

# *Symmetries of Feynman Integrals*

Barak Kol  
Hebrew Un., Jerusalem  
Indo-Israeli meeting  
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- Introduction
- Main Results
- Mechanism
- Applications
- More results
- Conclusions

based on work w. Philipp Burda, Subhajit Mazumdar,  
Amit Schiller and Ruth Shir in 8 papers since 2015.

# Introduction

# Motivation

At risk of stating the obvious:

Feynman diagrams are at the computational core of Quantum Field Theory.

Necessary for experimental design and data analysis including at the LHC.

# Background

- Feynman diagrams 1948.
- Schwinger parameters.
- Landau singularities 1959.
- Cutkosky discontinuities 1960.
- Dimensional regularization, Giambiagi & Bollini, 't Hooft & Veltman 1972.

# Background - well known, but less

- Integration By Parts (IBP), Chetyrkin & Tkachov 1981
- Differential Equations (DE), Kotikov 1991, Remiddi 1997.
- Dimensional Recurrence, Tarasov 1996.
- Canonical Basis for DE, Henn 2013.

# Modern perturbative field theory

- On shell unitarity, Bern, Dixon & Kosower 1994.
- Scattering amplitudes: Witten 2003; Britto, Cachazo, Feng & Witten (BCFW) on-shell recursion relations 2005; Cachazo, He & Yuan amplitudes 2013; Amplituhedron, Arkani-Hamed & Trnka 2014.

A general theory  
is still lacking

# Personal trajectory

- Got interested through binary problem in Einstein's gravity; Effective Field Theory for Post-Newtonian approx; Higher loop; Integration By Parts.
- Is it possible to contribute to a longstanding issue?



# Main results

# Definitions

A Feynman integral

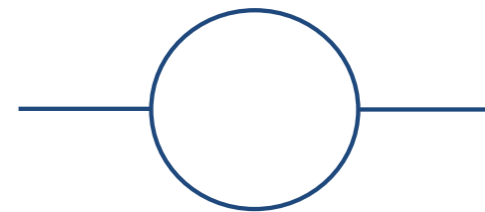
$$I(\mu_1, \dots, \mu_P, p_1^\mu, \dots, p_X^\mu) = \int \frac{d^d l_1 \dots d^d l_L}{\prod_{i=1}^P (k_i^2 - \mu_i + i0)}$$

$\mu := m^2$

Parameter space: masses  
and kinematical invariants

$$X = \{\mu_i, p_r \cdot p_s\}$$

Example:  
bubble diagram



$$I(\mu_1, \mu_2, p^2) = \int \frac{d^d l}{(l^2 - \mu_1)((l + p)^2 - \mu_2)}$$

# Definitions

A Feynman integral

$$I(\mu_1, \dots, \mu_P, p_1^\mu, \dots, p_X^\mu) = \int \frac{d^d l_1 \dots d^d l_L}{\prod_{i=1}^P (k_i^2 - \mu_i + i0)}$$

$\mu := m^2$

Comment: Numerators, spacetime dimension

# Main Result: reduction

Reduction 
$$\hat{I}(x) = \hat{I}(x_0) + \int_{x_0}^x J^\alpha(\xi) d\xi_\alpha$$

Leading singularity normalization  $\hat{I} := I/I_0$

$I_0$  - a homogeneous soln., incl. maximal cut.

$x_0$  - conveniently chosen base point within G orbit (more below).

Integrand depends on simpler diagrams, namely, with an edge contracted.

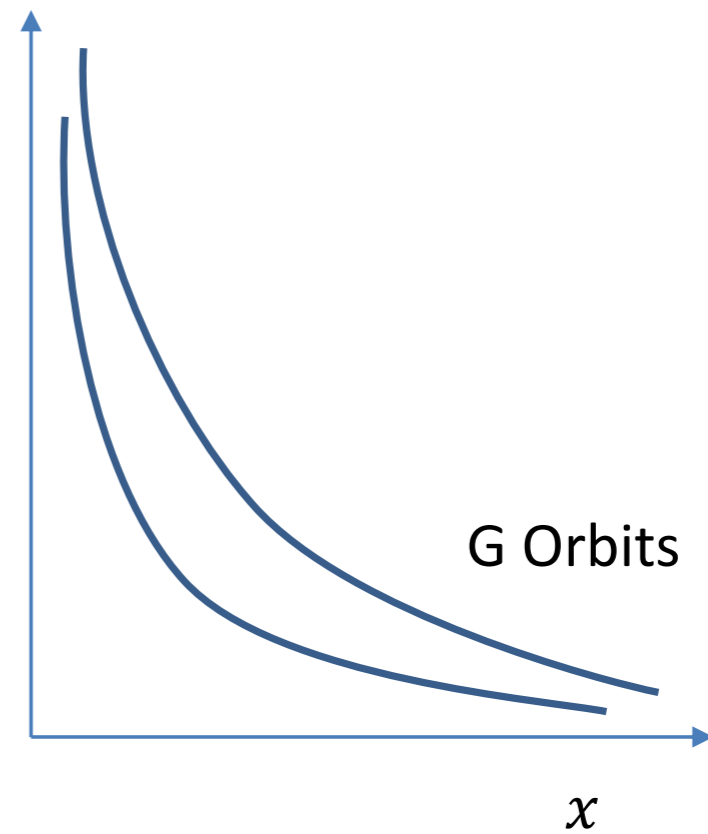
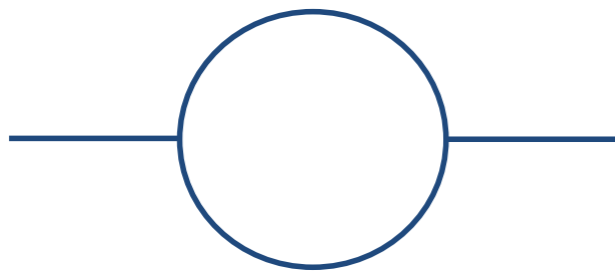
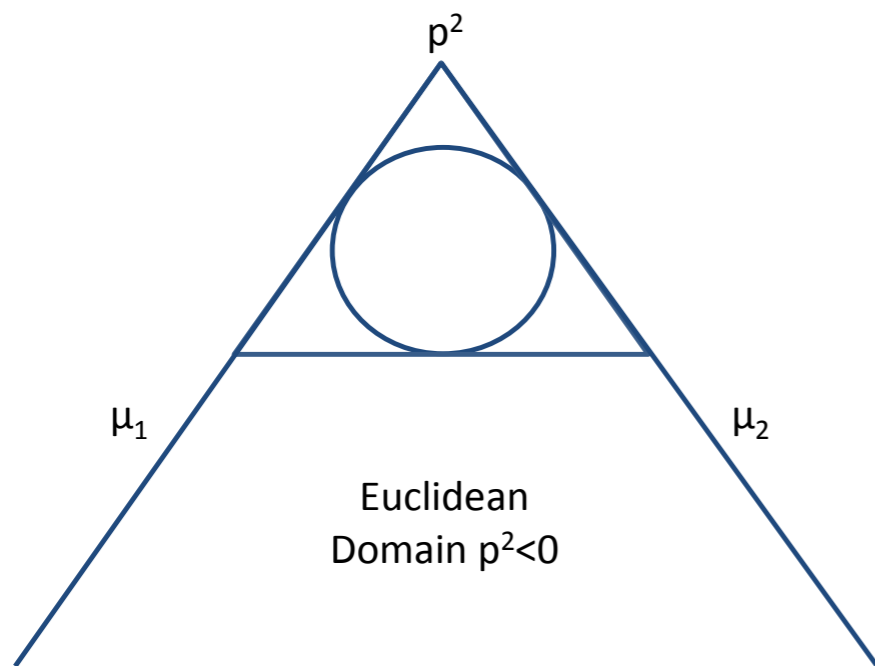
# SFI equation system

$$c^a I + (T^a)_j^i x_i \frac{\partial}{\partial x_j} I = J^a$$

$$a = 1, \dots, \dim G$$

- $c^a$  - x-indep. const's
- $T^a$  - group generators, acting on  $X$
- $J^a$  - simpler diagrams
- Symmetries of Feynman Integrals - SFI
- Closely related to Integration By Parts (IBP) and Differential Equations (DE) methods.

# Foliation of $X$ into $G$ orbits



# Singularity locus

$$I(x) = \sum_{\alpha} c_{\alpha}(x) J^{\alpha}(x)$$

A linear combination of simpler diagrams.

# General theory - SFI

## Summary of main results

$$\hat{I}(x) = \hat{I}(x_0) + \int_{x_0}^x J^\alpha(\xi) d\xi_\alpha$$

$$c^a I + (T^a)^i_j x_i \frac{\partial}{\partial x_j} I = J^a$$

$$I(x) = \sum_{\alpha} c_{\alpha}(x) J^{\alpha}(x)$$



# Mechanism

# Current freedom

$$G \subset T_{L,n-1}$$

upper block triangular,

$L = \#$  loops,  $n = \#$  external legs

$$\delta l = l + p$$

$$\delta p = p$$

$l$  – loop currents

$p$  – external currents

Geometric characterization:  $G$  is defined by preserving the subspace of squares inside quadratics

$$Q = Sp\{l_a \cdot l_b, l_a \cdot p_r, p_r \cdot p_s\}$$
$$a, b = 1, \dots, L, \quad r, s = 1, \dots, n - 1$$

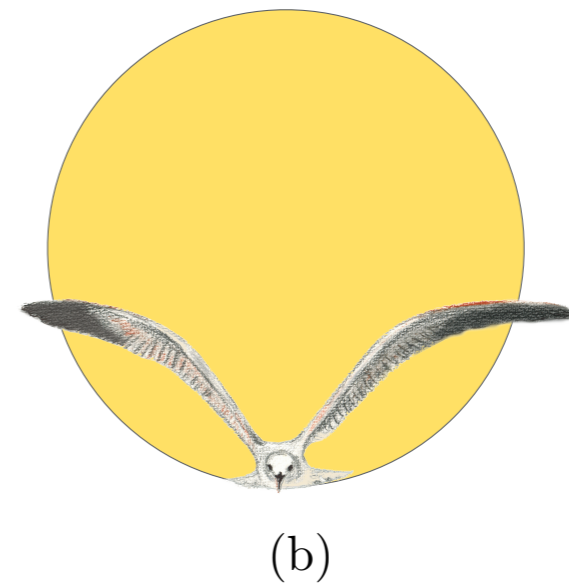
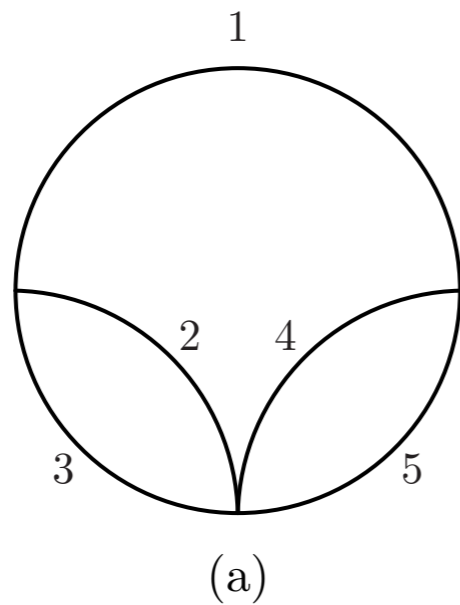
$$S \subset Q$$

$$S = Sp\{k_i^2, p_r \cdot p_s\}$$
$$i = 1, \dots, P$$

$$Num = Q/S$$

# Applications

# Vacuum seagull



evaluated a sector w. 3 mass scales (novel).

$$\begin{aligned}
I(m^2, m^2x, 0, m^2y, 0) &= (m^2)^{\frac{3d}{2}-5} I_0(x, y, d) \\
&(I_1(x, y, d) + I_1(y, x, d) + I_2(x, d) + I_2(y, d) \\
&+ I_3(x, y, d) + c_4(d))
\end{aligned} \tag{53a}$$

where

$$I_0(x, y, d) = i\pi^{\frac{3d}{2}}((1-x)(1-y))^{d-3}, \tag{53b}$$

$$\begin{aligned}
I_1(x, y, d) &= \left[ c_{1a}(d)xy^{\frac{3d}{2}-5} {}_2F_1\left(5 - \frac{3d}{2}, 4 - d, 3 - \frac{d}{2} \middle| \frac{x}{y}\right) \right. \\
&+ c_{1b}(d)x^{\frac{d}{2}-1}y^{d-3} {}_2F_1\left(3 - d, 2 - \frac{d}{2}, \frac{d}{2} - 1 \middle| \frac{x}{y}\right) \left. \right] \\
&\times F_1(3d/2 - 4, d - 2, d - 3, 3d/2 - 3 | x, y),
\end{aligned} \tag{53c}$$

$$c_{1a}(d) = -3\Gamma\left(3 - \frac{3d}{2}\right)\Gamma(4 - d)\Gamma\left(\frac{d}{2} - 2\right)\Gamma\left(\frac{d}{2} - 1\right), \tag{53f}$$

$$c_{1b}(d) = -\frac{4\pi\text{Csc}\left(\frac{\pi d}{2}\right)\Gamma(2 - d)\Gamma(2 - \frac{d}{2})}{3d - 8}, \tag{53g}$$

$$I_2(x, d) = c_2(d)x^{\frac{d}{2}-1} {}_2F_1(d/2 - 1, d - 2, d/2 | x), \tag{53d}$$

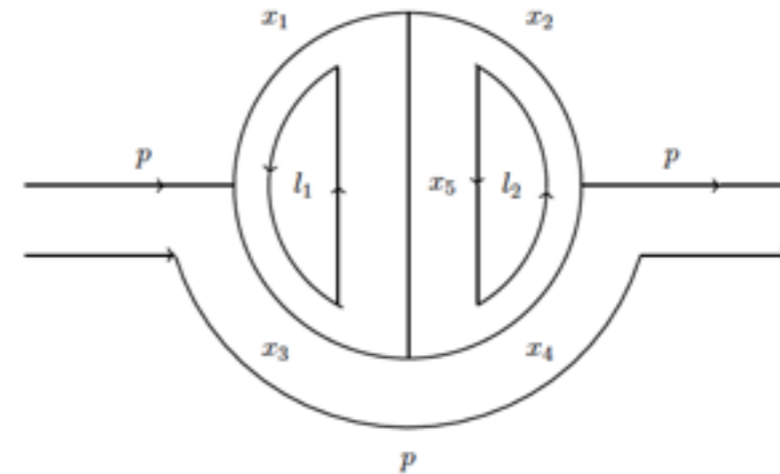
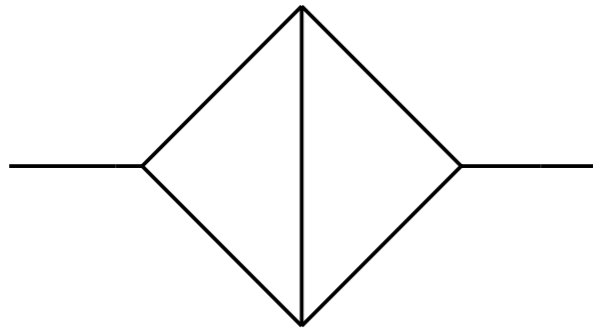
$$c_2(d) = \pi d \text{Csc}\left(\frac{\pi d}{2}\right)\Gamma(2 - d)\Gamma\left(-\frac{d}{2}\right), \tag{53h}$$

$$\begin{aligned}
I_3(x, y, d) &= c_3(d)(xy)^{\frac{d}{2}-1} {}_2F_1(d/2 - 1, d - 2, d/2 | x) \\
&\times {}_2F_1(d/2 - 1, d - 2, d/2 | y),
\end{aligned} \tag{53e}$$

$$c_3(d) = \Gamma\left(1 - \frac{d}{2}\right)^3. \tag{53i}$$

$c_4(d)$  was defined in (48e).

# Kite - singularity locus



$$B_3(x) = 0$$

$$\begin{aligned}
 B_3 = & x_1 x_4(x_1 + x_4) + x_2 x_3(x_2 + x_3) + x_5 x_6(x_5 + x_6) + \\
 & + x_1 x_2 x_5 + x_1 x_3 x_6 + x_2 x_4 x_6 + x_3 x_4 x_5 + \\
 & - (x_1 x_4(x_2 + x_3 + x_5 + x_6) + x_2 x_3(x_1 + x_4 + x_5 + x_6) + x_5 x_6(x_1 + x_2 + x_3 + x_4))
 \end{aligned} \tag{3.2}$$

$$I|_{B_3} = -\frac{1}{d-4} \frac{\vec{u} \cdot \vec{J}}{\partial^5 B_3} = \dots$$

Generalizes the massless case [ChetyrkinTkachov1981], and the more general “diamond rule” [Ruijl, Ueda and Vermaseren 2015].

Singularity locus



More results

# Symmetry adapted basis

Equation basis chosen to be compatible with the standard discrete symmetry of the diagram

# Maximal minors

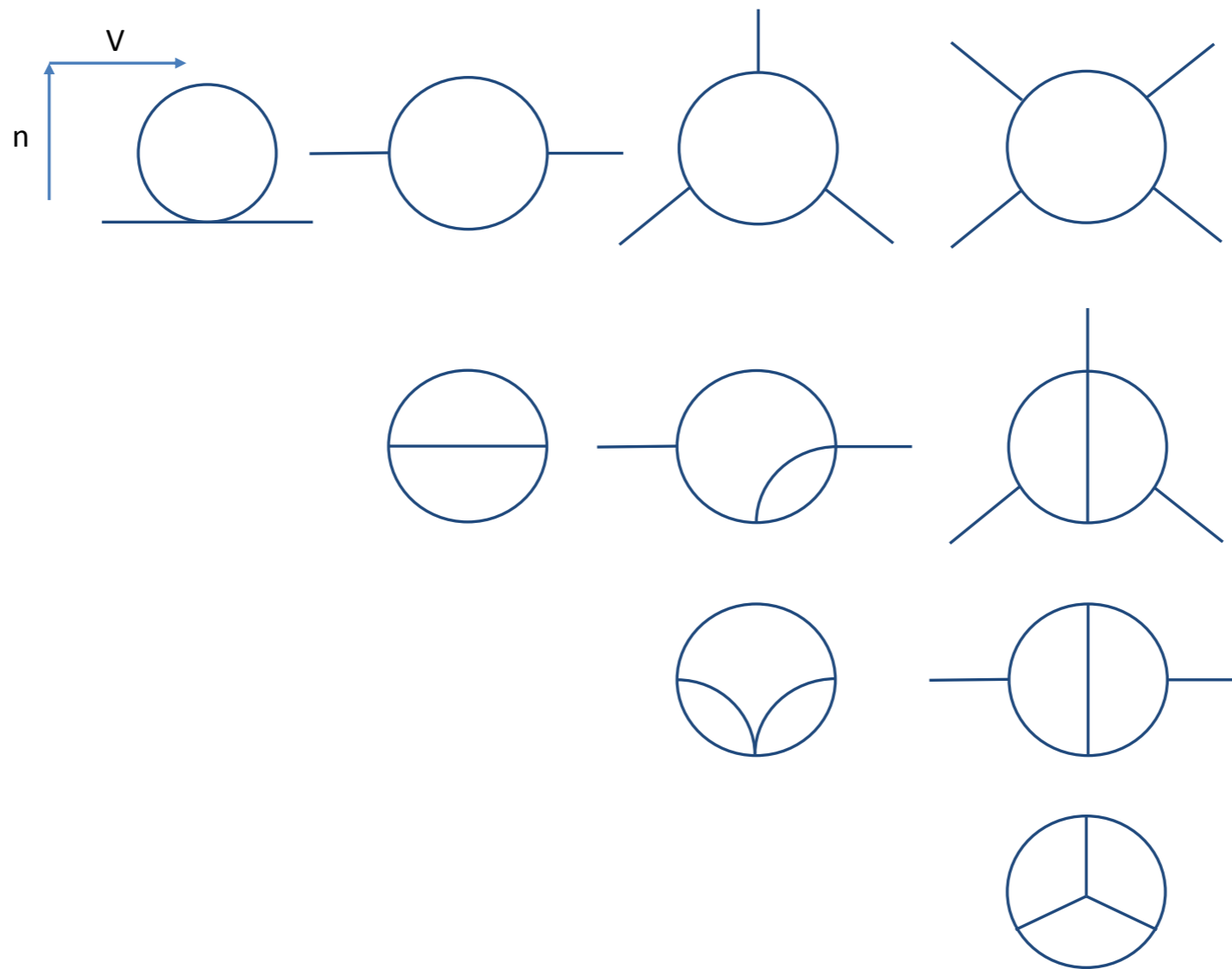
- Factorisation of maximal minors

$$M^I_A(x) = \epsilon_{Aa_1 \dots a_r} \epsilon^{Ii_1 \dots i_r} T x_{i_1}^{a_1} \dots T x_{i_r}^{a_r}$$

$$M^I_A(x) = S(x) Orb^I(x) Stb_A(x)$$

Partial group invariants, stabilizers

# Hierarchy of diagrams



# Conclusions

# Conclusions

- Every diagram defines an SFI group  $G$ .
- The group acts on parameters and allows to change parameters at the price of a line integral over simpler diagrams.
- Achieved a novel and general theory.
- Continuing to apply and refine the method.

# Conclusions

Even 70 years later innovation is possible.

Thank you for your  
attention!