





Sunny Itzhaki

The 5<sup>th</sup> Indian-Israeli conference on strings



That is based on something: [1808.02259](#) and [1811.12117](#) with [Karinne Attali](#) and work in progress with [Amit Giveon](#) and [David Kutasov](#)



Related also to

[Lior Liram's](#)

later today.

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and work in progress with [Amit Giveon](#) and [David Kutasov](#)

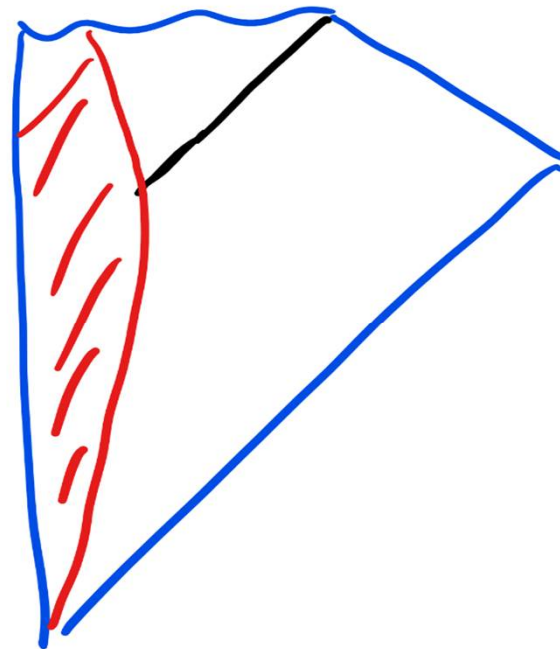
The question that we wish to ask is:

What are black holes made of?

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In GR the answer is: **nothing**

The BH is empty.



# What are black holes made of?

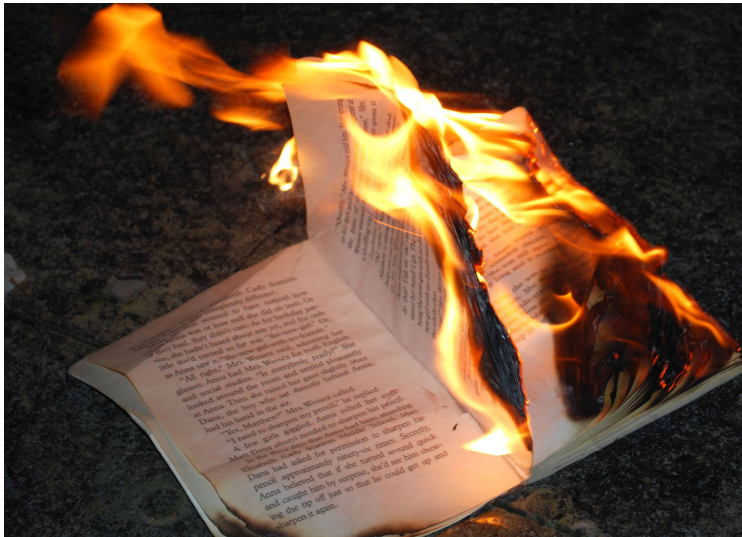
In GR the answer is: nothing. The BH is empty.

**A comment:**

This answer is very much related to the BH info puzzle.

## Related to the BH info. puzzle

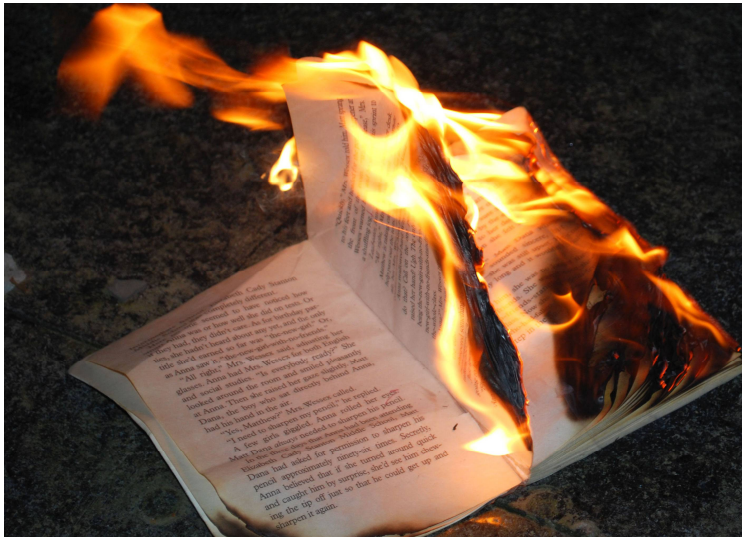
When **something** is burned info  
is lost in practice, but not in principle.



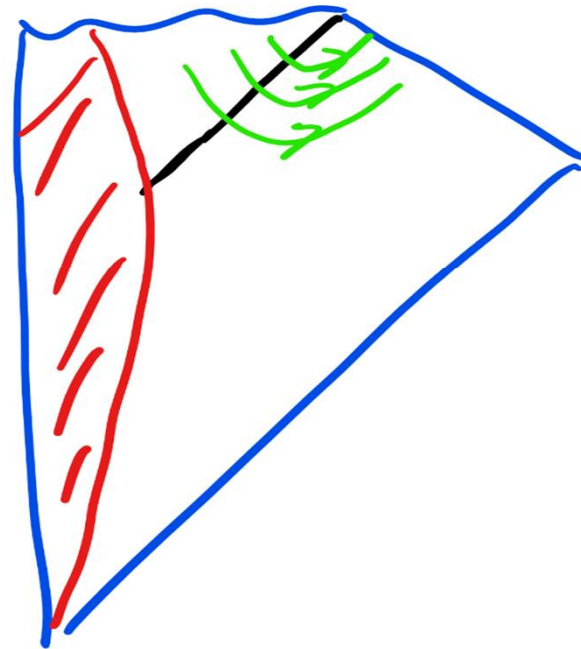


## Related to the BH info puzzle

When **something** is burned info  
is lost in practice, but not in principle.



When **nothing** is burned info  
is lost in principle.



# What are black holes made of?

In GR the answer is: nothing. The BH is empty.

There is a simple argument why we expect this

to be the answer in quantum gravity (whatever that is):

**everything** falls to the BH.

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1. Black NS5-branes are filled with **folded strings**.
2. We'll see why they don't fall to the BH.

Our starting point is an observation due to [Maldacena](#) about long strings in linear dilaton backgrounds (his motivation was a different aspect of BH physics).

Let's start by considering a 2D linear dilaton background:

$$ds^2 = -dt^2 + dr^2, \quad \Phi = -Qr,$$

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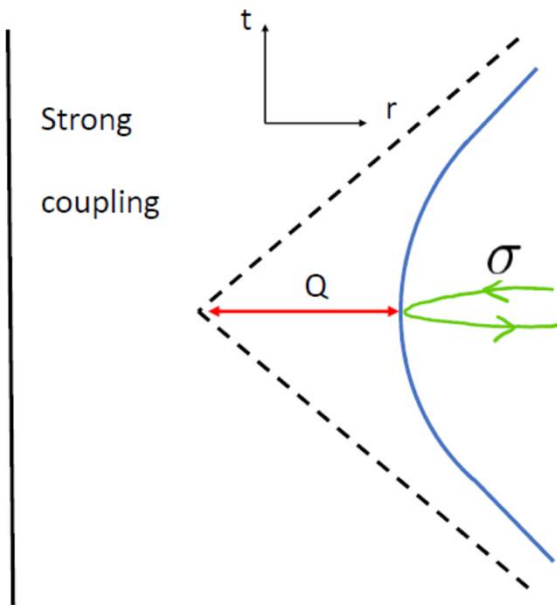
field  $\partial_+ \partial_- r = 0$    $r = r_+ + r_-$  only with a non standard energy

momentum tensor  $T_{++} = -(\partial_+ r)^2 + Q\partial_+^2 r$

Since now  $T_{++} = -(\partial_+ r)^2 + Q\partial_+^2 r$  there are new classical (non-perturbative) solutions [Maldacena, 2005] that describe long strings

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$$r = r_0 + Q \log \left( \frac{1}{2} \left( \cosh \left( \frac{t - t_0}{Q} \right) + \cosh \left( \frac{\sigma}{Q} \right) \right) \right)$$



- Can fold only towards weak coupling.

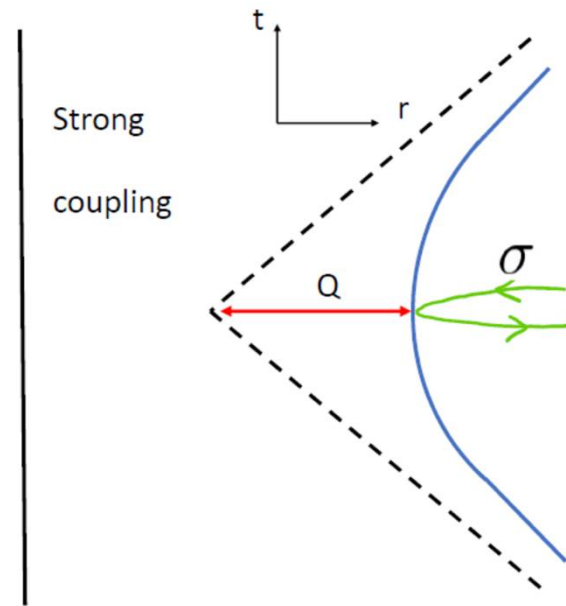


**No short strings.**

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- The length scale is  $Q$ .

(compared with the perturbative length scale  $1/Q$ )

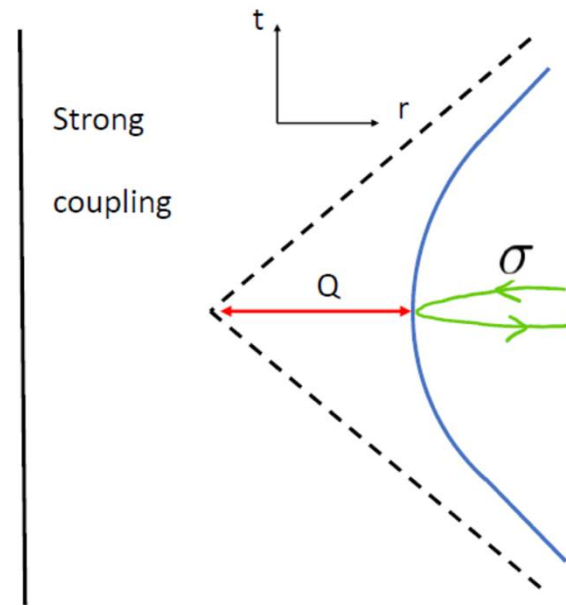


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No short strings whenever we have a local space-like linear dilaton.

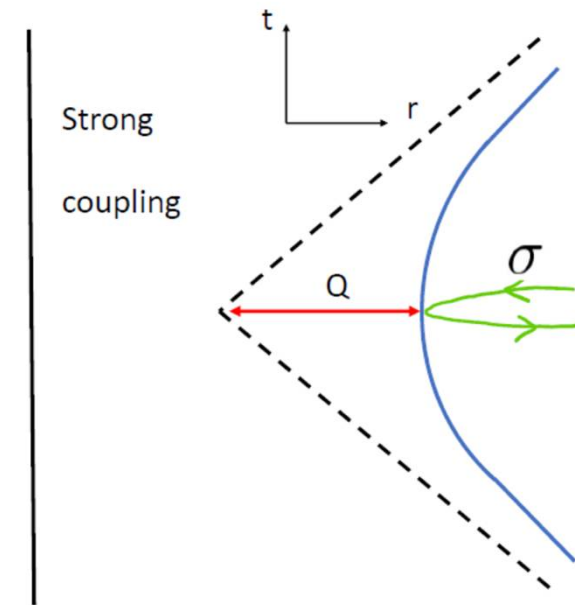


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For example: black NS5-branes

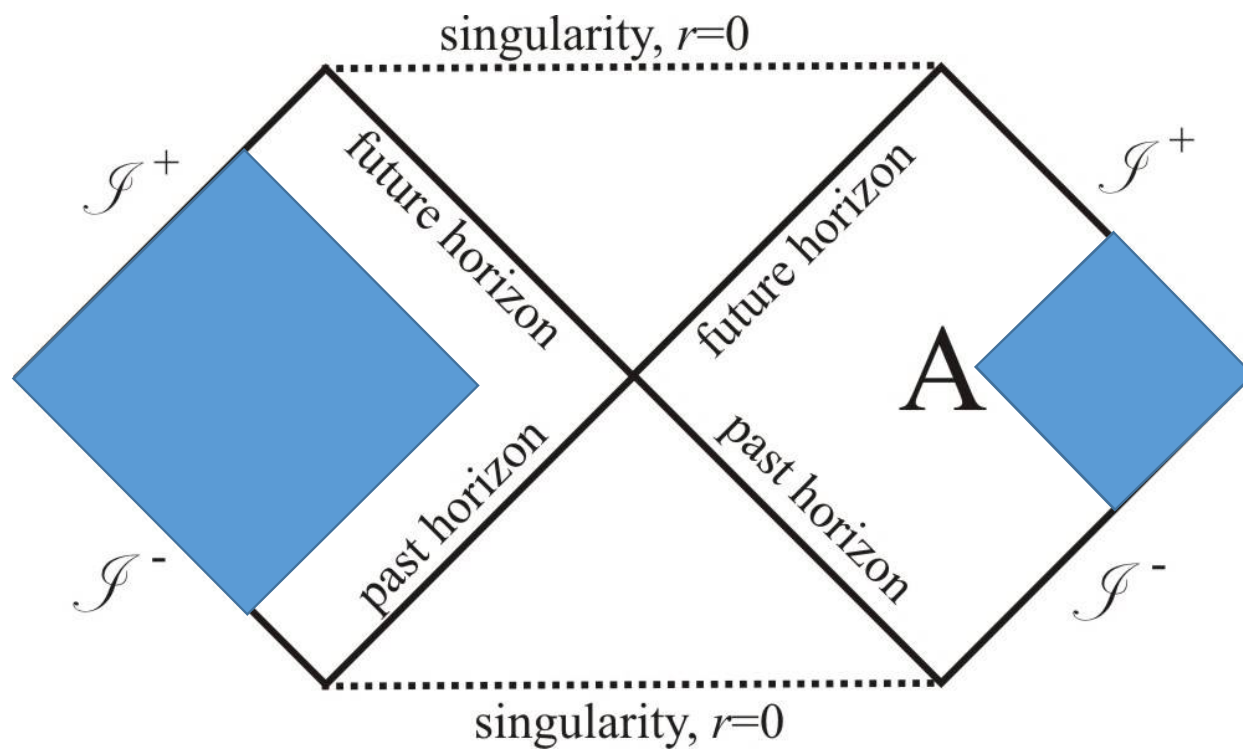
that are described by  $SL(2)/U(1)$  CFT



The folding is local.

↳ No short string outside the  $SL(2, \mathbb{R})_k/U(1)$  BH.

$$ds^2 = -f(r)dt^2 + \frac{dr^2}{f(r)}, \quad \Phi(r) = \phi_0 - Qr$$

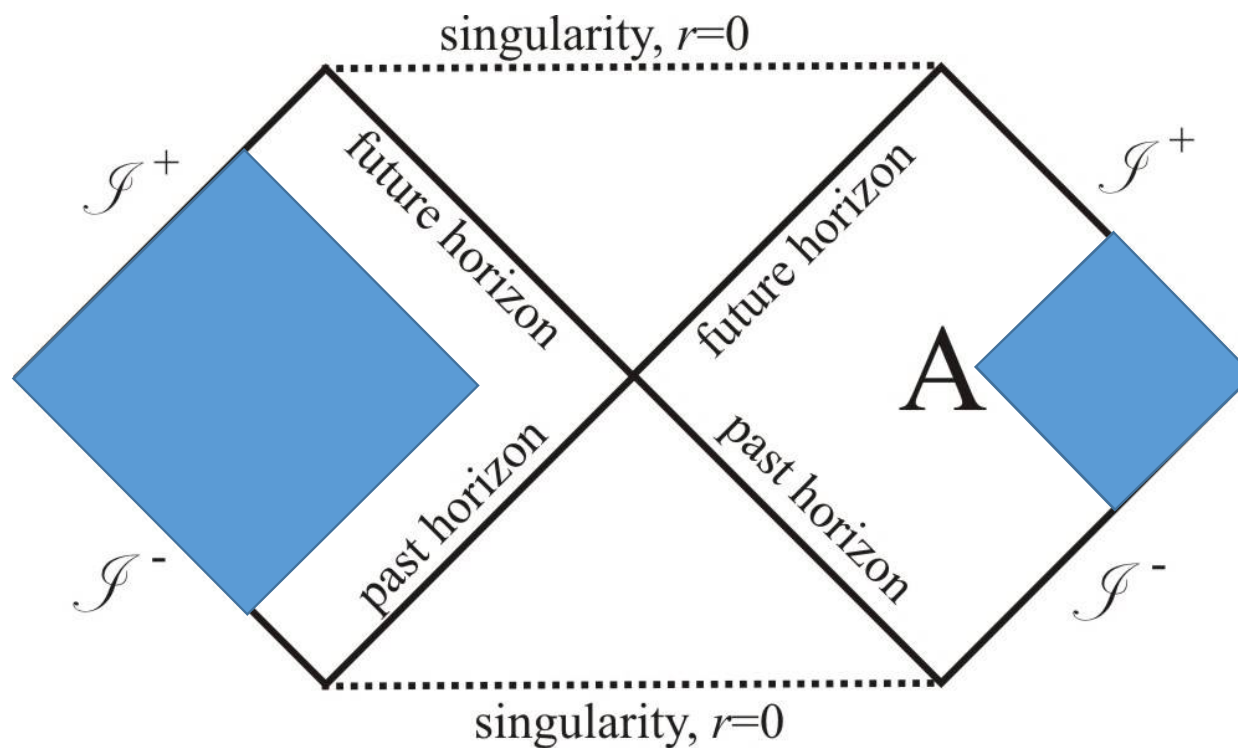


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- $k$  is the # of NS5-branes  
and  $Q = 1/\sqrt{k}$  .

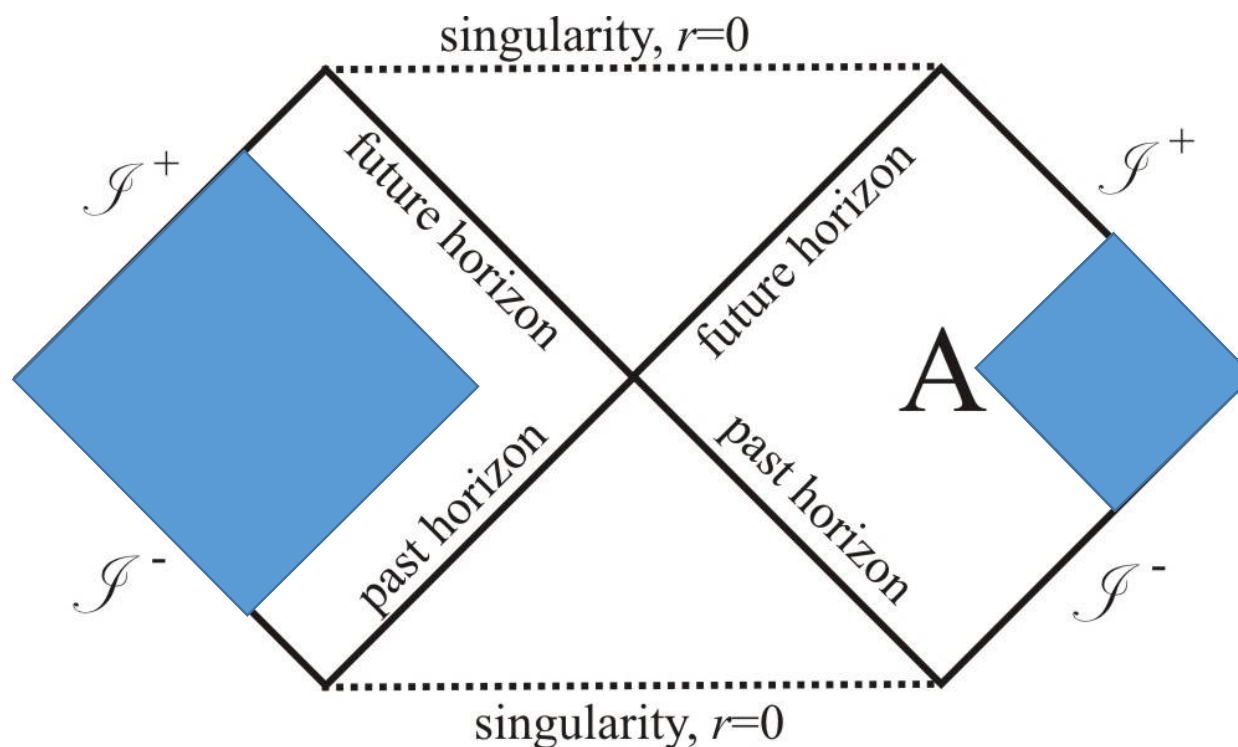


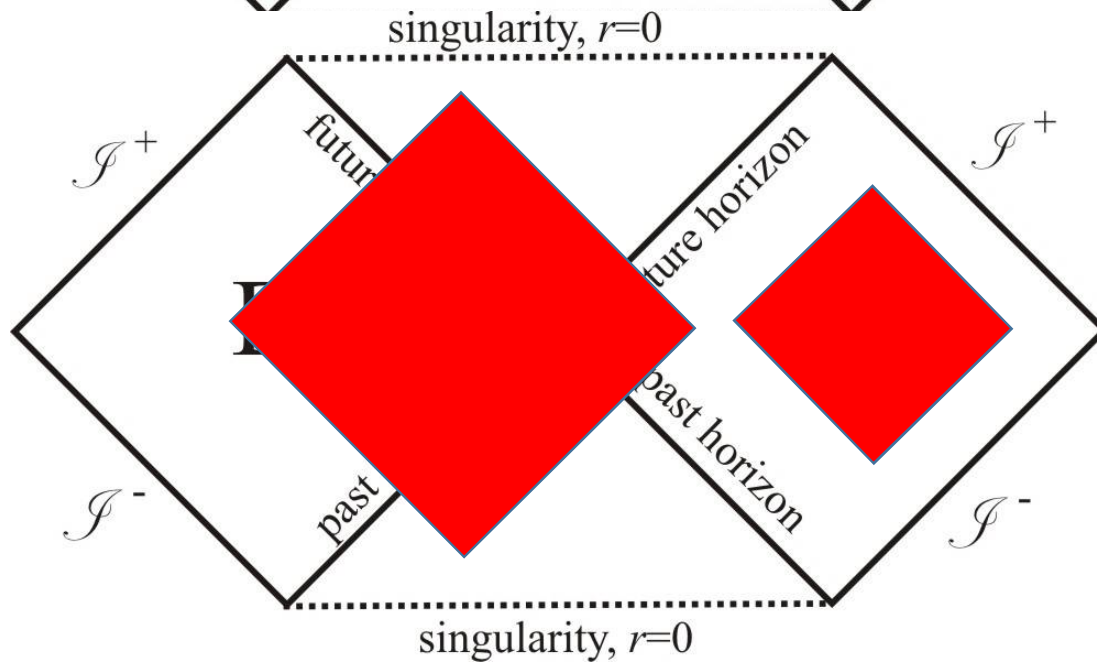
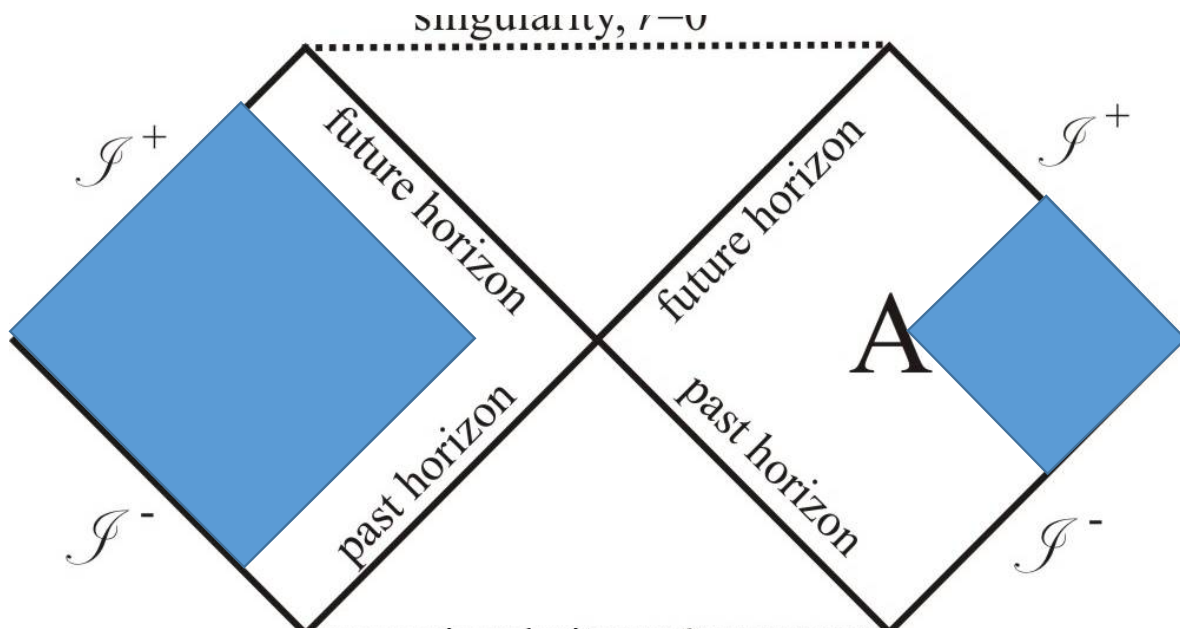
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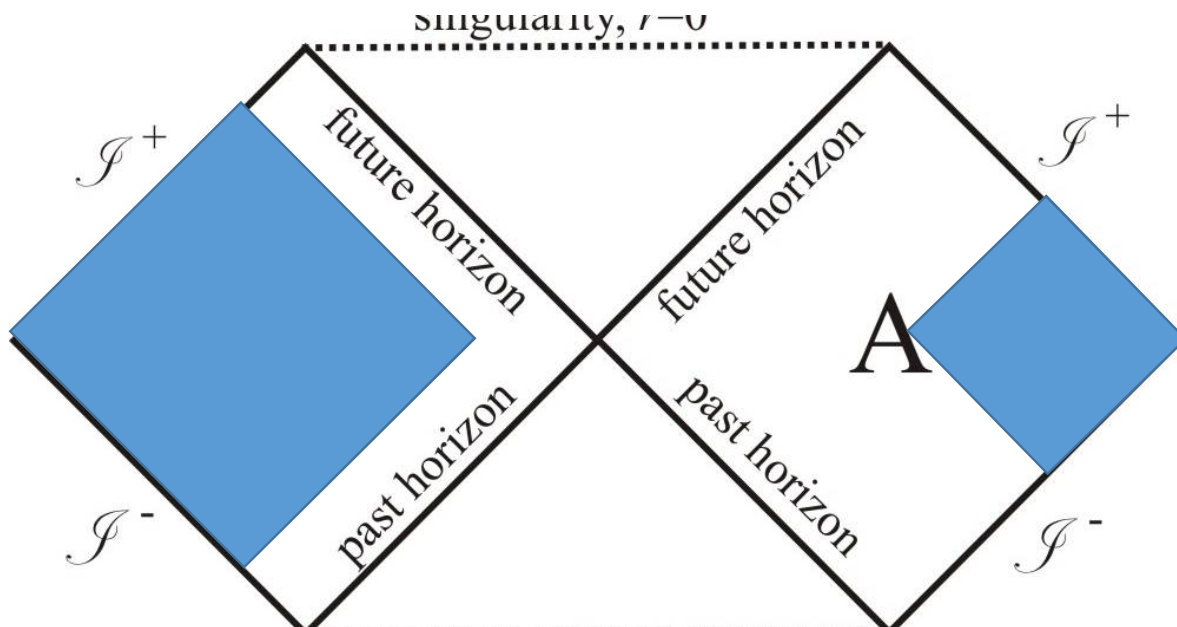
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- $k$  is the # of NS5-branes  
and  $Q = 1/\sqrt{k}$ .
- The string can fold (for large  $k$ ) only away from the BH.

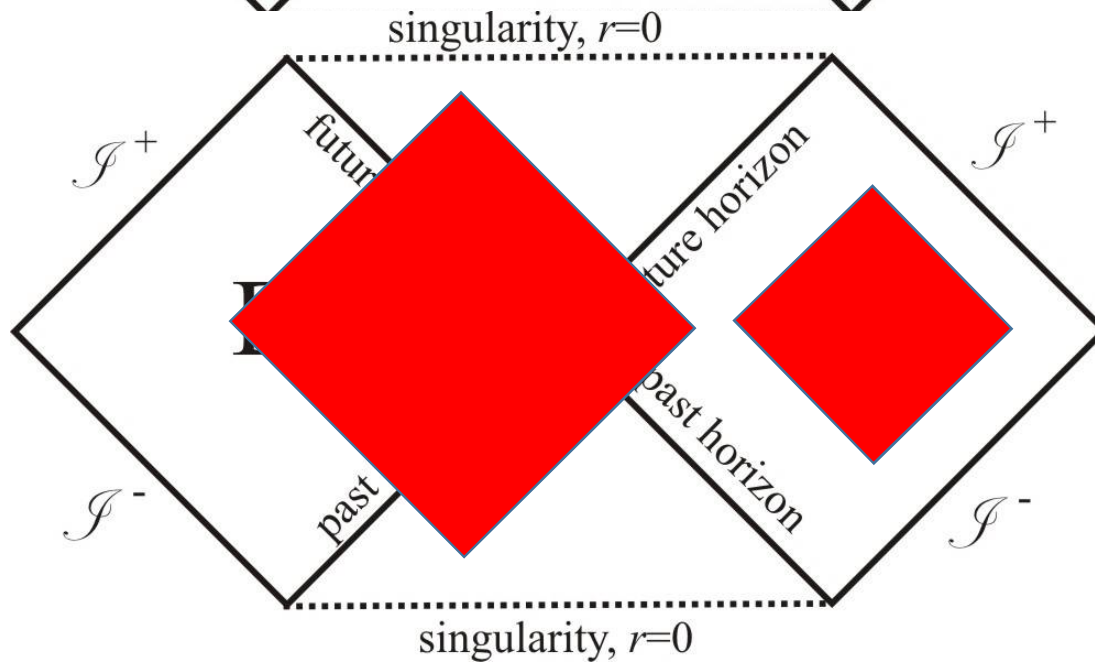




- No short string.
- No finite energy folded strings outside the  $SL(2)/U(1)$  BH.



## What about inside the BH?



- No short string.
- No finite energy folded strings outside the  $SL(2)/U(1)$  BH.

What happens inside the BH?

We have a time like linear dilaton  $ds^2 = -(dX^0)^2 + (dX^1)^2$ ,  $\Phi = QX^0$ .

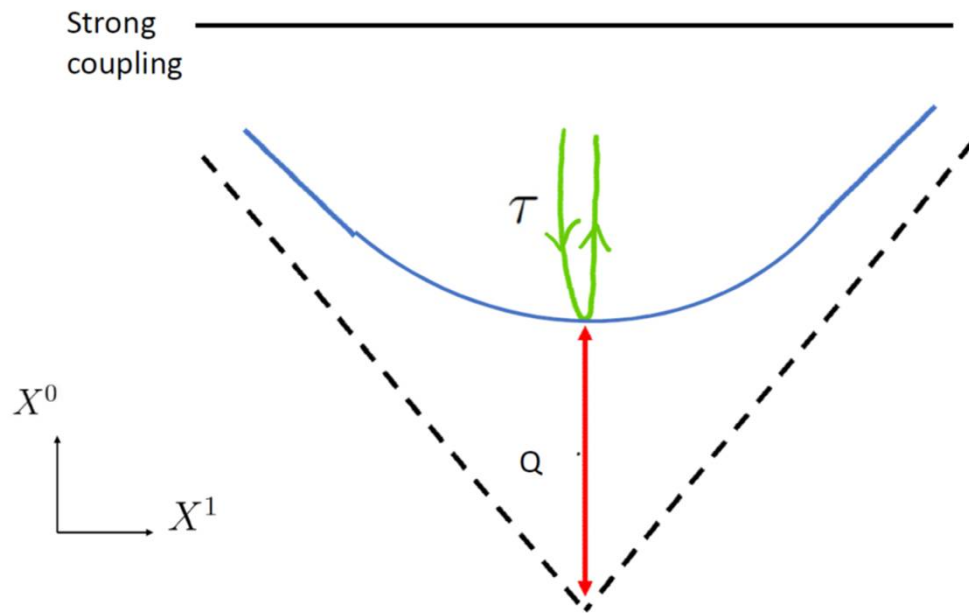
Now the classical solutions is

$$X^0 = x^0 + Q \log \left( \frac{1}{2} \left( \cosh \left( \frac{X^1 - x^1}{Q} \right) + \cosh \left( \frac{\tau}{Q} \right) \right) \right).$$

## What happens beyond the horizon?

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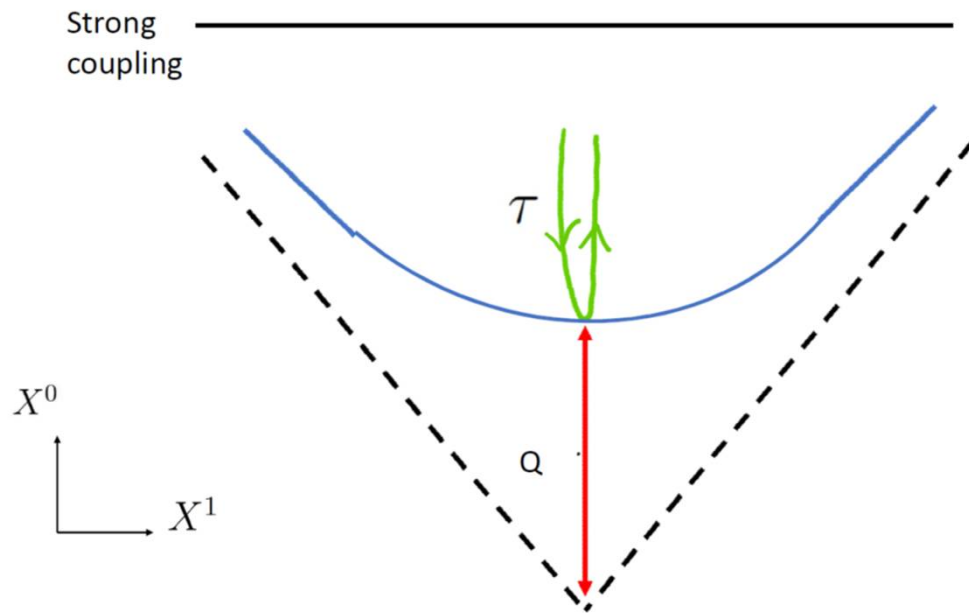


A closed (folded) string is created at a certain point in space-time.

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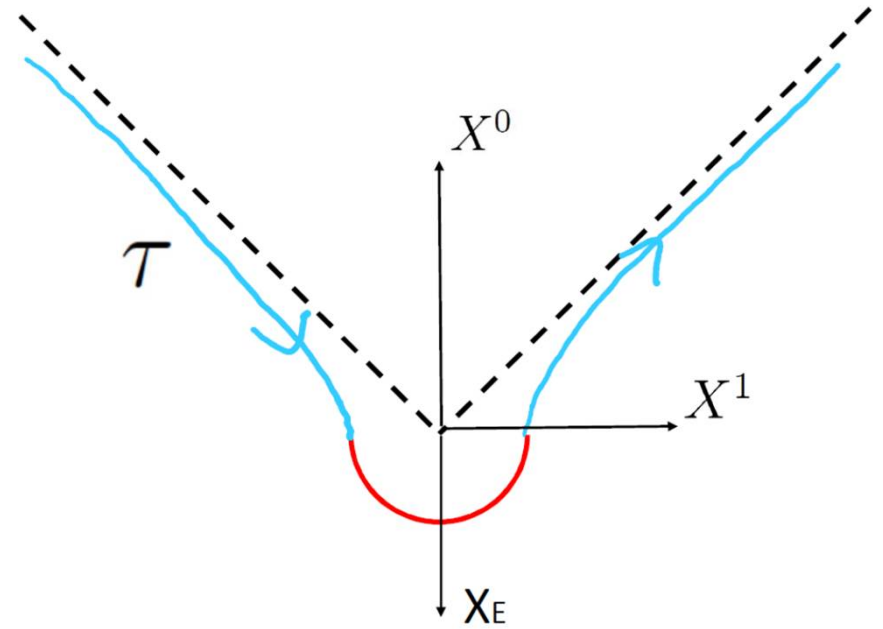
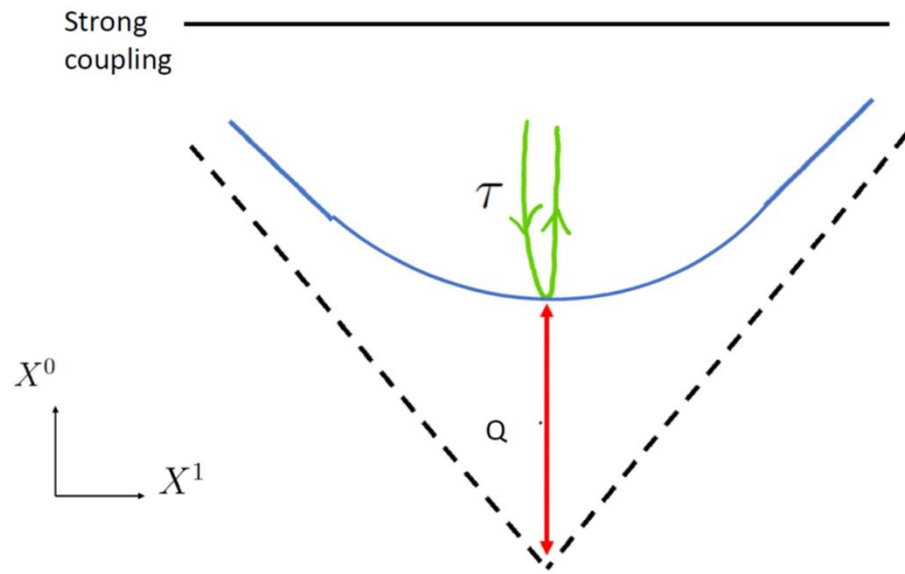


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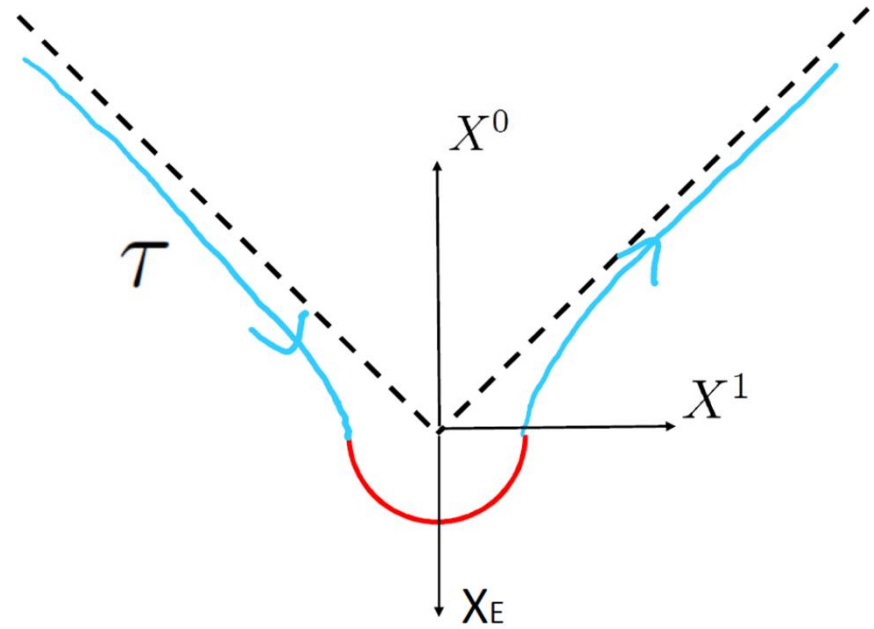
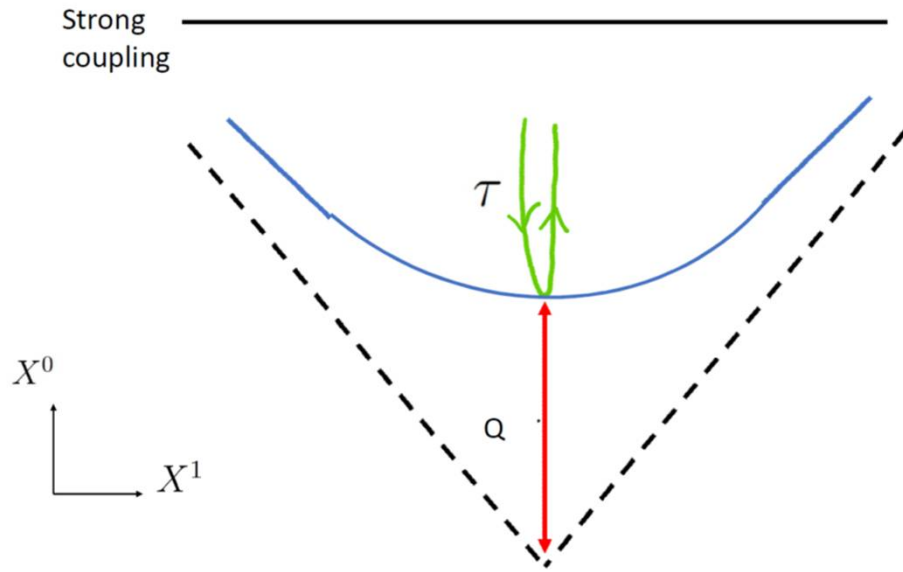
Very different than standard spontaneous creation.



## Comparison to the Schwinger mechanism

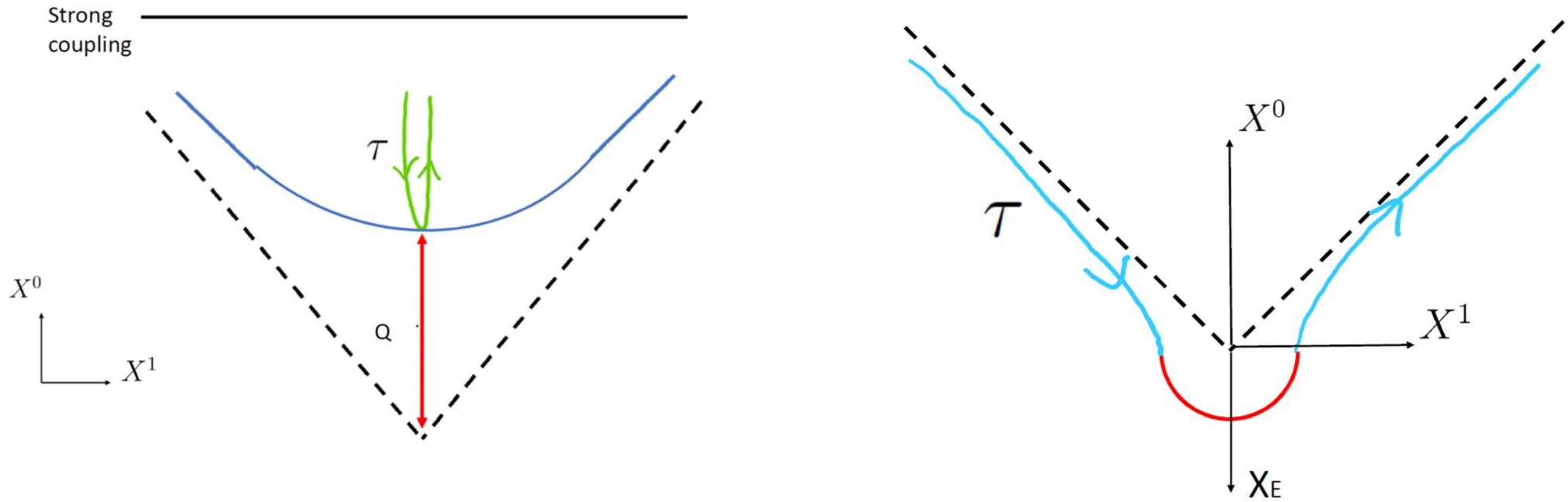


## Comparison to the Schwinger mechanism



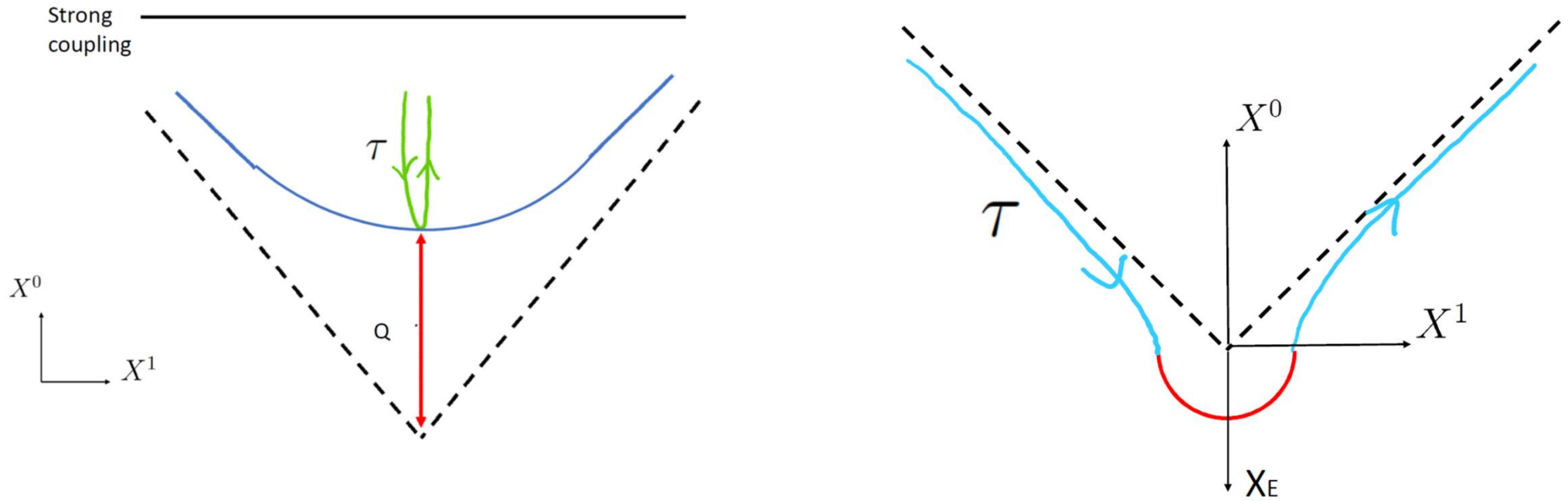
- Classical solution in Min. space

## Comparison to the Schwinger mechanism



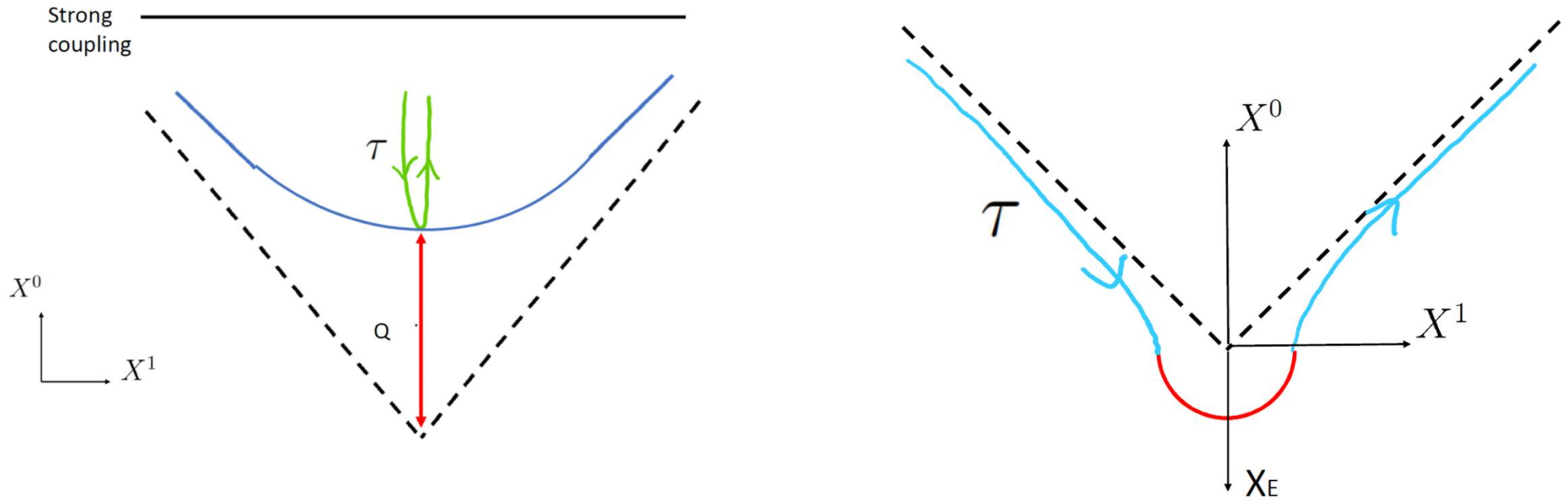
- Classical solution in Min. space ➡ no tunneling ➡ no exp. suppression.

## Comparison to the Schwinger mechanism

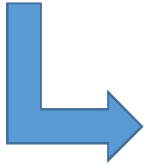


- Classical solution in Min. space → no tunneling → no exp. suppression.
- No suppression in the string coupling constant either.

## Comparison to the Schwinger mechanism



- Classical solution in Min. space  $\Rightarrow$  no tunneling  $\Rightarrow$  no exp. suppression.
- No suppression in the string coupling constant either.
- The point where the string folds is traveling faster than light.



So we have a lot of these folded strings, but what can we say about them?

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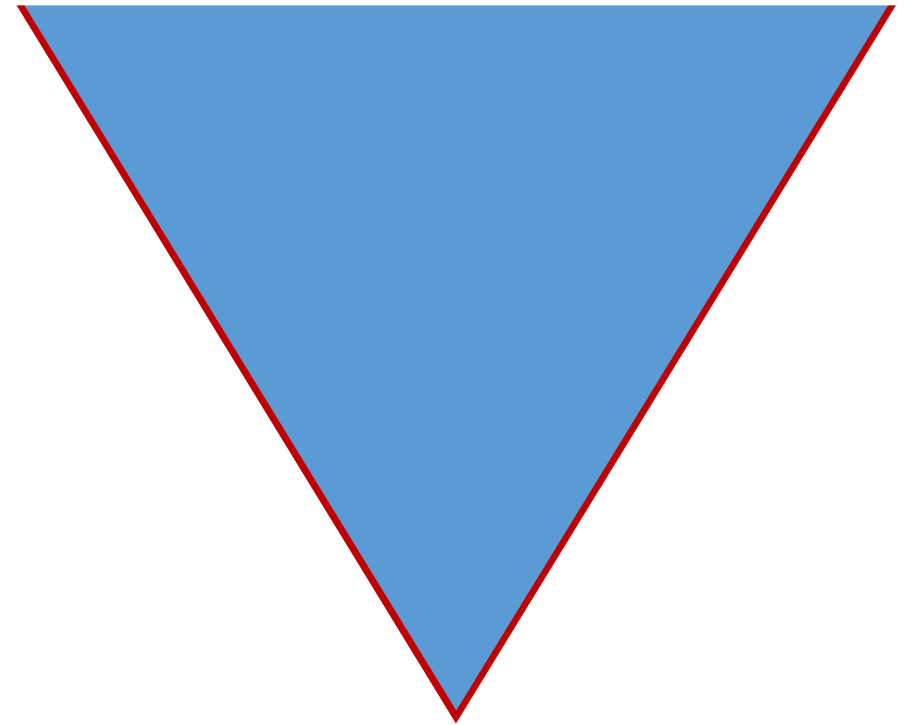
Since they are created at a certain point in space-time their total energy (and momentum) must vanish.

So we have a lot of these folded strings, but what can we say about them?

Since they are created at a certain point in space-time their total energy (and momentum) must vanish. **Indeed**

$$T^{00} = \frac{1}{2\pi\alpha'} \frac{-4 \cosh\left(\frac{x-x_0}{Q}\right) + 4e^{\frac{t-t_0}{Q}}}{\sqrt{-1 + \left(2e^{\frac{t-t_0}{Q}} - \cosh\left(\frac{x-x_0}{Q}\right)\right)^2}}$$

and **E=0**.



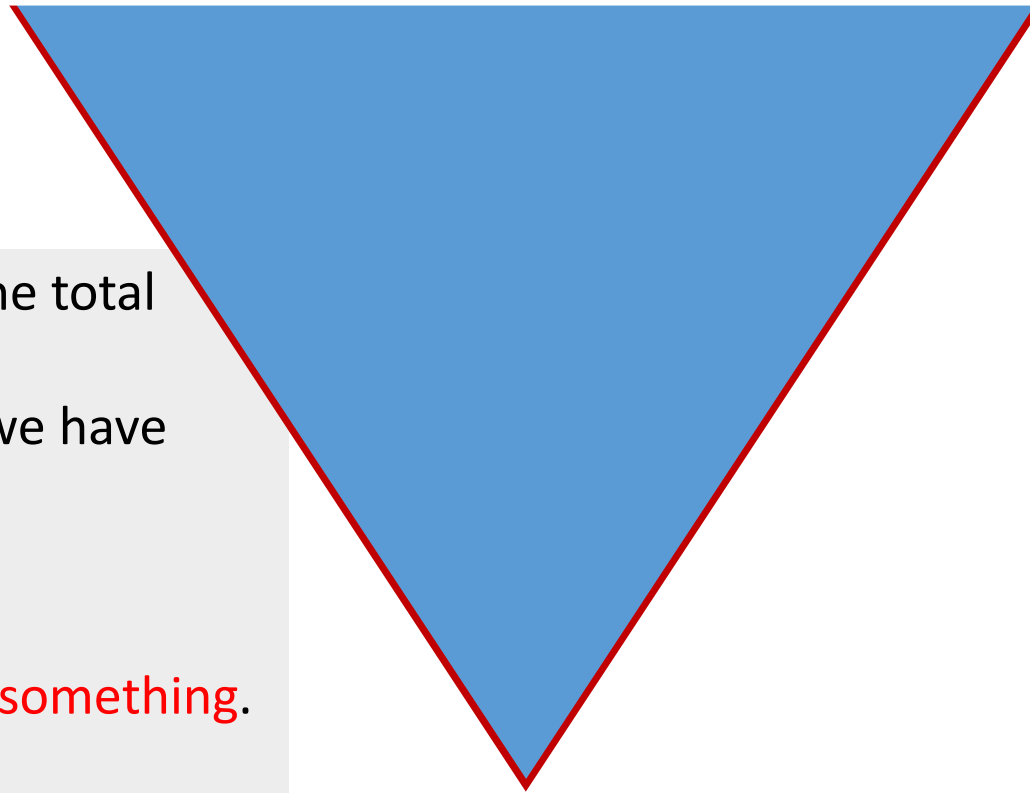


## Total energy vanishes

Normally when the total  
energy vanishes we have

**nothing.**

But this is clearly **something.**



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In short: they violate the ANEC.

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They do so in the right way (as if they know they have to solve the BH info puzzle).

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It is the statement that

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Useful in QFT (where it can be proved),

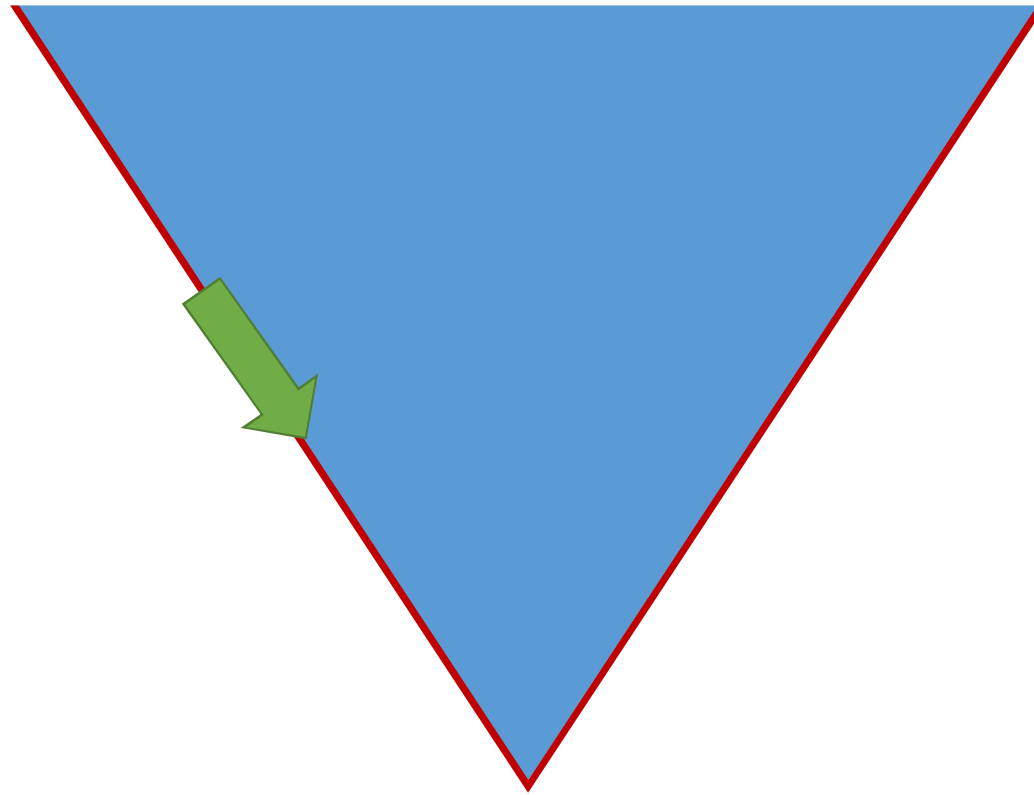
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Useful in QFT (where it can be proved) and is sufficient to prove various fundamental properties of BH and Cosmology. [For example that the BH is empty.](#)

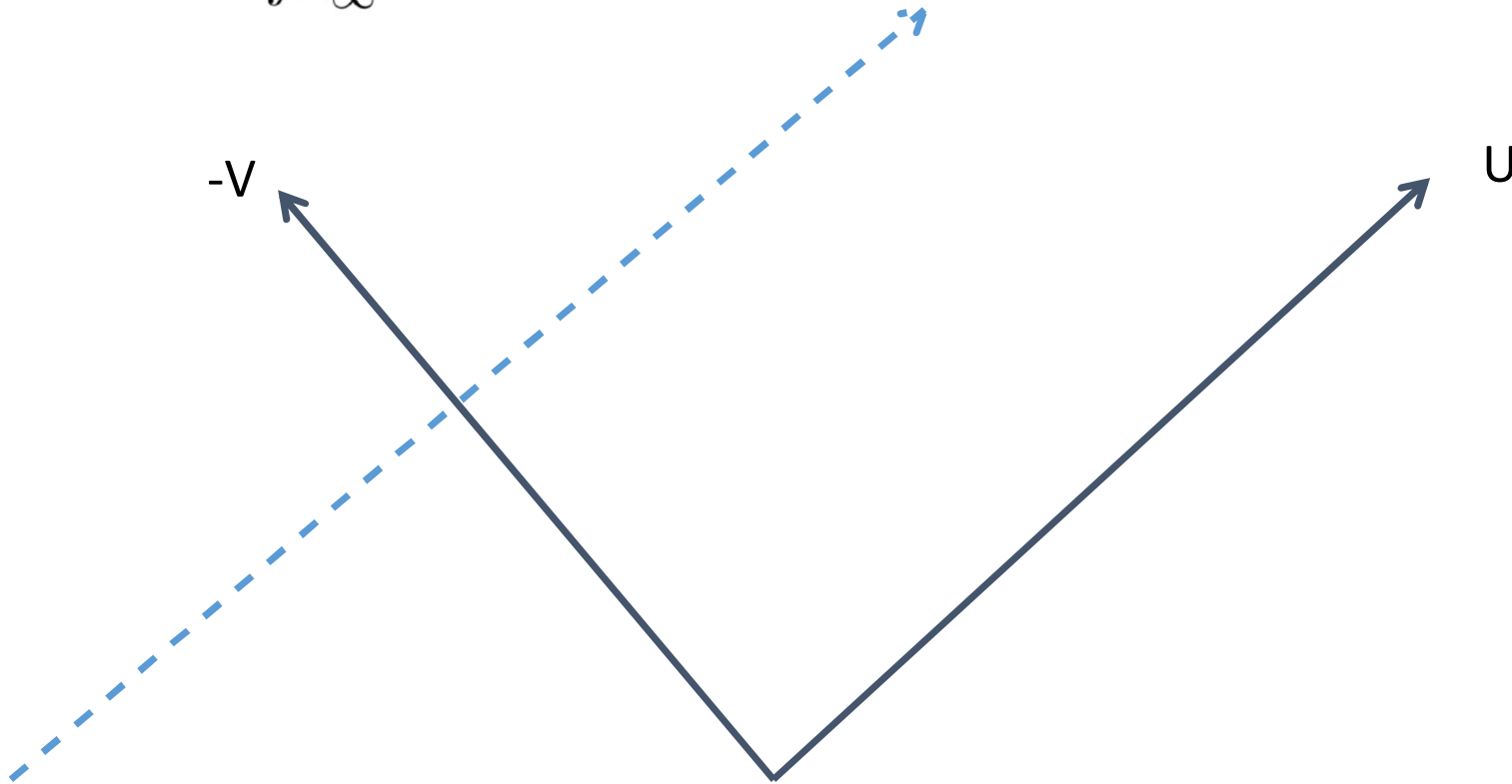
$$T_{uu} = \frac{1}{2\pi\alpha'} \frac{-e^{-\frac{x-x_0}{Q}}}{\sqrt{-1 + \left(2e^{\frac{t-t_0}{Q}} - \cosh\left(\frac{x-x_0}{Q}\right)\right)^2}} < 0$$



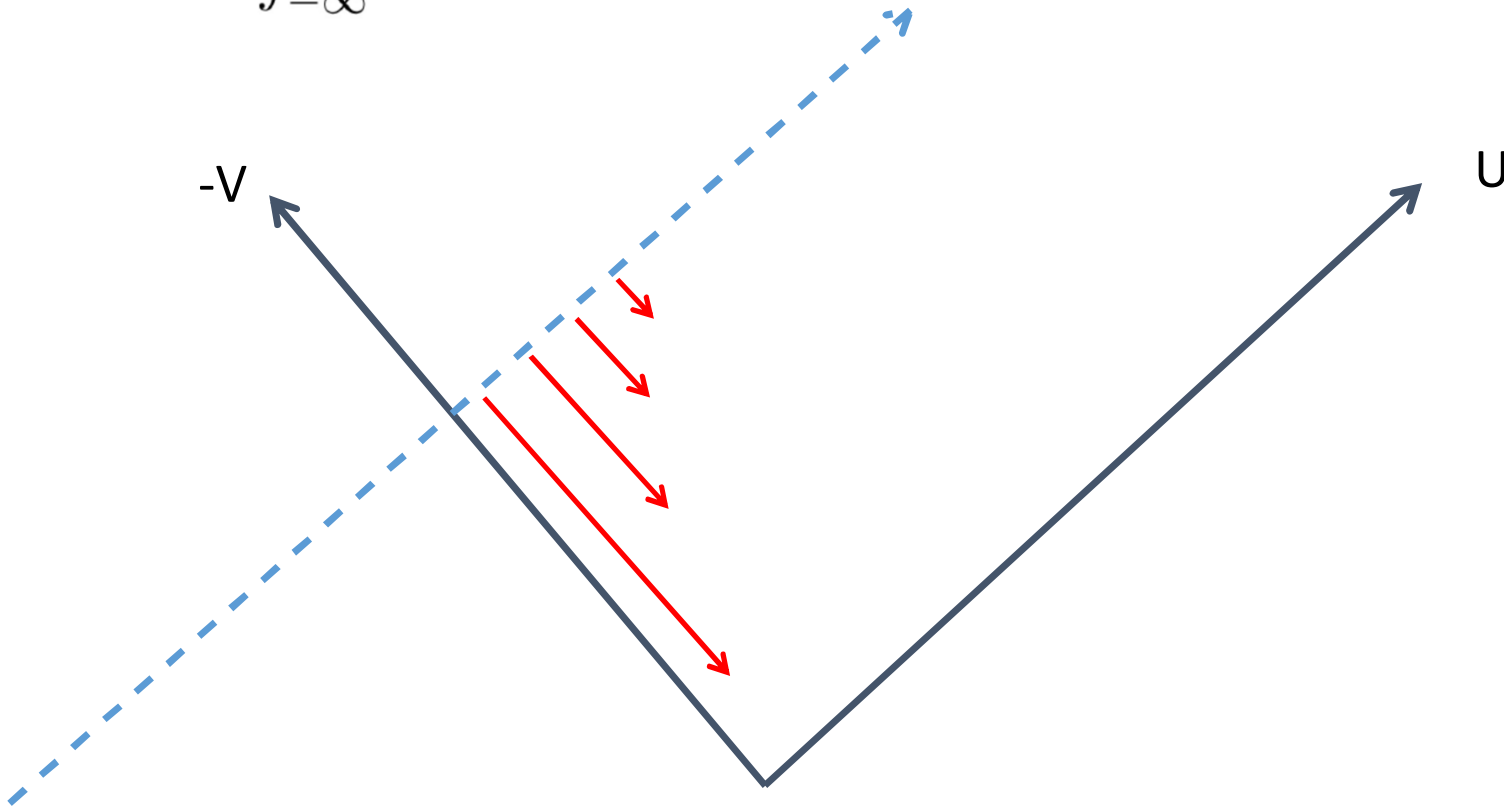


Same with ANEC

$$\int_{-\infty}^{\infty} du T_{uu} = (v - Q \log(4)) < 0$$

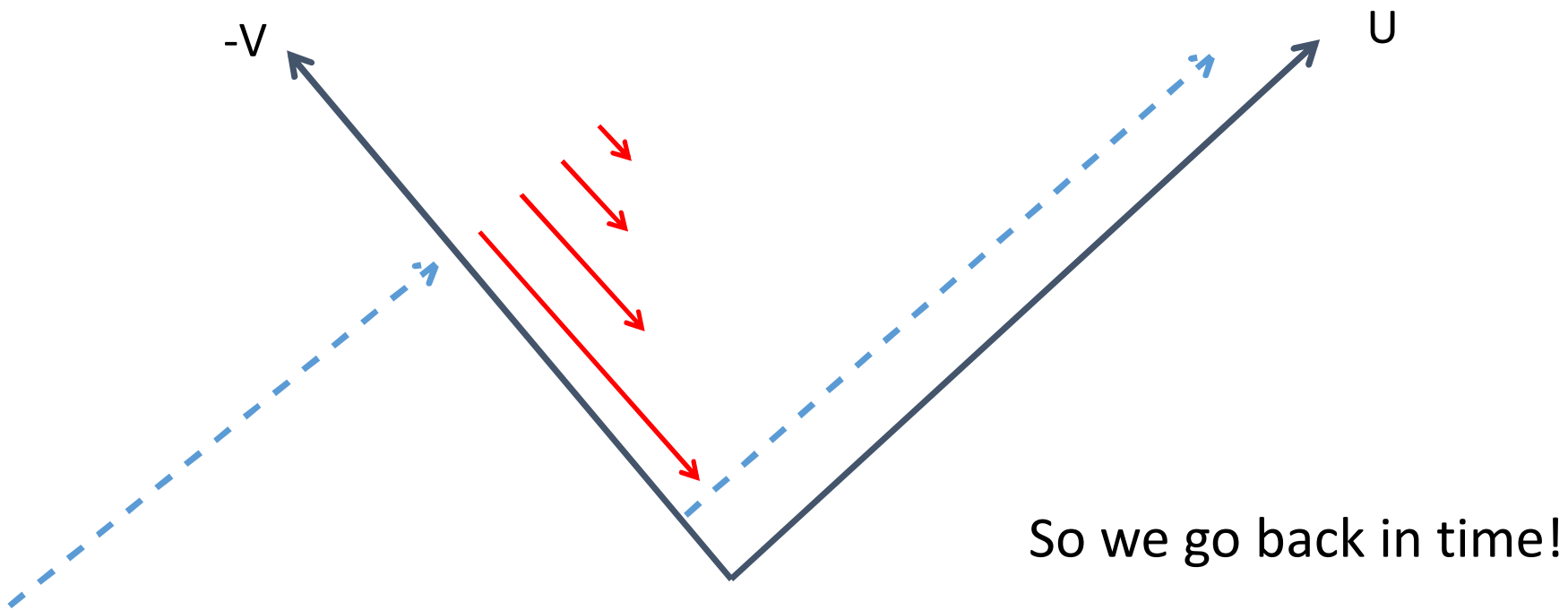


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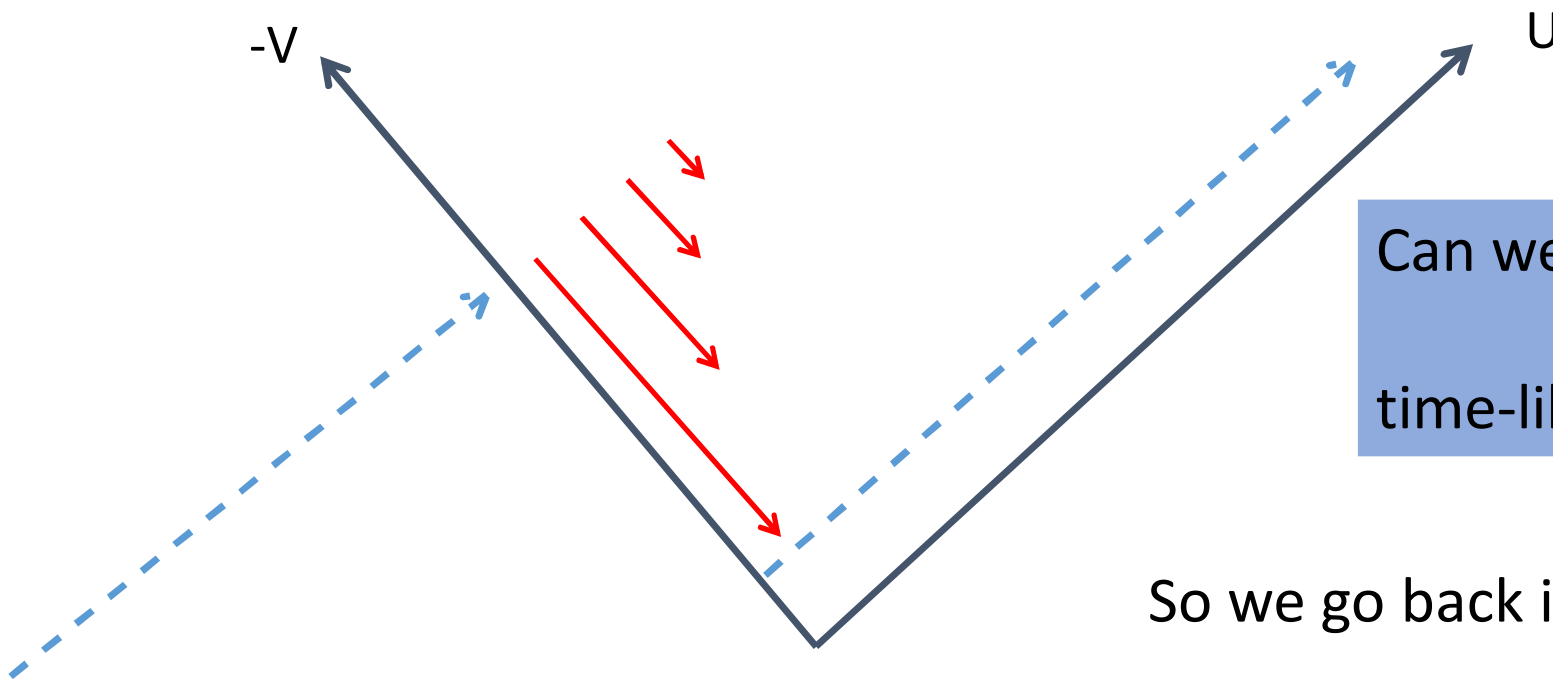


What about causality?

$$\int_{-\infty}^{\infty} du T_{uu} = (v - Q \log(4)) < 0$$



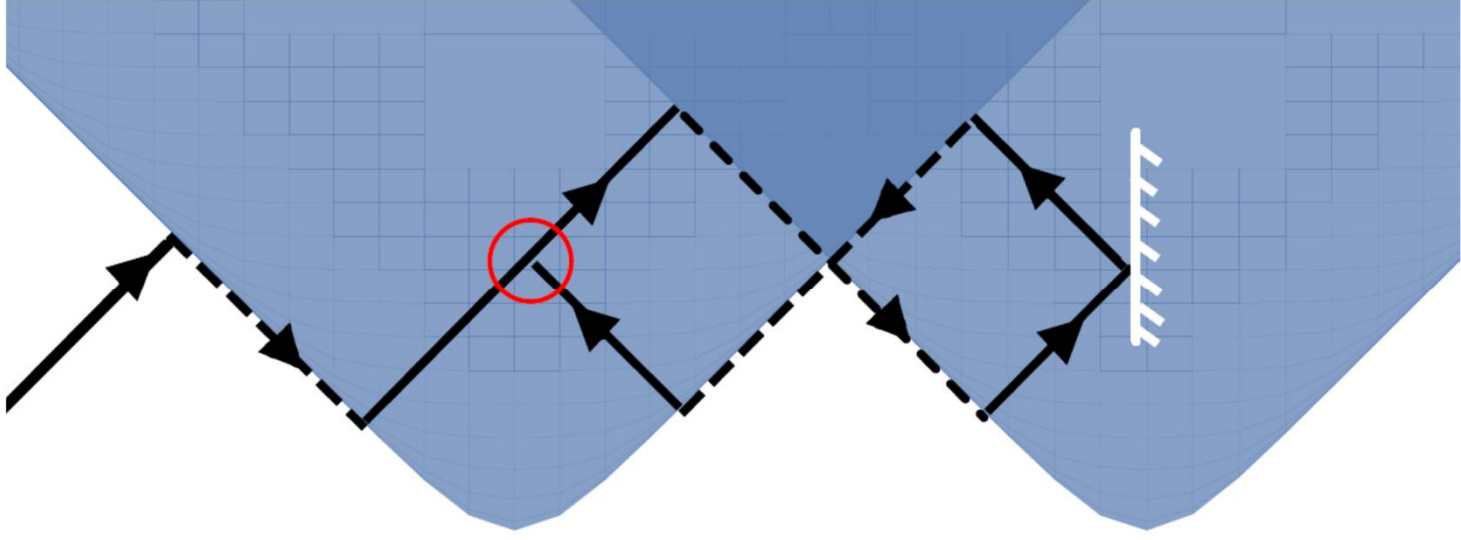
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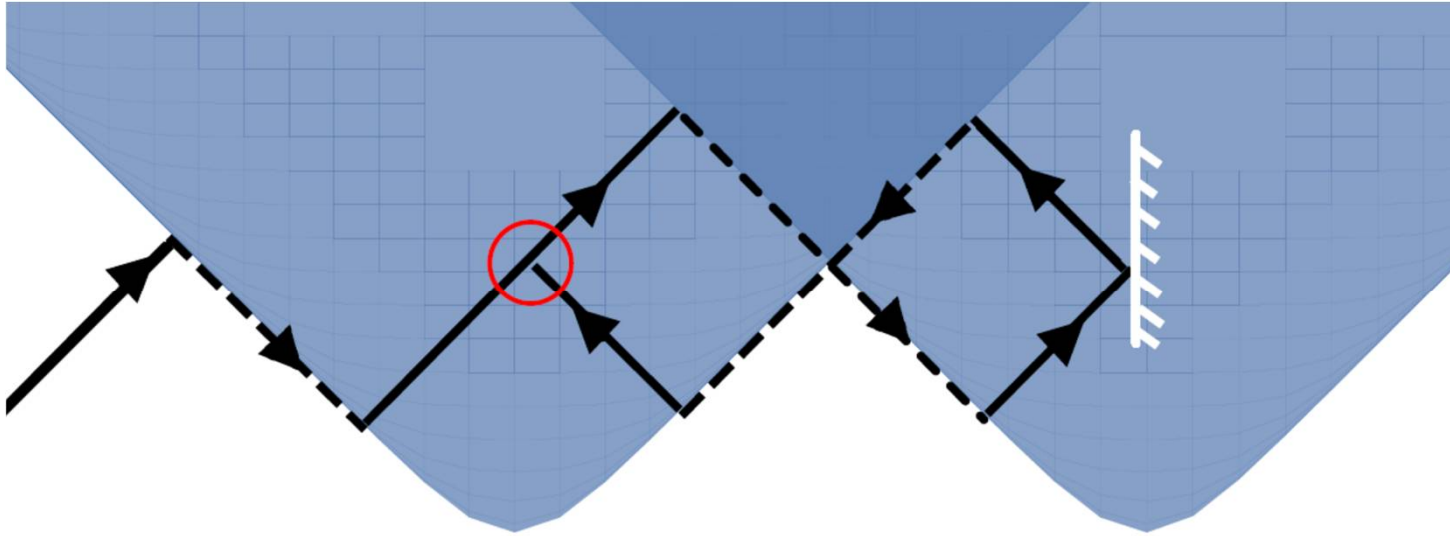
Can we have a closed  
time-like curve?

So we go back in time!

Yes we can

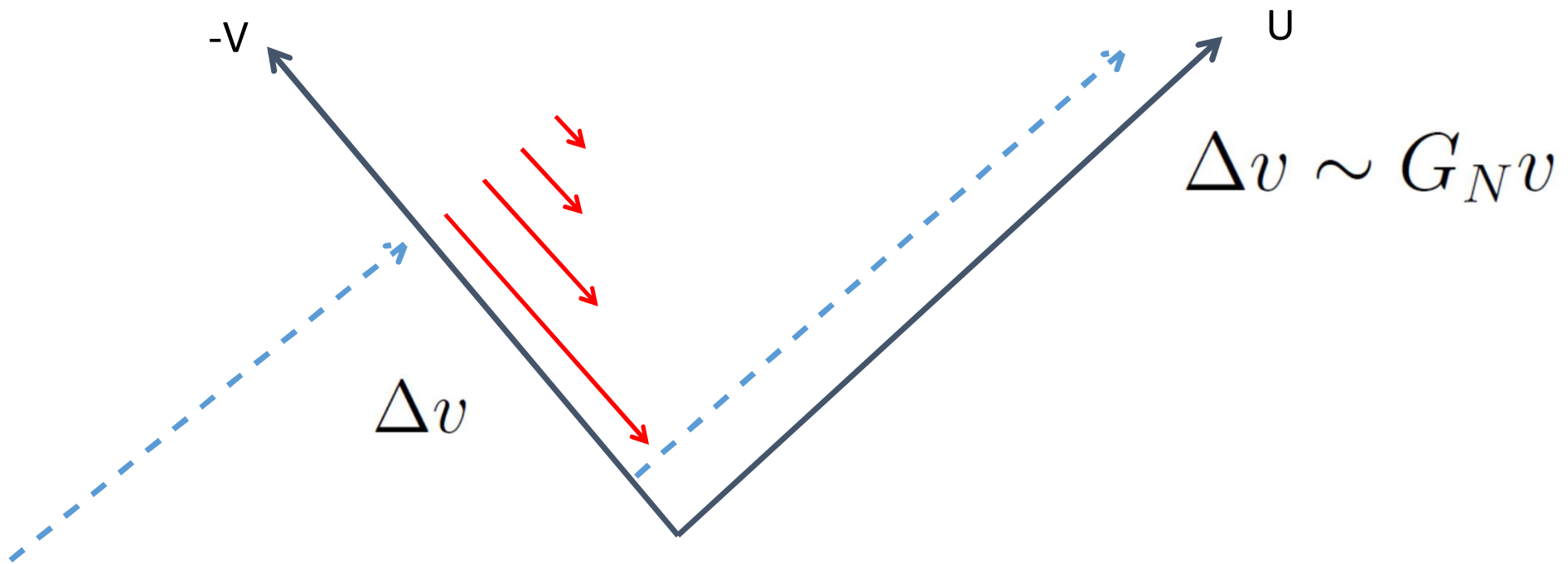


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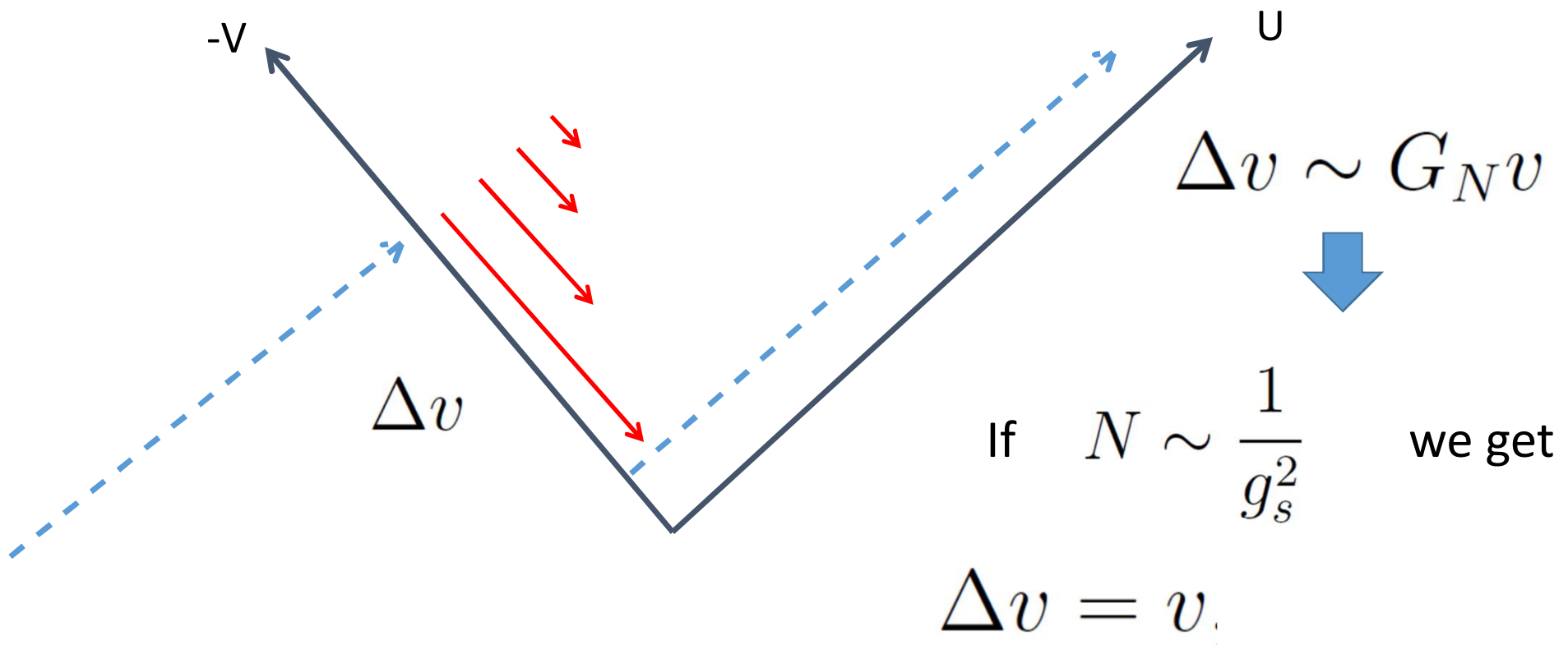


- Clear tension with causality.
- In fact also with unitarity.

Just like, **you can't be half pregnant**, if you violate the ANEC you have to fully violate it:

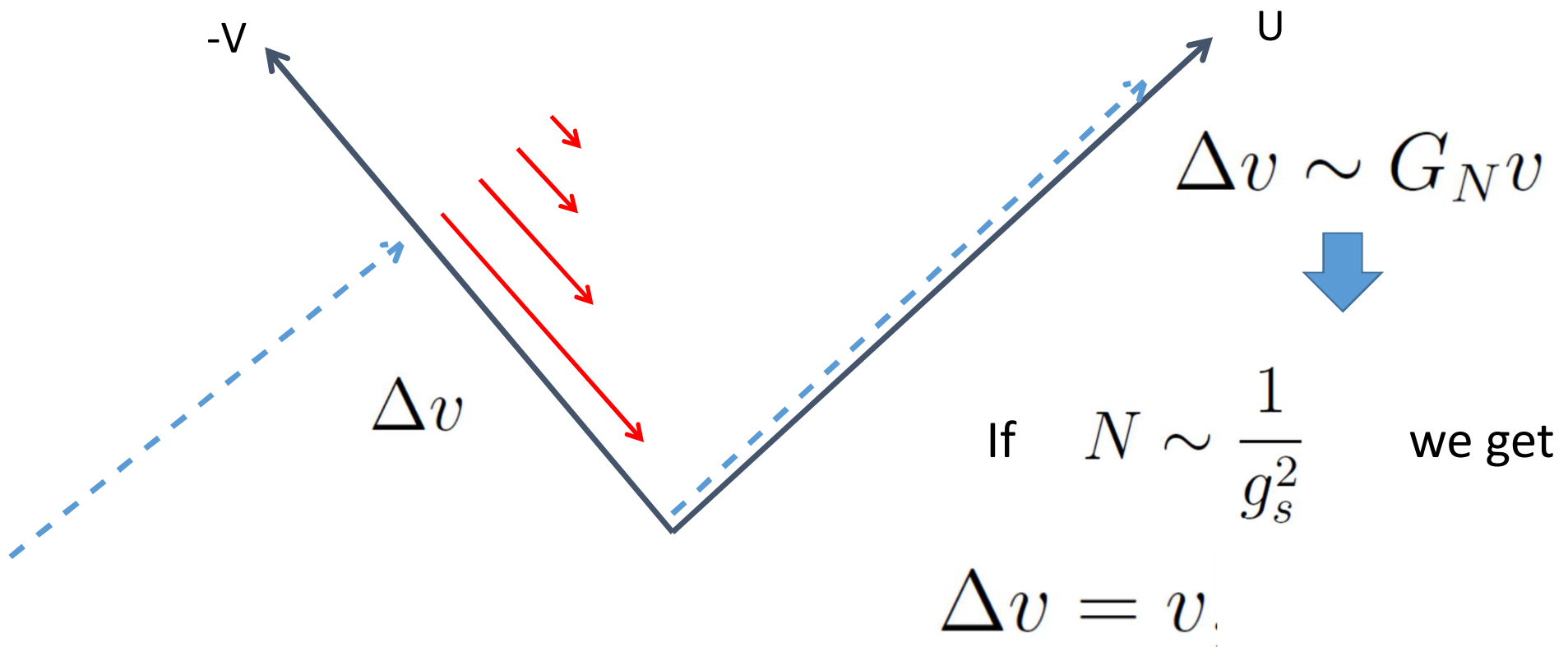


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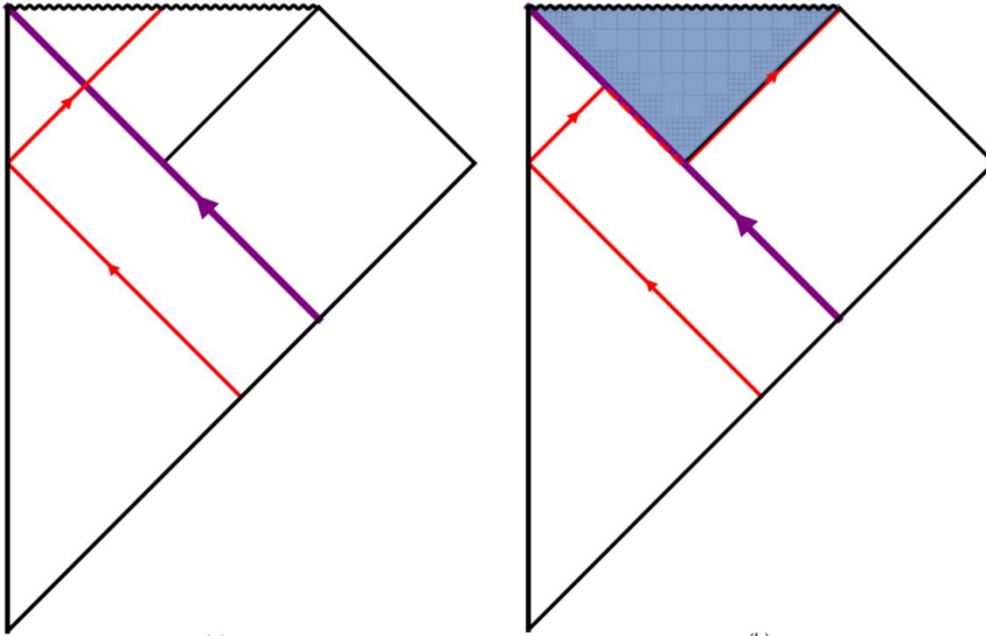
A conjecture:  $N \sim \frac{1}{g_s^2}$



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The BH is cloaked by the folded strings



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Lior will discuss the reflection coefficient that an observer at infinity can measure and conclude that the horizon is singular **from the inside**.

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But can we see more directly the folded string at infinity?

Any imprint of this story at infinity?

Classically (in  $\alpha'$ ) the answer is no (the string can fold towards the BH only behind the horizon).

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Same when perturbative corrections (in  $\alpha'$ ) are taken into account.



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This is due to a non-trivial operator (Gerasimov, Marshakov, Morozov, Olshanetsky and Shatashvili, 1990) that is invariant under the  $SL(2,R)$  symmetry.

So again something that looks like nothing.

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**Schwarzschild** is expected to be filled with all of the above.

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And what about Cosmology?

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Objects that violate the ANEC can completely change the rules of the game  
in Cosmology.

Thank you