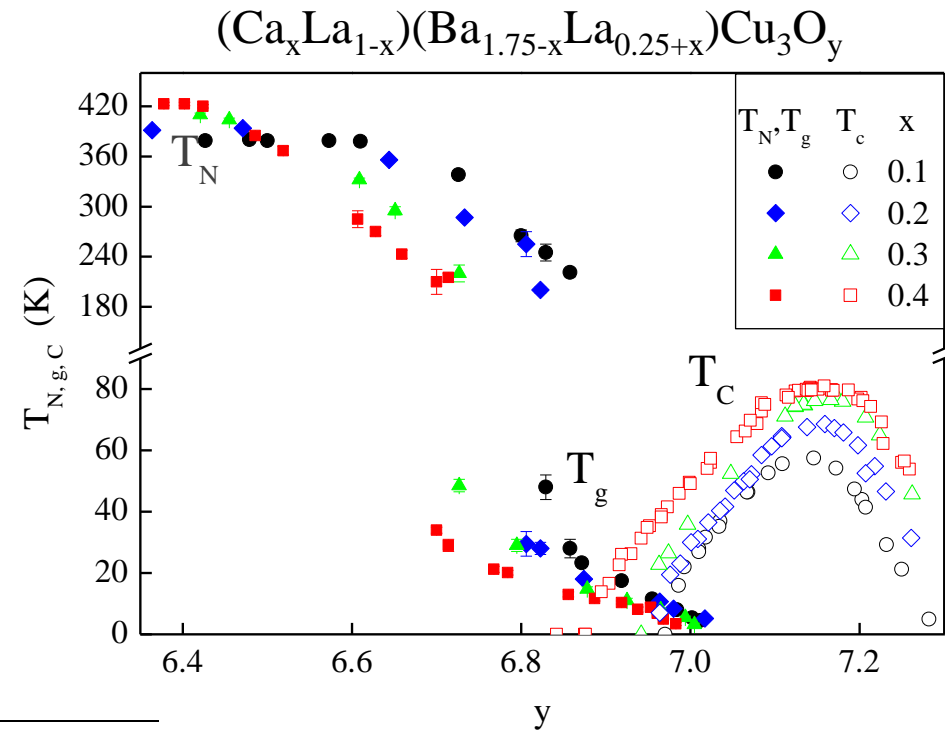
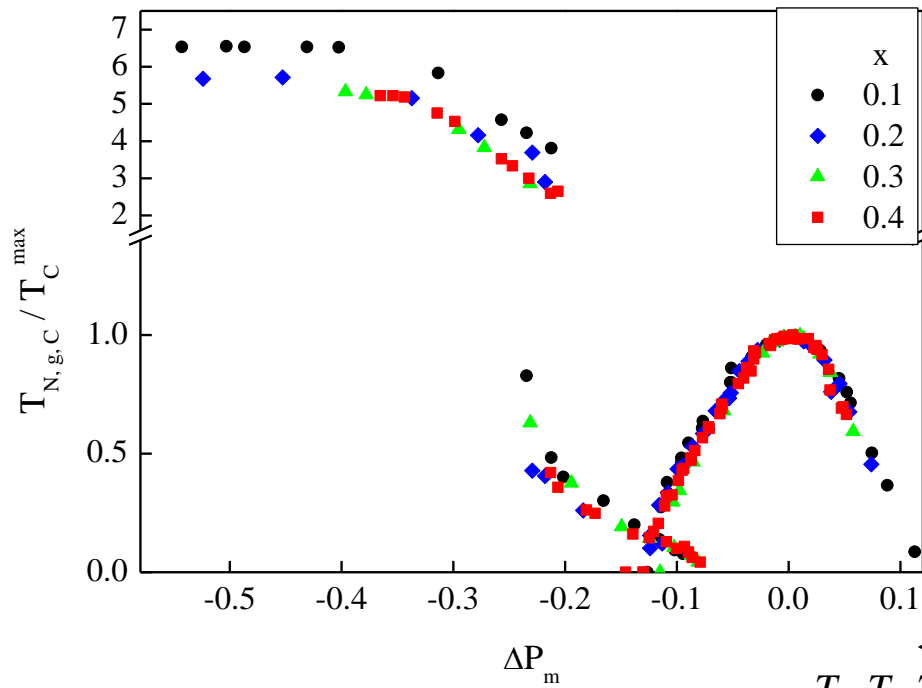






Transformation of the entire doping range.

$$\Delta p_m = p - p^{opt} = K_x (y - y_x^{opt})$$



$$T_N, T_g, T_c \Rightarrow T_N, T_g, T_c / T_c^{max}$$

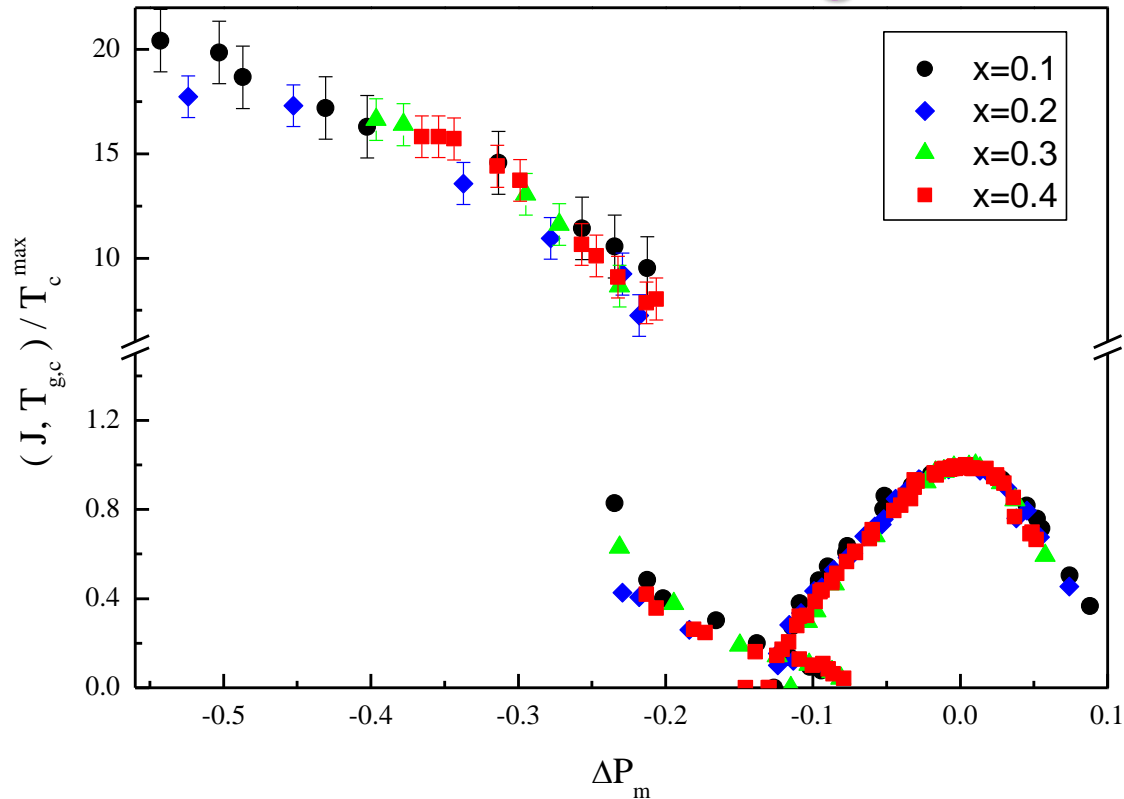
$$y \Rightarrow K(x)\Delta y$$

The scaling works in the entire doping range apart for $x=0.1$?

The role of anisotropies

- T_N is determined by the in-plane J and out of plane J_{\perp} coupling.
- We extracted J out of T_N .

Unified Phase Diagram



The in-plane J is extracted from T_N .

Scaling Conclusion

- We found that T_c scale like the in-plane J therefore is a consequence of a 2D magnetic interaction.

$$T_c \propto J$$

- **Question: Does T^* scales with J as T_c does, or with some other magnetic parameter?**

The experimental methods

- The SQUID
(Superconducting
QUantum Interference
Device)
- The temperature range
is 1.2K to 310K
- The field range is up to
6.5T.



Susceptibility

- Definition
$$\chi_0 = \lim_{H \rightarrow 0} \frac{\partial M}{\partial H}$$

- Practice
$$M = \chi_{dc} H$$

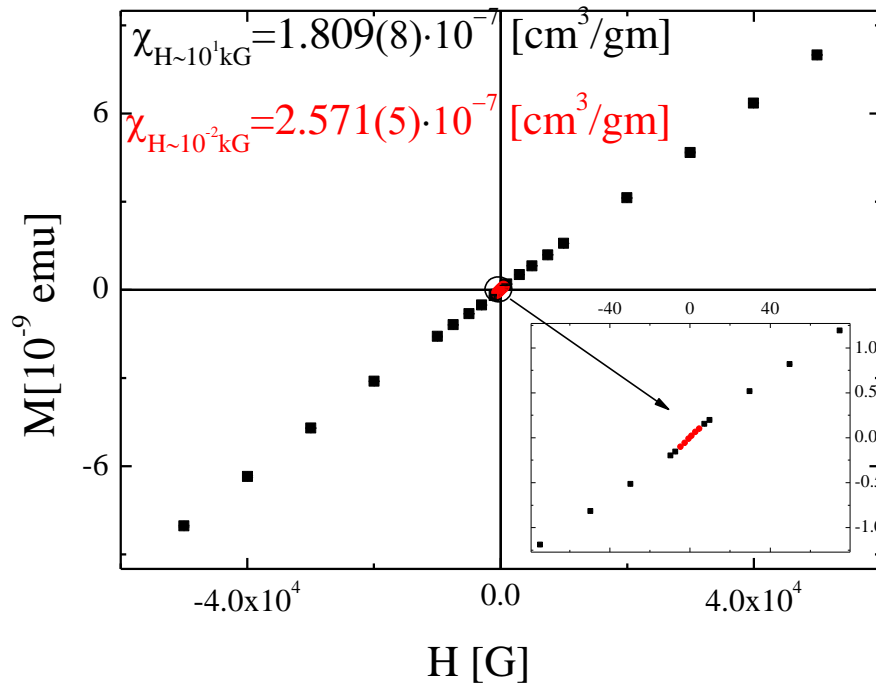
- Where D is known as the demagnetizing factor, and it get different values for different geometries.

$$\chi_{dc} = \frac{\chi_0}{1 + D\chi_0}$$

- For needle like sample $D=0$, then:

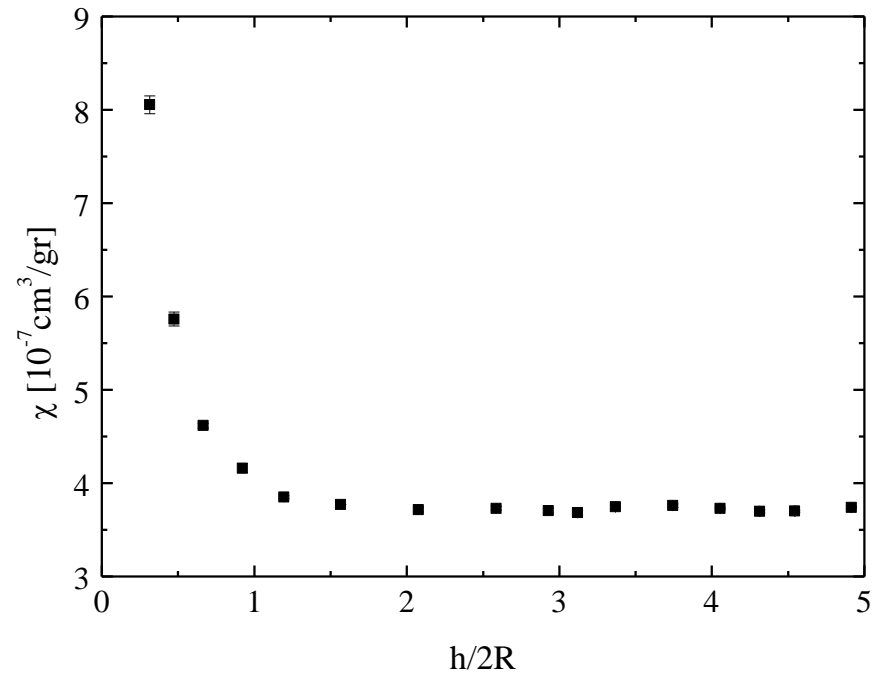
$$\chi_{dc} \approx \chi_0$$

Measurement condition

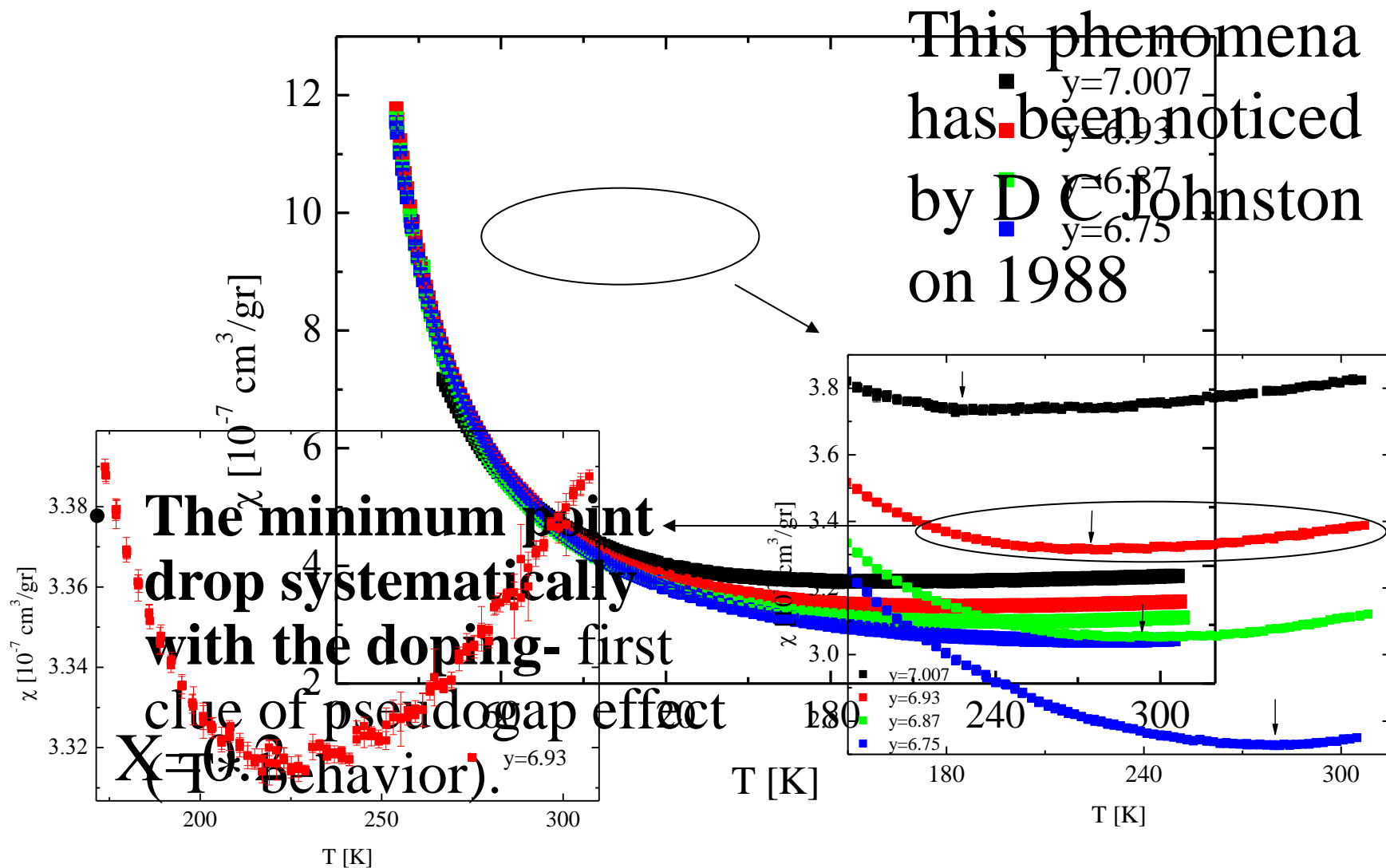


Field dependence

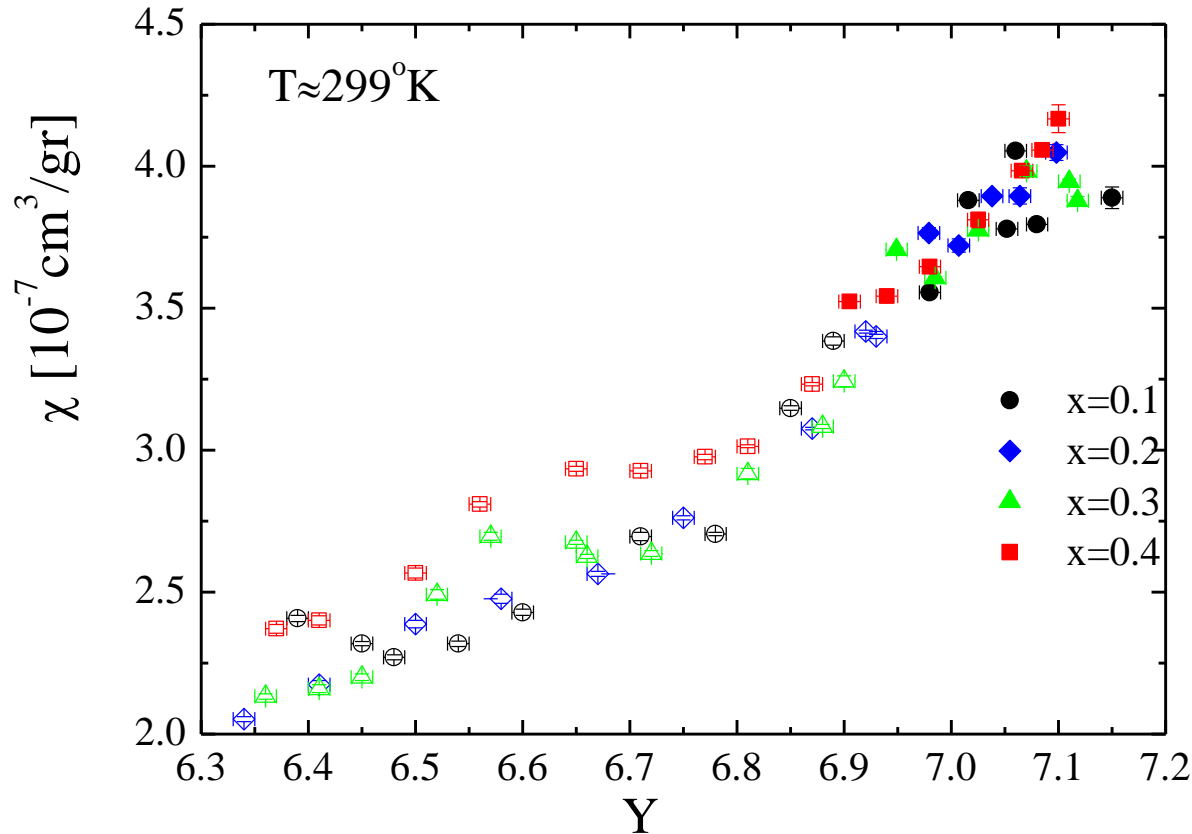
Geometric dependence



Raw data



$$\chi_0 = f(y)$$



- The value of χ is increasing with the doping (Pauli susceptibility).

Susceptibility types

- Isolated spin: Langevin paramagnetism, Curie law

$$\chi_0 = \frac{N \mu_B^2}{3k_B T} = \frac{C}{T}$$

- Weakly coupled spins: Curie-Weiss

$$\chi_0 = \frac{C}{T + \theta}$$

- Pauli spin (Landau) :

$$\chi_0(T) = \text{const} = \mu_B^2 \mathbf{D}(\varepsilon_f)$$

- Core: Van Vleck and Langevin

$$\chi_0(T) = \text{const}$$

There is no traditional theory about increasing susceptibility with T

Strongly coupled spins

- Two coupled spins according to Heisenberg model.

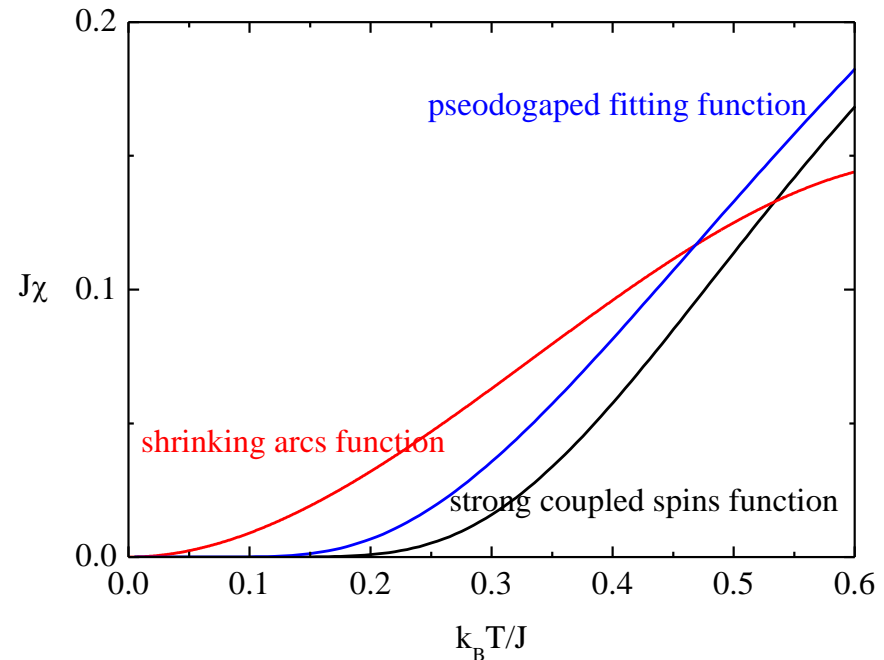
$$\chi_0 = \beta \left[e^{\frac{\beta J}{2}} \cosh\left(\frac{\beta J}{2}\right) \right]^{-2}$$

- shrinking arcs phenomena.

$$\chi_0 = A(T) \left(\frac{2T}{T^*} - \left[\frac{T}{T^*} \right]^2 \right)$$

- The fitting term.

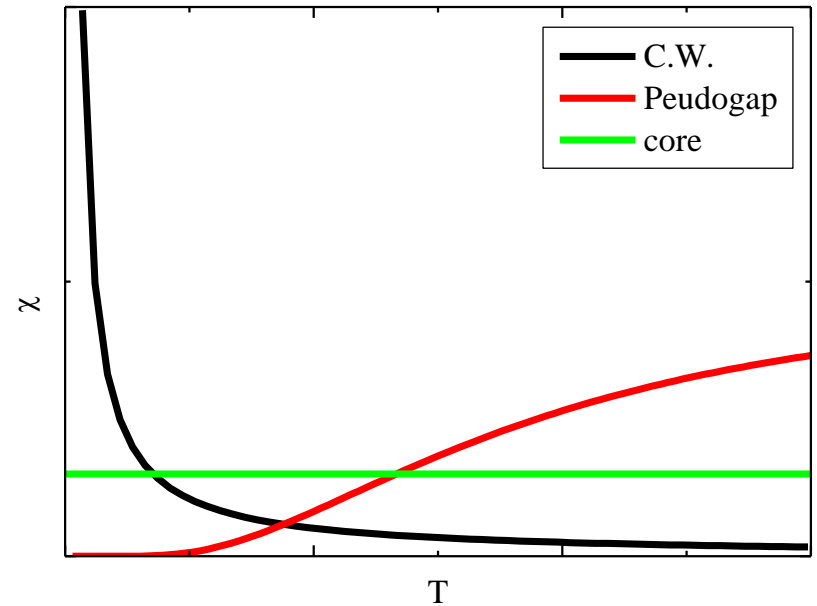
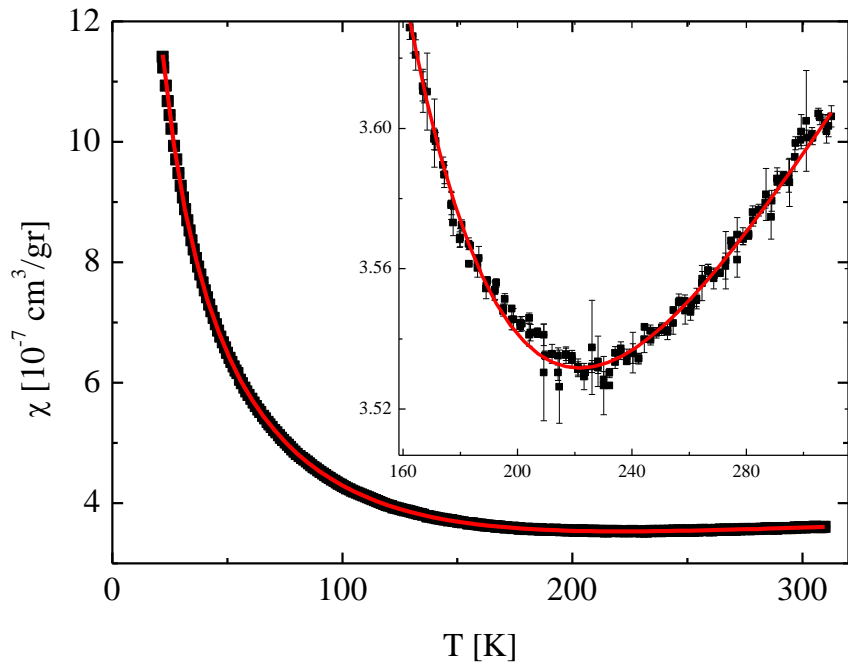
$$\chi_0 = \frac{const}{\cosh\left(\frac{T^*}{T}\right)}$$



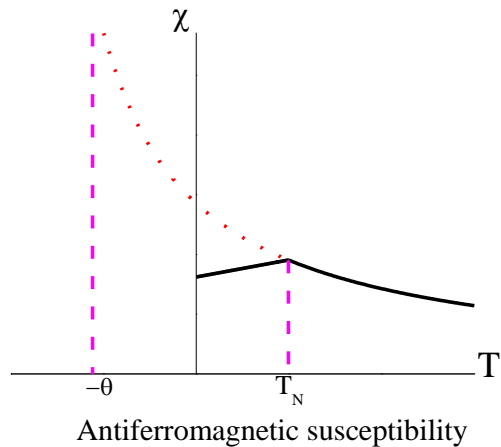
The fitting function

$$\chi_0 = \frac{C_1}{T + \theta} + \frac{C_2}{\cosh\left(\frac{T^*}{T}\right)} + C_3$$

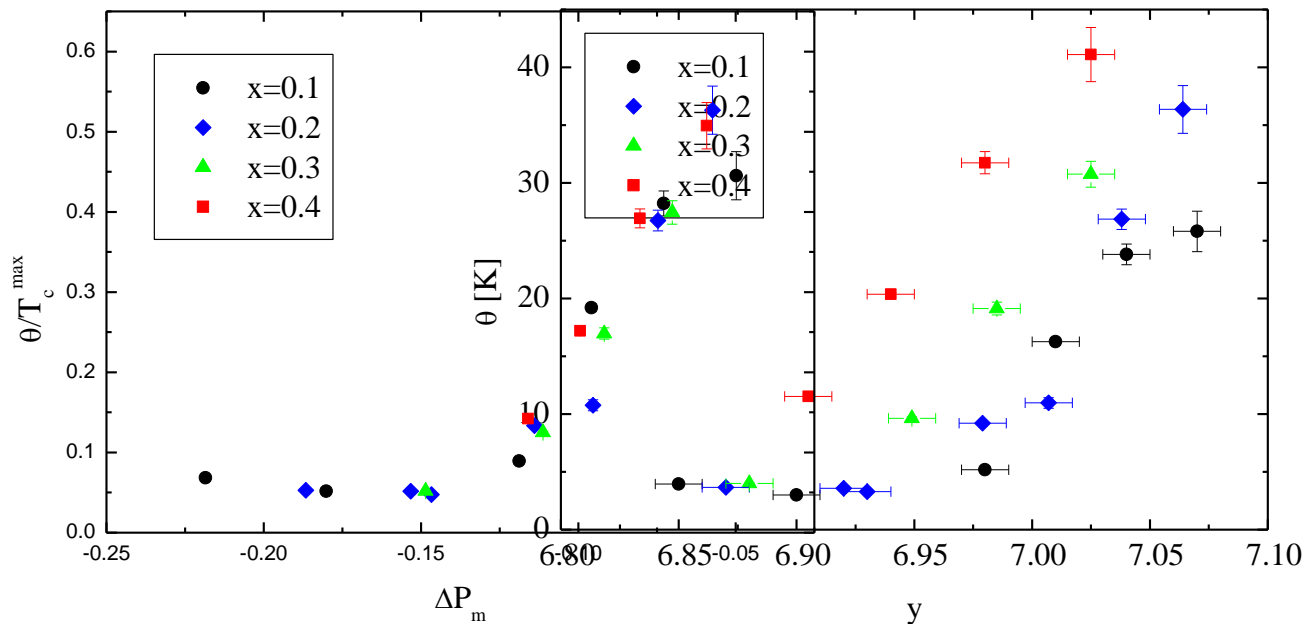
C.W. + PG + CORE



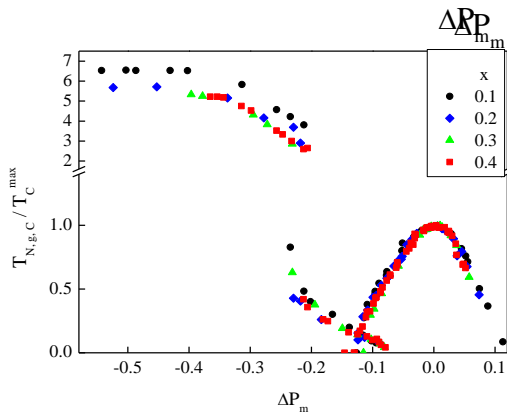
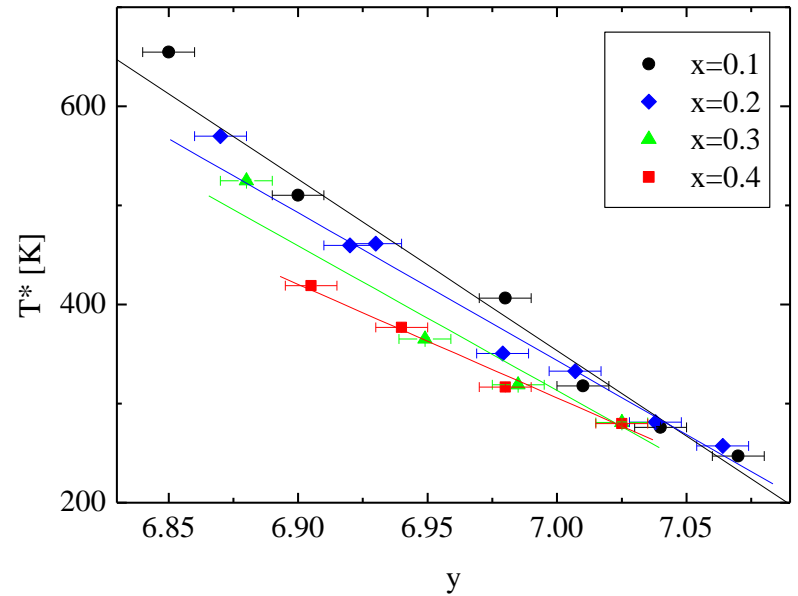
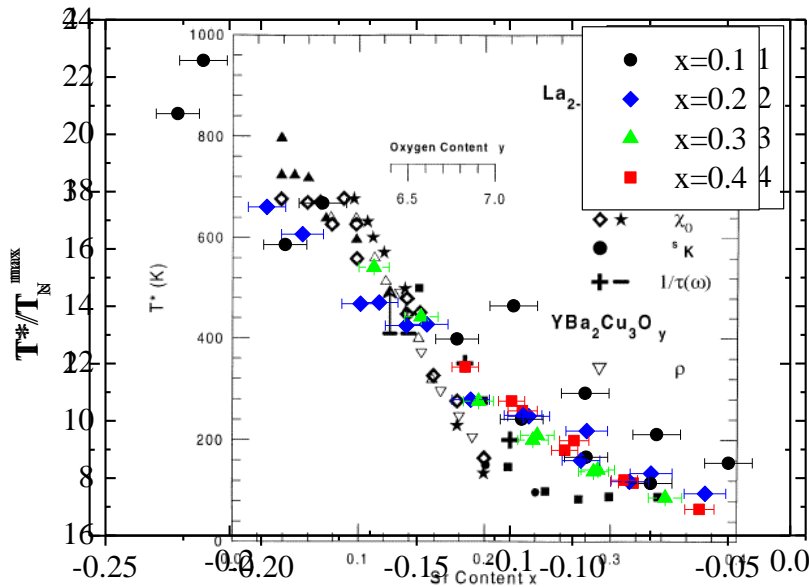
Curie-Weiss temperature



$$\theta = \left[\frac{2S(S+1)}{3K_B} \right] \sum_i Z_i J_i$$



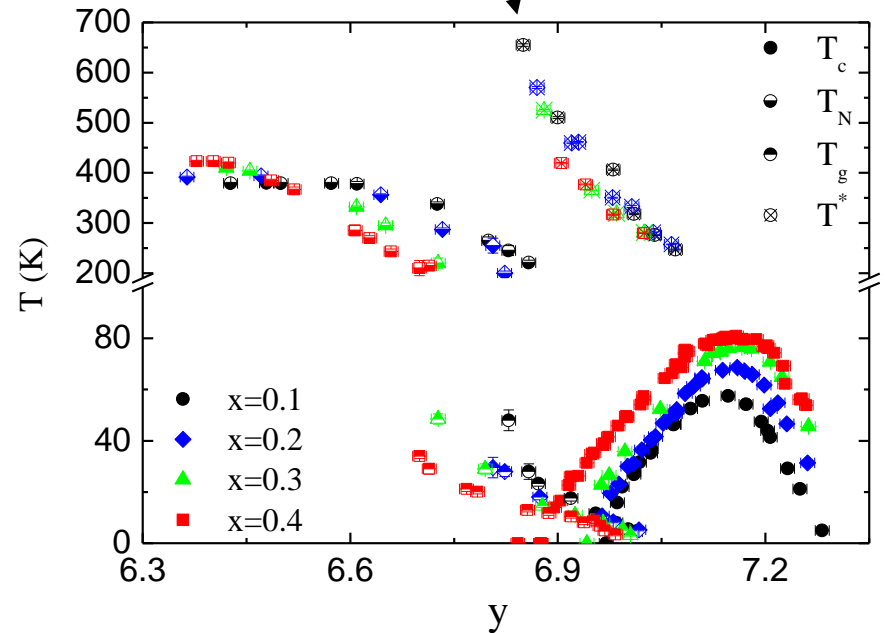
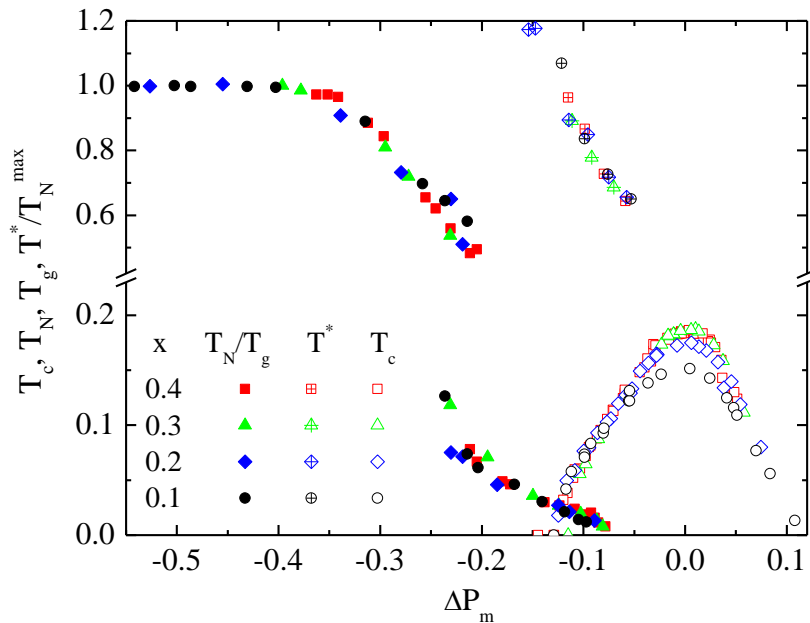
T^*



T^* The T^* doesn't scale well with T_C .
 scale with T_N
 Very similar to the T_C/T_N scaling.

Conclusions

We added the T^* to the phase diagram
 T^* scale like T_N , and it is a 3D magnetic phenomena.
 T_c is a 2D magnetic phenomena.



Acknowledgment

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