# Outflows in Tidal Disruption Events

Julian Krolik, JHU with Roseanne Cheng, Tsvi Piran, Hotaka Shiokawa, Gilad Svirski

### What Are Tidal Disruption Events?

Conceptual definition:

Star passes within tidal radius of a supermassive black hole; Much of its material eventually accreted onto the black hole

Operational definition:

Optical or UV or X-ray flare; Generally caught while declining; Detectable ~few weeks — ~year; In a galactic center; Distinguishable from a supernova

### Best Observed Example: ASASSN 14li



Brown et al. 2017



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## TDEs Are Just Like AGN

- Accretion onto a supermassive black hole is the basic engine
- Expect T ~ few x  $10^4$  few x  $10^5$  K
- If black hole rotates, why not a jet?

## TDEs Are **NOT** Like AGN

- Accretion non-steady, possibly super-Eddington, non-circular, fed relatively close to the black hole
- Missing much of the usual phenomenology: no NLR, obscuring torus, coronal X-rays; no broad CIII], MgII, sometimes no Balmer lines; line widths few x AGN widths, and change over time
- Indications that much of the visible light **not** from local turbulent dissipation

## **Basic Mechanics**

• Tidal radius from density/frequency matching

 $R_T \simeq R_* (M_{\rm BH}/M_*)^{1/3}$ 

 $R_T \simeq 15[(k/f)/0.08]^{1/6} (M_*/M_{\odot})^{2/3-\xi} M_{\rm BH,6.5}^{-2/3} R_g$  (main sequence)

- Number of stars with <R> this small << 1 —> stars come from far out on nearly-parabolic orbits
- Within R<sub>T</sub>, "independent fluid elements"

Half stellar mass bound, half unbound  $\max |E| \sim GM_{\rm BHI}R_*/R_T^2$ 

### **Basic Mechanics**

Most-bound energy implies

 $a_{\min} \sim R_T (M_{\rm BH}/M_*)^{1/3} \sim 2000[(k/f)/0.08]^{1/6} (M_*/M_{\odot})^{1/3-\xi} M_{\rm BH,6.5}^{-1/3} R_g$  $t_0 \equiv P_{\min} \simeq 20[(k/f)/0.08]^{1/2} (M_*/M_{\odot})^{(1-3\xi)/2} M_{\rm BH,6.5}^{1/2} d$  $e_{\min} \simeq 1 - 2(M_{\rm BH}/M_*)^{1/3}$ 

- Most-unbound energy implies  $v_\infty \simeq 11,000[(k/f)/0.08]^{-1/6}(M_*/M_\odot)^{-1/6+\xi/2}M_{\rm BH,6.5}^{1/6}~{\rm km/s}$
- Lack of another energy scale implies  $dM/dE \simeq const.$  for  $-\max |E| \le E \le +\max |E|$

### Consequences for Stellar Debris

Mass-return rate rises to  $\sim M*/(3t_0)$  at t  $\sim t_0$ 

 $\max\left(\dot{M}_{\rm return}\right) \simeq 60(\eta/0.1)[(k/f)/0.08]^{-1/2}(M_*/M_{\odot})^{(1+3\xi)/2}M_{\rm BH,6.5}^{-3/2}$ 

Mass-return rate then falls ~  $(t/t_0)^{-5/3}$ 

But mass-return rate is **NOT** the same as mass-accretion rate

 $E_B(a_{min}) << E_B(R_T)$  and orbital energy-loss is slow:

Glancing convergence makes pericenter shocks weak; small velocities make apocenter shocks weak;

Orbital plane oblique to black hole spin can precess.

## Putting It All Together

#### Shiokawa, K, Cheng, Piran & Noble 2015



## Immediate Result

- ~1/3 bound mass deflected inward near  $R_T$  by t ~ 10t<sub>0</sub>
- Most bound mass in an extended, messy elliptical flow
- Unbound mass coasts outward, slowing from  $\sim$  c/4 to  $\sim$   $\sim$  c/30

Varieties of Outflows

### **Radiation-driven Winds**

(Strubbe & Quataert 2009, 2011; Metzger & Stone 2016)

If mass-return rate super-Eddington, maybe  $L_{acc} > L_{E}$ ?

Assume luminosity ~  $(R_{ISCO}/2R_T)L_{acc}$  from fallback shock at ~2R<sub>T</sub>; Guess fraction of returning mass to expel; Guess fraction of v<sub>ff</sub>(R<sub>T</sub>) for terminal speed.

Fallback shock photons diffuse out through outflow; Disk radiation (filtered by outflow?) reprocessed by unbound matter

Transfer through radiation-driven outflow + unbound matter makes optical/UV continuum + emission lines; a very extended stellar atmosphere! (Roth et al. 2016)

Problems:

So much put in by hand; Shock near R<sub>T</sub> usually weak; Asymmetry of unbound matter + optical depth of outflow lead to shifted lines

### There Are Jets!

### Swift discovered two, both in 2011

#### Dramatically variable

Very hard spectrum



SwJ1644+57: left (K & Piran 2011); right (Burrows et al. 2011)

### Maybe There Aren't Jets, After all

- VLBI —> v < 0.3C (Yang et al. 2016)
- Fe Ka lags continuum by ~ 100 s ~ 10 r\_g (Kara et al. 2016)



### Maybe There Are Jets, After All (Lu, K, Kumar & Crumley 2017)

- Close in and without relativistic motion, thermal photons from disk keep electrons too cool to produce hard spectrum
- Continuum dilution —> true Kα lag ~1000 s
- Relativistic beaming, larger lengthscale needed for low enough ionization to permit Ka emission
- Beamed X-rays accelerate disk atmosphere
- Multiple Compton scatters in cool medium create red tail; continuum dilution shortens the apparent lag



Unbound Matter (Guillochon et al. 2016; K, Piran, Svirski & Cheng 2016)

Unbound mass carries as much energy as a supernova

 $E_{\text{unbound}} \simeq 6 \times 10^{50} [(k/f)/0.08]^{-1/3} (M_*/M_{\odot})^{-1/3+\xi} M_{BH,6.5}^{1/3} \text{ erg}$ 

Spherically-expanding ejecta slow down only after

 $10(n_{\rm ext}/10^4 {\rm cm}^{-3})^{-1/3}[(k/f)/0.08]^{1/6}(M_*/M_{\odot})^{1/6-\xi/2}M_{\rm BH,6.5}^{-1/6} {
m yr}$ 

Actual unbound ejecta form a thin wedge, ~1 rad in azimuthal extent; drive a wider wedge-shaped bow shock:

 $\Delta \theta \sim 0.2 [R_s/R_s(\Delta \theta)]^{2/3} \left[1 - R_s(\Delta \theta)/R_s\right]^{1/3} (M_*/M_{\odot})^{1/9} M_{BH,6.5}^{-1/9}$ 

If external density moderately high and bow shock leads to equipartition magnetic field and relativistic electrons, detectable synchrotron emission

### Example: ASASSN 14li

Observed multiple times at several radio frequencies

Each spectrum —> peak frequency, flux at peak frequency; self-absorbed synchrotron model determined by R, n<sub>e</sub>, and B; energy minimization fixes all three.





- Outflows in TDEs can be rather different from AGN outflows
- Best evidence for jets (in some instances) and the unbound debris
- Winds due to super-Eddington luminosity much discussed and plausible, but luminosity may not reach those levels, and not observationally supported