



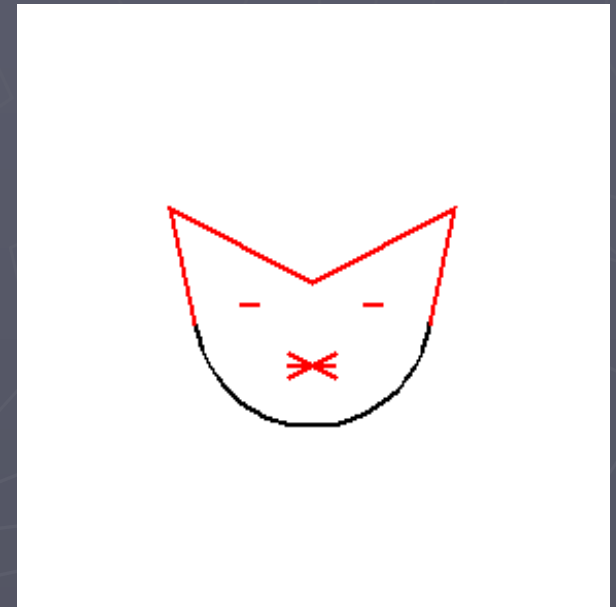
# The Baron and the cat



Joint work with Oded Kenneth  
Original idea: J. Wisdom

# Baron and cat

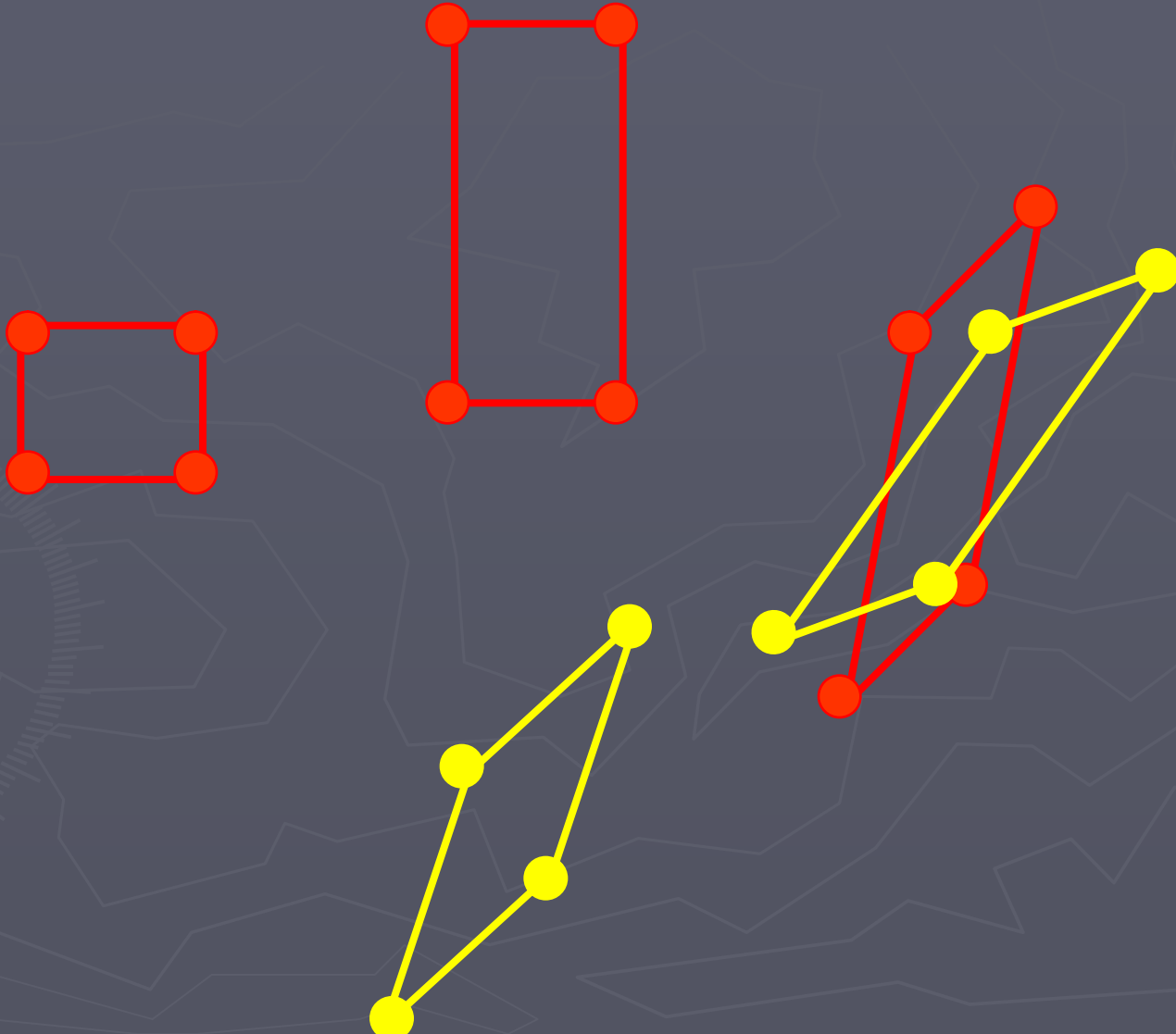
- ▶ A cat can rotate with zero angular momentum
- ▶ Can Baron von Munchausen lift himself?
- ▶ Translate with zero momentum?



Movie: omri Gat


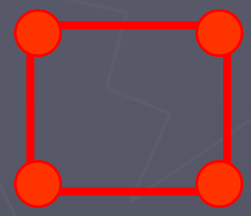
H. Knoerer: Cats always fall on their feet

# Rotations without angular momentum: Physics



# Rotations without angular momentum: Mathematics

Generators of deformations


$$X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$


$$[X, Z] = 2 \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} = 2R$$

The commutator of deformations= a rotation

# Newton's law

If a system experiences no external force,  
the center-of-mass of the system will remain at rest.

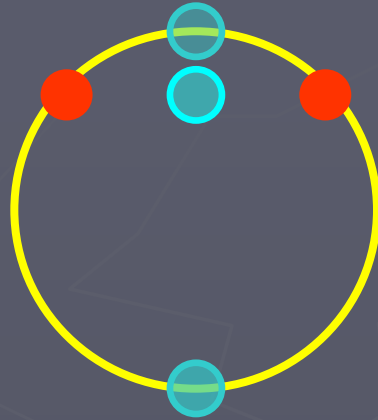
[http://webphysics.davidson.edu/physlet\\_resources/bu\\_semester1/c12\\_cofm\\_motion.html](http://webphysics.davidson.edu/physlet_resources/bu_semester1/c12_cofm_motion.html)

Newton:

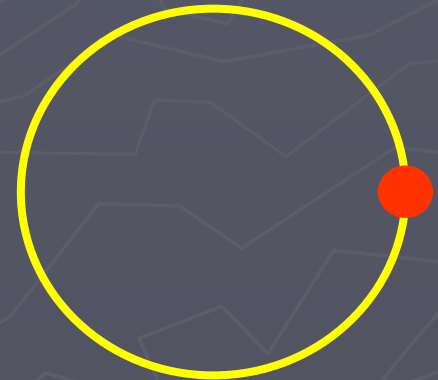
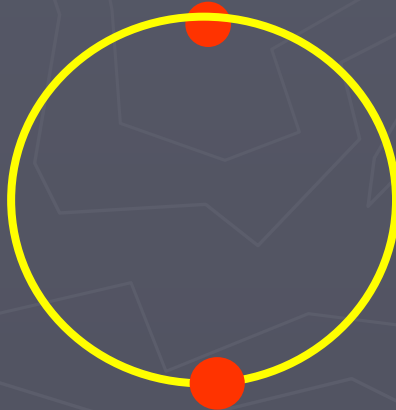
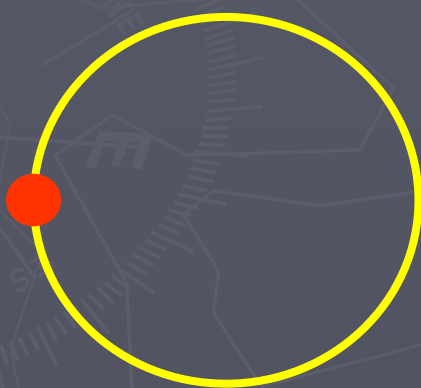
Lex I: Corpus omne perseverare in statu suo quiescendi vel  
Movendi uniformiter in directum, nisi quatenus a viribus  
impressis cogitur statum illum mutare.

Every body perseveres in its state of being at rest or of moving  
uniformly straight forward, except insofar as it is compelled to  
change its state by force impressed

# Ambiguous center of mass: Topology

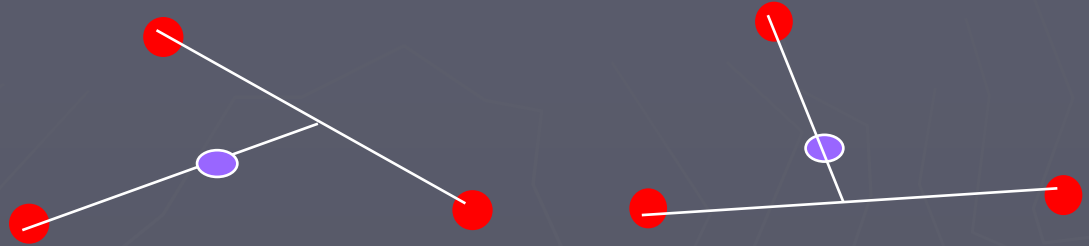


Ambiguous center of mass allows the Baron to move itself

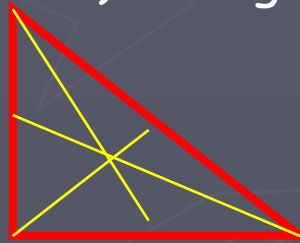


# Ambiguous center of mass: Geometry

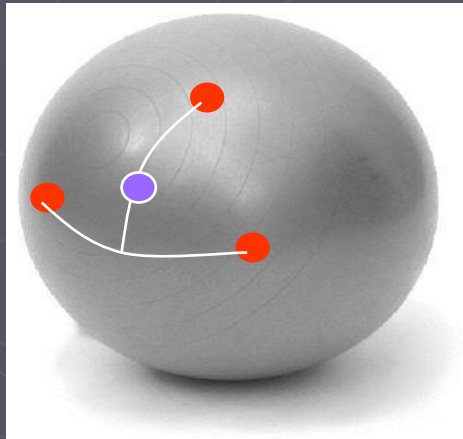
Euclidean plane



The three medians of a (Euclidean) triangle intersect at one point

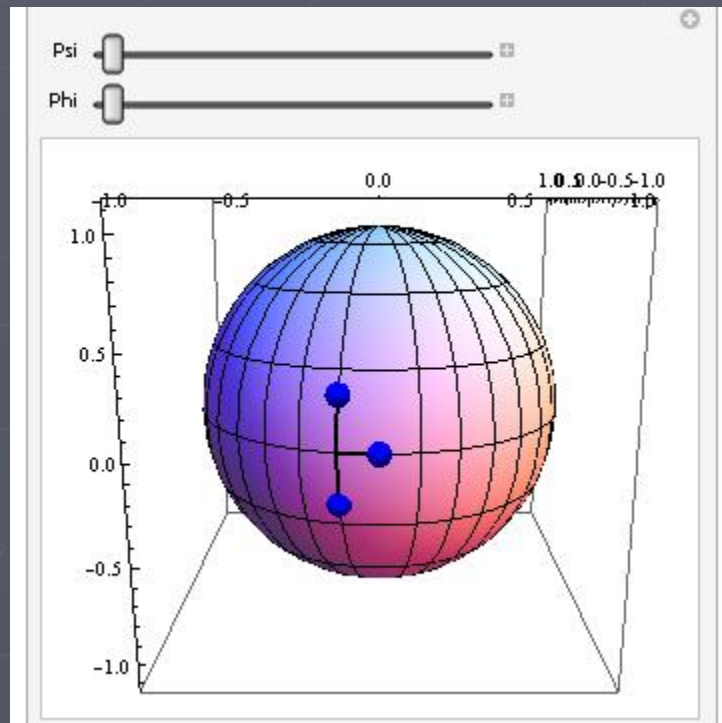


On a sphere



# Swimming triangles

On sphere

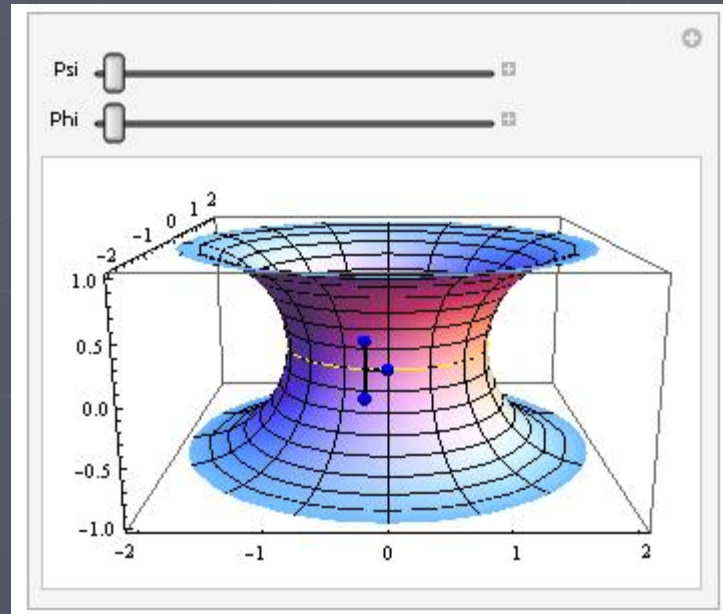


Movies: Oren Raz



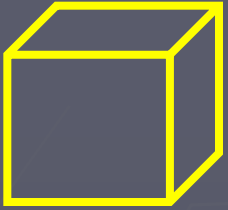
# Swimming triangles

Inside a torus



Movies: Oren Raz

# Shape fixes location



Shape: All mutual distances  $\binom{n}{2}$

Location:  $nd$  coordinates

$n > 2d + 1$  overconstrained.

Large rigid bodies generically fits nowhere

Generically, shape fixes a location

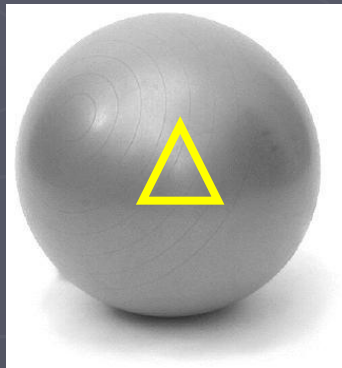


# When can rigid bodies move?

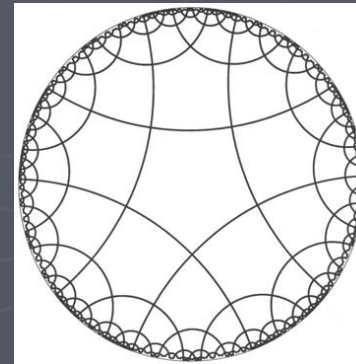
Euclidean space is a symmetric space: Homogeneous and isotropic.

In a symmetric space if a rigid body fits somewhere, it fits any where.

Examples:



Sphere



Lobachevski plane

In a symmetric space, shape does not determine location. You may try to swim.

# Killing forms

Killing for Euclidean translation

$$\xi_t = dx$$



Killing for Euclidean rotation

$$\xi_R = xdy - ydx$$



# Controls: Deformations

Example: Euclidean case



$$\eta_1 = x\partial_x - y\partial_y$$



$$\eta_2 = x\partial_y + y\partial_x$$

Vector field for deformations depend on  
choice of origin: Ambiguity in Killing

$$\eta_1 \rightarrow \eta_1 + a\xi_x - b\xi_y$$

# Main result

Swimming distance

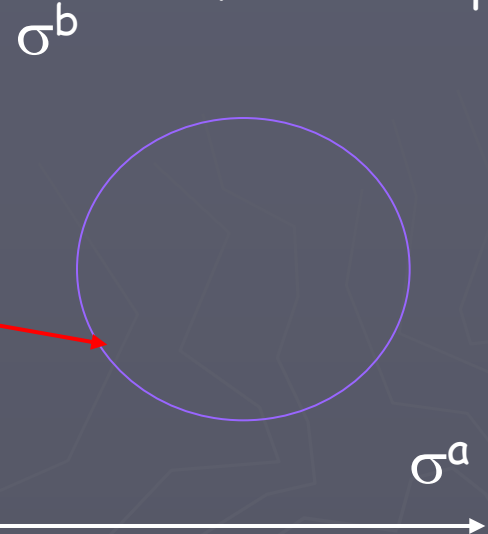
$$\dot{d}\tau = \langle d\xi | \eta_a, \eta_b \rangle d\sigma^a \wedge d\sigma^b$$

Killing 2-form

Deformation fields

Stroke

Deformation space



Space allows for swimming if the Killing 1-form is not closed

# Cats spin, the Baron lies

Translations  $\xi_t = dx \longrightarrow d\xi_t = 0$

The Baron lies

Rotations

$$\xi_R = xdy - ydx \longrightarrow d\xi_R = 2dx \wedge dy$$

The cat falls on its feet

# Small swimmers on symmetric surface

Property of space

Size of stroke

$$\delta x \approx 8R \left( \frac{\sum m_n x_n y_n^2}{M} \right) dA$$

Geometry of a swimmer

Wisdom

Swim away from Black hole  
Relativistic motions

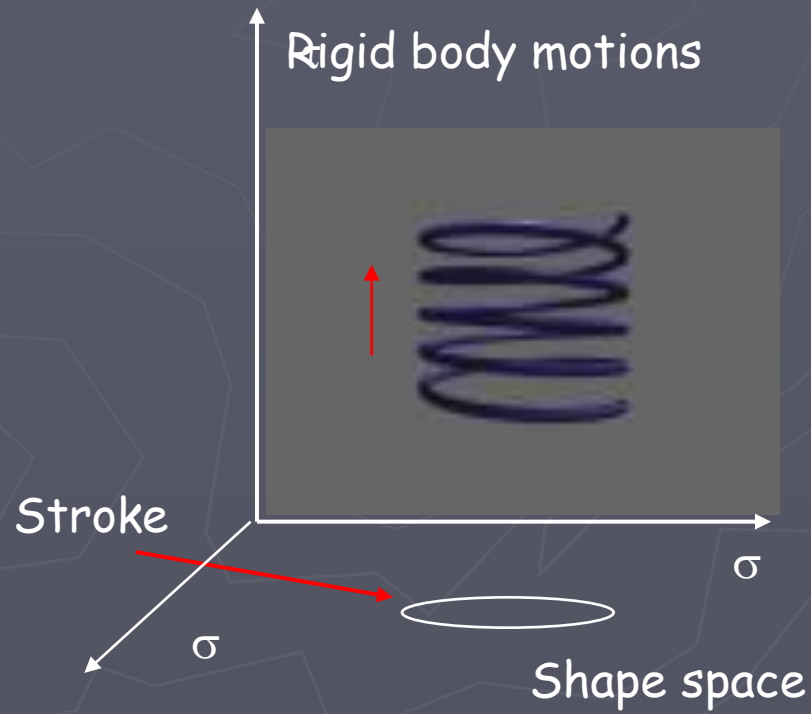
AK:

Symmetric spaces  
Non relativistic bodies



# The swimming problem

Swimming with periodic strokes:



# Conservation of momentum

A system of particles

$$P_\xi = \sum_n p_n \cdot \xi(x_n), \quad p_n = \frac{\partial L}{\partial \dot{x}_n} = m\dot{x}_n$$

Killing vector field

Swimming equations

$$P_\xi = 0 \longrightarrow \sum m_n \xi(x_n) \cdot dx_n = 0$$

As many equations as Killing fields:  
Can't move in other directions

Eq. independent of time parameterization: Geometric

# Killing fields: Translation & rotations

Defining property: No strain

$$\xi_{\mu;\nu} + \xi_{\nu;\mu} = 0$$

Relation to Riemann

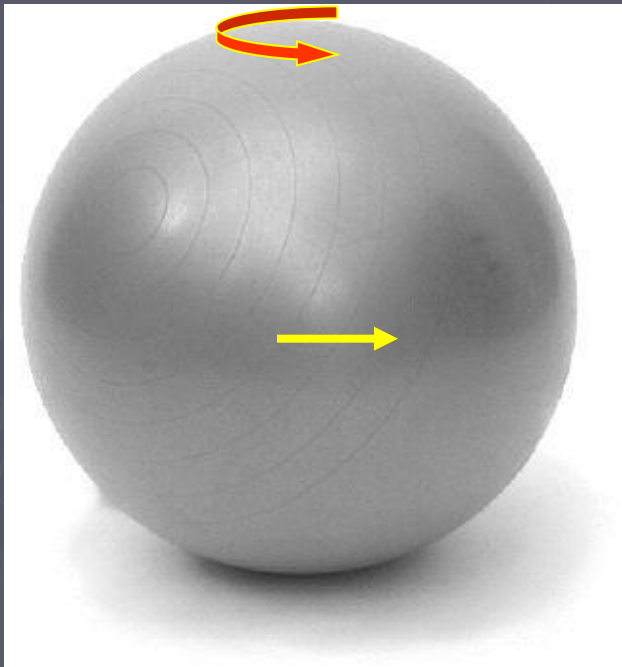
$$\xi_{\ell;i;j} = -R_{\ell j i k} \xi^k$$

Determined by initial data at a point

$$\xi_{\mu}(0), \quad \xi_{[\mu,\nu]}(0)$$

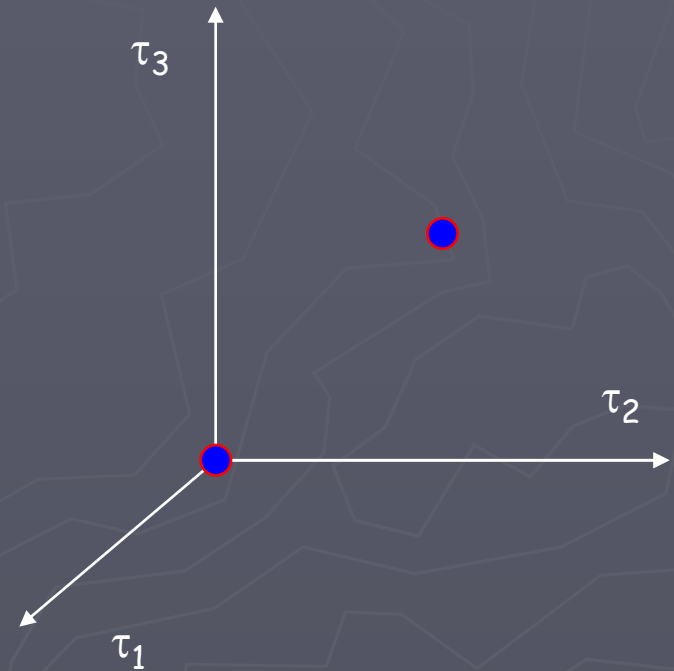
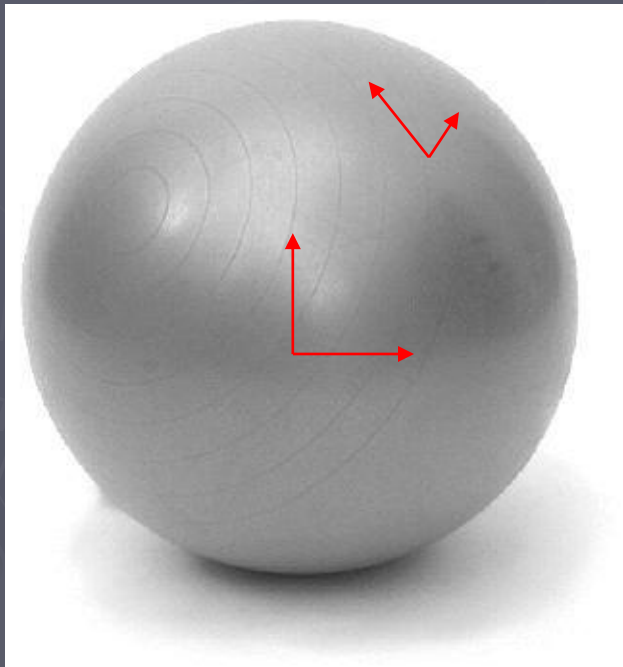
translation

Rotation



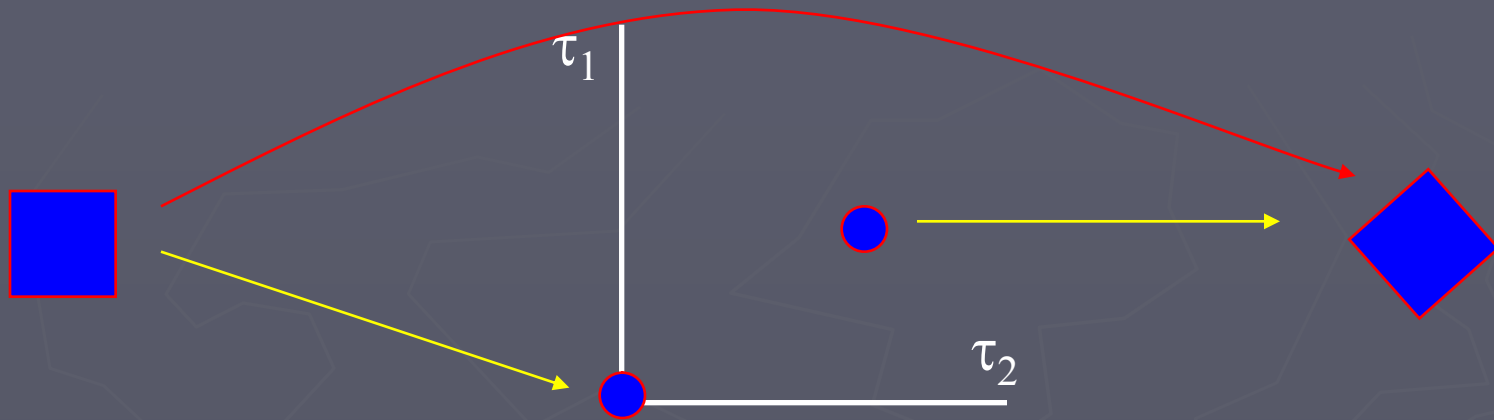
Rotation about z axis:  
Looks like translation near equator  
And like rotation at poles

# Rigid body motion: Coordinates



Putting coordinates for  
The analog of Euclidean  
motions

# Coordinates for rigid body motion



$$\dot{x}(t) = \tau \cdot \xi(x(t)), \quad x(0) \rightarrow x_\tau(1)$$

$$S(\tau) = e^{\tau \xi} S$$

Symbolic notation for transported shape

# Gauge choice instead of CM

$$\langle \xi | \eta \rangle = \frac{1}{M} \sum m_n \xi(x_n) \cdot \eta(x_n)$$

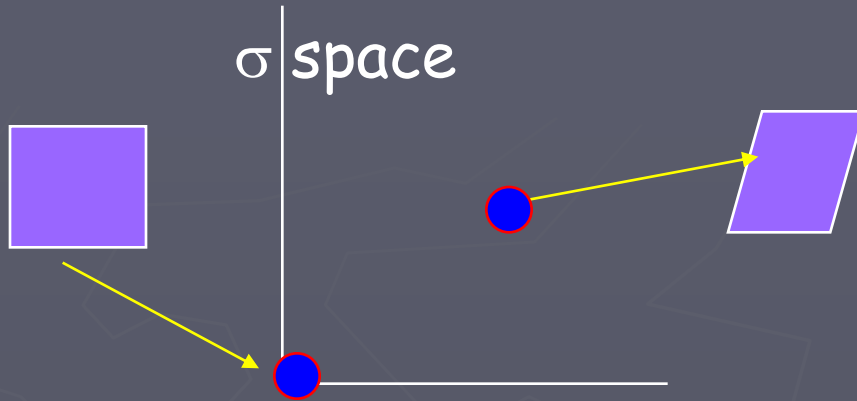
Euclidean example:

$$\langle \xi_1 | \eta_1 \rangle = \frac{1}{M} \sum m_n x_n$$

Gauge choice: Analog to choosing cm as fiducial pt  
for translations

$$\langle \xi_j | \eta_a \rangle = 0$$

# Coordinates for deformed shapes



$$\dot{x}(t) = \sigma \cdot \eta(x(t)), \quad x(0) \rightarrow x_\sigma(1)$$

Symbolically

$$S(\sigma) = e^{\sigma \eta} S$$

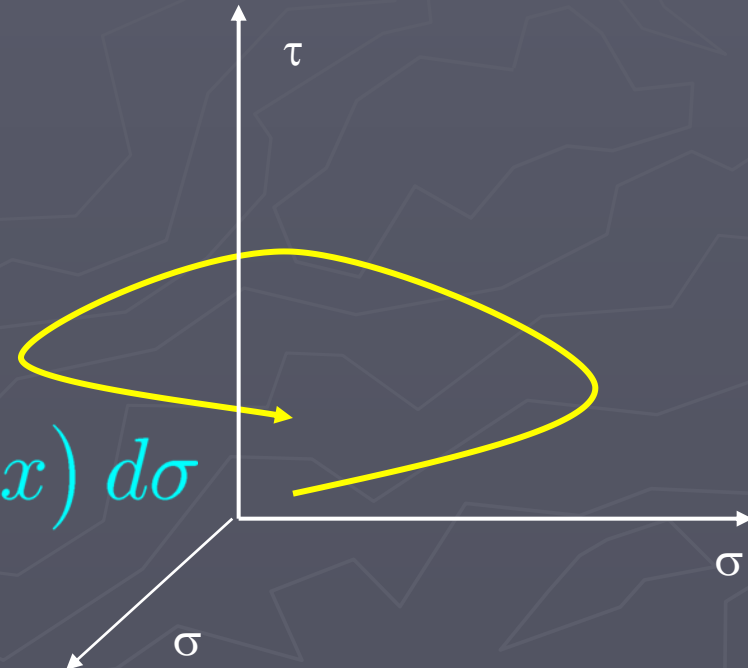
# Total space

Coordinates for deformations + Euclidean motions

$$S(\tau, \sigma) = e^{\tau\xi} e^{\sigma\eta} S$$

$$x_0 \longrightarrow x(\tau, \sigma; x_0)$$

$$dx(\tau, \sigma; x_0) = (\partial_\tau x) d\tau + (\partial_\sigma x) d\sigma$$



Motions is legit if it consistent with zero total momentum



# Equation of Motion

To leading order with  $\sigma$  and  $\tau$  small

$$dx = \xi d\tau + \eta d\sigma, \quad \xi(x(\sigma, \tau)), \quad \eta(x(\sigma, \tau)),$$

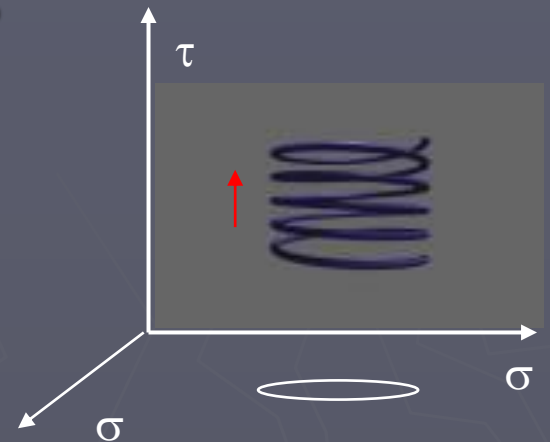
Substitute in conservation law

$$\sum m_n \xi(x_n) \cdot dx_n = 0$$

Gives a linear (system) of equations  
for  $d\tau$

# Main result

Swimming distance



$$\vec{d}\tau = \langle d\xi | \eta_a, \eta_b \rangle d\sigma^a \wedge d\sigma^b$$

Killing fields

Deformation fields

Shape space

$$\vec{d}\tau^\alpha = \langle \partial_{[k} \xi_j^\alpha | \eta_a^k \eta_b^j \rangle d\sigma^a \wedge d\sigma^b$$

# Role of Curvature

Killing field for symmetric spaces can be  
Explicitly expressed in terms of the curvature

$$\xi_{\ell;i;j} = -R_{\ell j i k} \xi^k$$

In a symmetric space

$$R_{ijkl} = R(\delta_{ik}\delta_{jl} - \delta_{il}\delta_{jk})$$

Swimming in the direction of the k-th Killing field

$$\delta x^k = 2R \left( Q^k d\sigma^a \wedge d\sigma^b \right)$$

Geometry of embedding space

$$Q^k = \frac{1}{M} \sum m_n x_n^i \eta_a^i(x_n) \eta_b^k(x_n)$$

Stroke and swimmer data

# Application to astrophysics?

How does swimming affect the motion of galaxies of size  $x$ ?

$$\frac{\Delta x}{x} \propto R (\delta x^2) \approx R v^2 (dt)^2$$

Hubble constant 75 km/sec/Mpc gives a measure of the curvature

$$R \sim H_0^2 \approx 10^{-20} [\text{year}]^{-2}$$

Perhaps not surprisingly, it takes the age of the universe for a significant swimming

# The hero: Oded Kenneth



Thanks to Amos Ori

