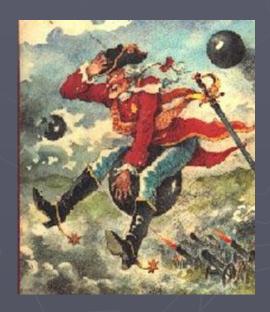


The Baron and the cat



Oded Kenneth

Baron and cat

- A cat can rotate with zero angular momentum
- Can Baron von Munchausen lift himself?
- Translate with zero momentum?



Movie: omri Gat

H. Knoerer: Cats always fall on their feet

Newton's law

If a system experiences no external force, the center-of-mass of the system will remain at rest.

http://webphysics.davidson.edu/physlet_resources/bu_semester1/c12_cofm_motion.html

Newton:

Lex I: Corpus omne perseverare in statu suo quiescendi vel Movendi uniformiter in directum, nisi quatenus a viribus impressis cogitur statum illum mutare.

Every body perseveres in its state of being at rest or of moving uniformly straight forward, except insofar as it is compelled to change its state by force impressed

Wikipedia

Respecting the third law

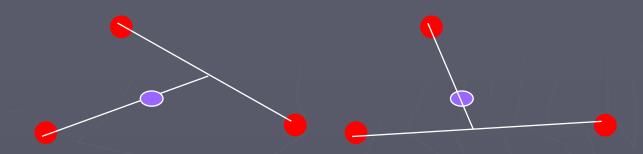


- · Internal forces equal and opposite
- · Guarantee conservation of total momentum

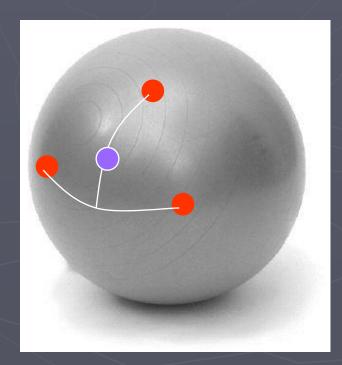
It is possible to swim in empty curved space at zero total momentum

Ambiguity in center of mass

Euclidean plane

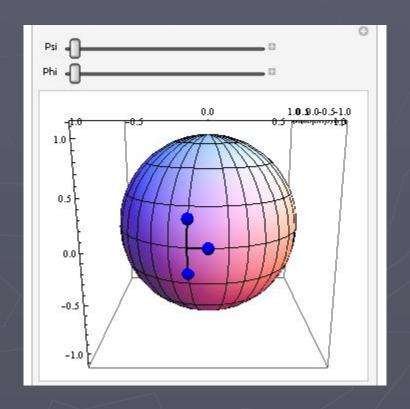


On a sphere



Swimming triangles

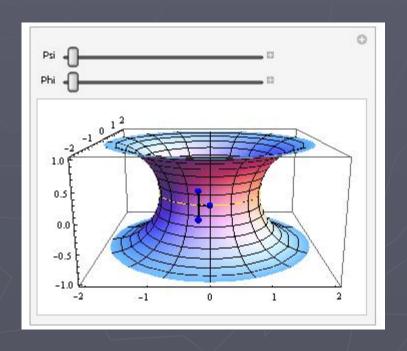
On sphere



Movies: Oren Raz

Swimming triangles

Inside a torus



Movies: Oren Raz

Main result

Deformation space

 σ^{a}

Stroke

Swimming distance

 $\overline{d\tau} = \langle d\xi | \eta_a, \eta_b \rangle \ d\sigma^a \wedge d\sigma^b$

Killing fields

Deformation fields

Why cats rotate while Baron lies

$$\xi_t = dx \longrightarrow d\xi_t = 0$$

The Baron lies

Rotations

$$\xi_R = xdy - ydx \longrightarrow d\xi_R = 2dx \wedge dy$$

The cat falls on its feet

Small swimmers: Wisdom's type formula

Property of space

How hard swims

$$\delta x_{\mu} \sim R_{\mu\nu\alpha\beta} \ x^{\nu} \ dx^{\alpha} \wedge dx^{\beta}$$

Wisdom
Swim away from Black hole
Relativistic motions

AK: Symmetric spaces Non relativistic bodies

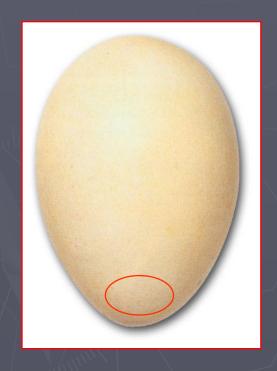
Outline

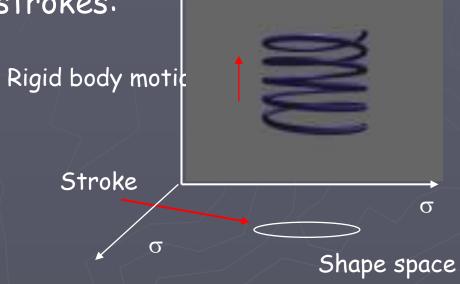
- Setting the problem
- Killing and deformation fields
- Swimming in empty space

Wisdom, Science AKNJP

The swimming problem

Swimming with periodic strokes:



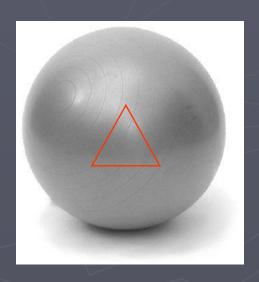


Shapes are hard to fit in a curved space

Periodic strokes bring you back home

Swimming in symmetric spaces

- Symmetric spaces: Isotropic and homogeneous
- Examples: Euclidean plane, Sphere, Hyperbolic plane



Killing fields: Euclidean case

Killing for Euclidean translation

$$\xi_k = \partial_k$$



Killing for Euclidean rotation

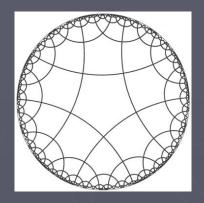
$$\xi_{\theta} = x_1 \partial_2 - x_2 \partial_1$$



Killing fields for Symmetric planes







Metric

$$(d\ell)^2 = \frac{|dz|^2}{(1+R|z|^2)^2}, \quad R = 0, \pm 1$$

Killing fields

$$\xi_1 = \partial_x + R \left((x^2 - y^2) \, \partial_x + 2xy \partial_y \right)$$

$$\xi_2 = \partial_y + R \left(2xy \, \partial_x + (y^2 - x^2) \, \partial_y \right)$$

$$\xi_3 = x \partial_y - y \partial_x$$

Conservation of momentum

A system of particles

$$P_{\xi} = \sum_{n} p_n \cdot \xi(x_n), \quad p_n = \frac{\partial L}{\partial \dot{x}_n} = m \dot{x}_n$$

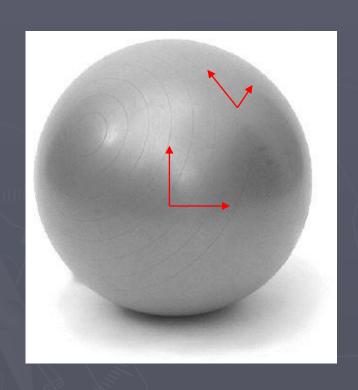
Swimming equations

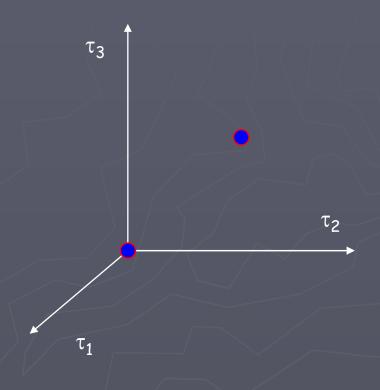
$$P_{\xi} = 0 \longrightarrow \sum m_n \xi(x_n) \cdot dx_n = 0$$

As many equations as Killing fields: Can't move in other directions

Eq. independent of time parameterization: Geometric

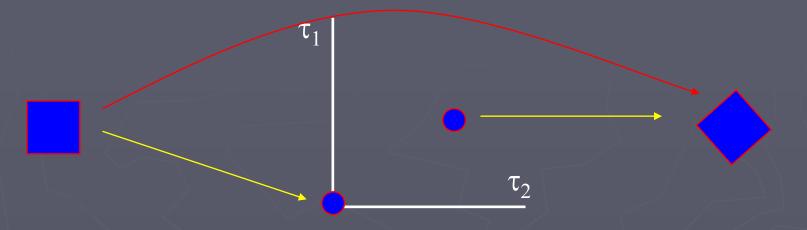
Rigid body motion: Coordinates





Putting coordinates for The analog of Euclidean motions

Coordinates for rigid body motion



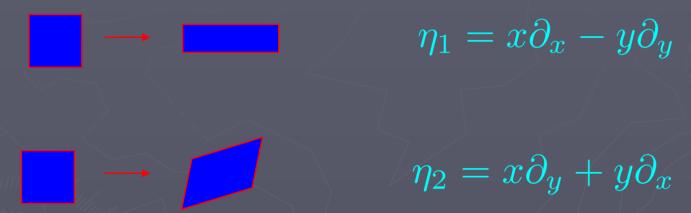
$$\dot{x}(t) = \tau \cdot \xi(x(t)), \quad x(0) \to x_{\tau}(1)$$

$$S(\tau) = e^{\tau \xi} S$$

Symbolic notation for transported shape

Controls: Deformations

Example: Euclidean case



Vector field for deformations depend on choice of origin: Ambiguity in Killing

$$\eta_1 \to \eta_1 + a\xi_x - b\xi_y$$

Gauge choice instead of CM

$$\langle \xi | \eta \rangle = \frac{1}{M} \sum_{n} m_n \xi(x_n) \cdot \eta(x_n)$$

Euclidean example:

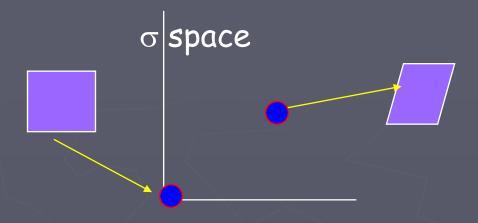
$$\langle \xi_1 | \eta_1 \rangle = \frac{1}{M} \sum m_n x_n$$

Gauge choice: Analog to choosing cm as fiducial pt

for translations

$$\langle \xi_j | \eta_a \rangle = 0$$

Coordinates for deformed shapes



$$\dot{x}(t) = \sigma \cdot \eta(x(t)), \quad x(0) \to x_{\sigma}(1)$$

Symbolically
$$S(\sigma) = e^{\sigma\eta}S$$

Total space

Coordinates for deformations + Euclidean motions

$$S(au, \sigma) = e^{ au\xi} e^{\sigma\eta} S$$

$$x_0 \longrightarrow x(au, \sigma; x_0)$$

$$dx(au, \sigma; x_0) = (\partial_{ au} x) d au + (\partial_{\sigma} x) d\sigma$$

Motions is legit if it consistent with zero total momentum

Eq. of Motion

To leading order with σ and τ small

$$dx = \xi d\tau + \eta d\sigma, \quad \xi(x(\sigma, \tau)), \quad \eta(x(\sigma, \tau)),$$

Substitute in conservation law

$$\sum m_n \xi(x_n) \cdot dx_n = 0$$

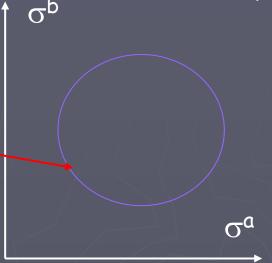
Gives a linear (system) of equations for $d\tau$

Main result

Deformation space

Swimming distance

Stroke



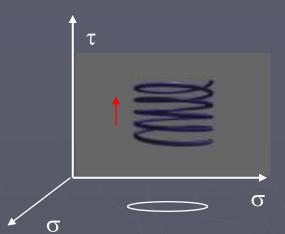
$$d\tau = \langle d\xi | \eta_a, \eta_b \rangle \ d\sigma^a \wedge d\sigma^b$$

Killing fields

Deformation fields

Main result

Swimming distance



$$d\tau = \langle d\xi | \eta_a, \eta_b \rangle \ d\sigma^a \wedge d\sigma^b$$

Killing fields

Deformation fields

Shape space

$$d\tau^{\alpha} = \langle \partial_{[k} \xi_{j]}^{\alpha} | \eta_a^k \eta_b^j \rangle \ d\sigma^a \wedge d\sigma^b$$

Riemann

Using

$$\nabla_i \nabla_j \xi_\ell = -R_{\ell j i k} \xi^k$$

One can write equations for transport for a small swimmer in terms of Riemann

$$M\delta x^{k} = R_{j\ell ik} \left(\sum_{i} m_{n} x_{n}^{i} \eta_{a}^{j}(x_{n}) \eta_{b}^{\ell}(x_{n}) \right) d\sigma^{a} \wedge d\sigma^{b}$$



The hero: Oded Kenneth



Thanks to Amos Ori