

# Physique mesoscopique des electrons et des photons -

## Structures fractales et quasi-periodiques

ERIC AKKERMANS  
PHYSICS-TECHNION



Aux frontieres de la physique mesoscopique,  
Mont Orford Quebec, Canada,

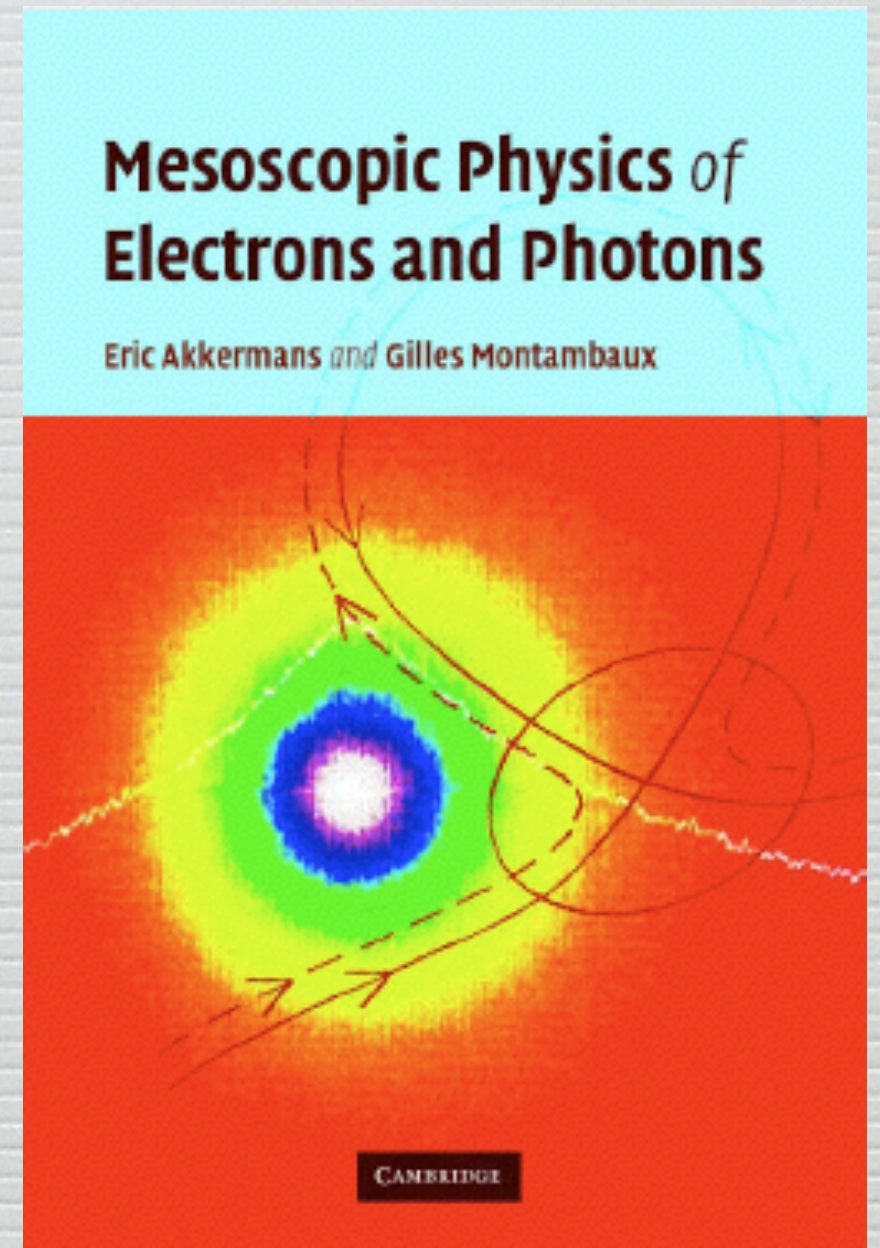


# Part 1

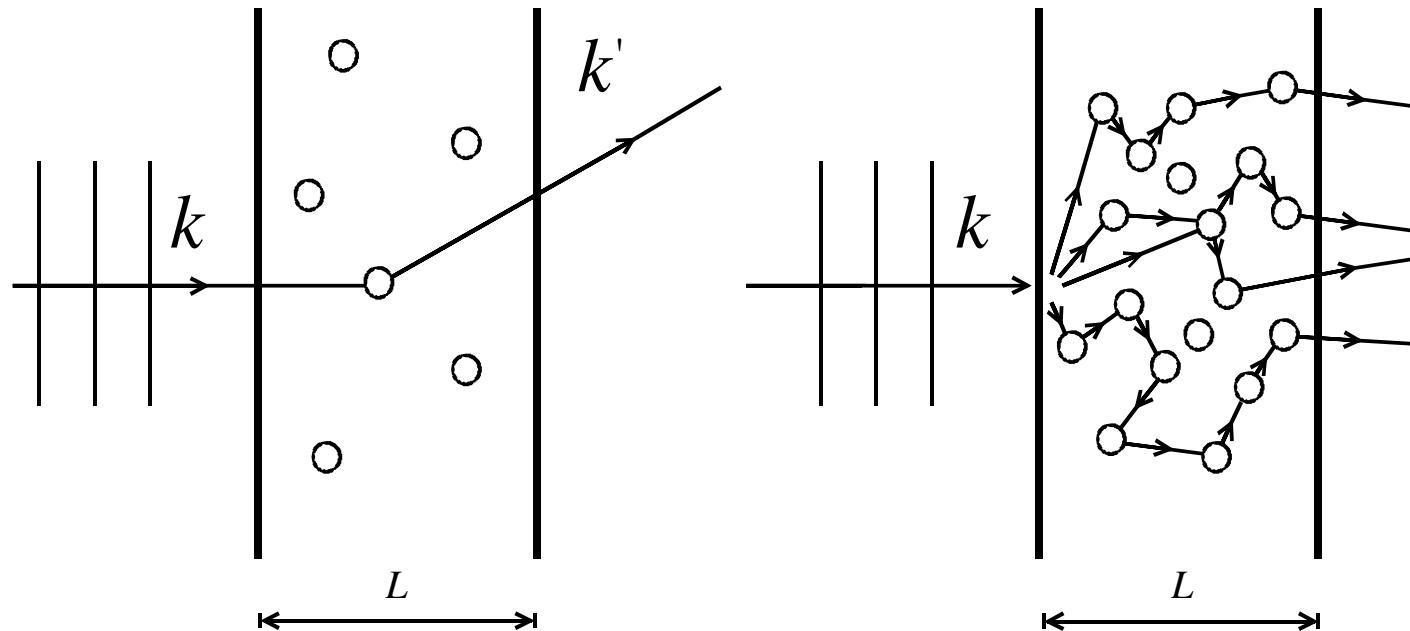
## What is quantum in mesoscopic physics ? Transport and interferences

(E.A., G. Montambaux)

more details in:



# Multiple scattering of electrons



2 characteristic lengths:

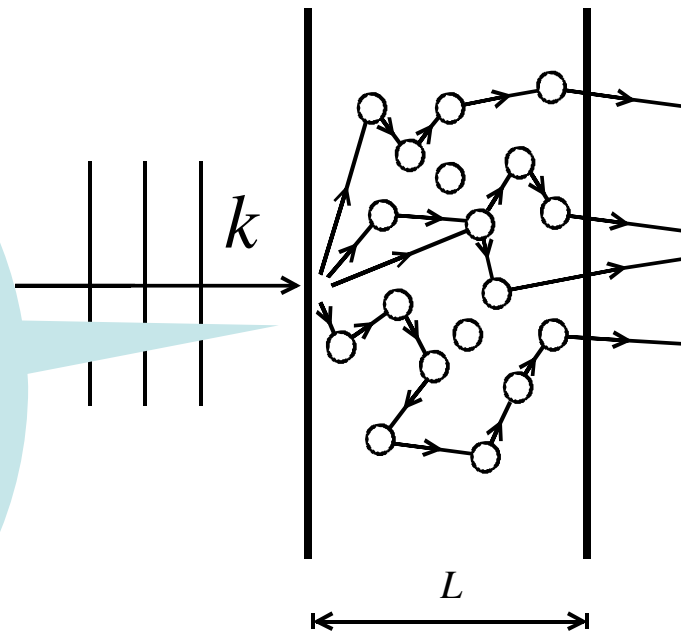
Wavelength:  $\lambda_F = k_F^{-1}$

Elastic mean free path:  $l$  (Disorder - Origin ?)

Weak disorder  $\lambda_F \ll l$  : independent scattering events

# Multiple scattering of electrons

We shall be interested only by this limit



2 characteristic lengths:

Wavelength:  $\lambda_F = k_F^{-1}$

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Some “canonical”  
mesoscopic effects



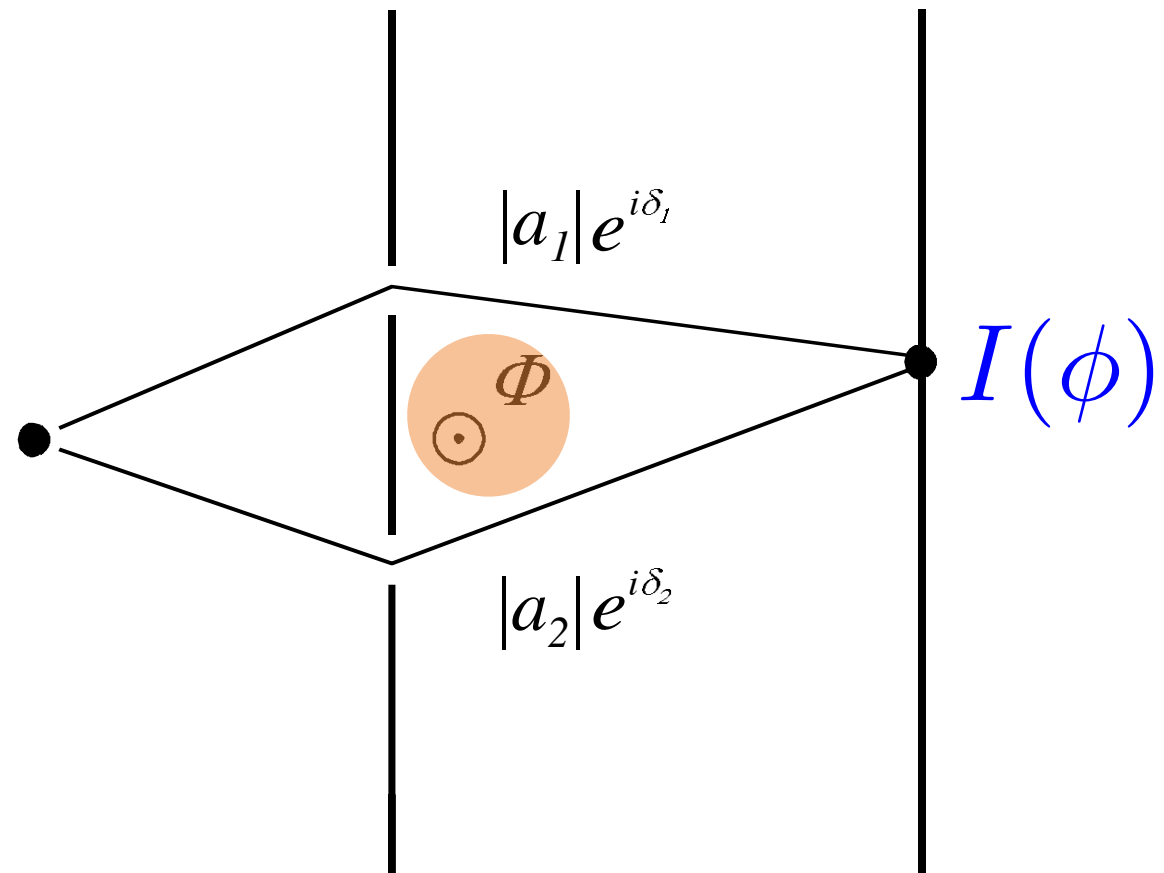
# Some “canonical” mesoscopic effects

## The Aharonov-Bohm effect

Aharonov-Bohm (1959)

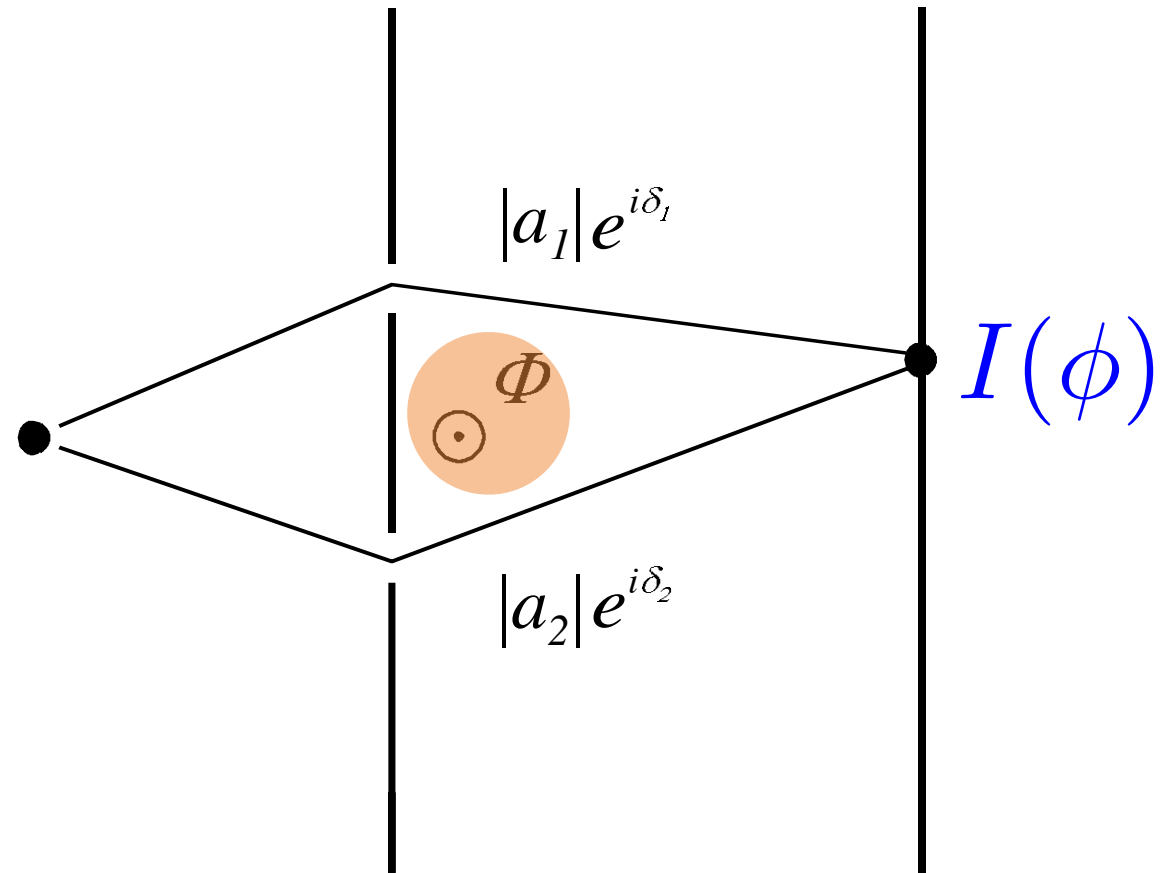
# The setup : Young slits

No magnetic field on  
the electrons : **no**  
**Lorentz force** and no  
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The quantum amplitudes  $a_{1,2} = |a_{1,2}|e^{i\delta_{1,2}}$  have phases:

$$\delta_1 = \delta_1^{(0)} - \frac{e}{\hbar} \int_1 \mathbf{A} \cdot d\mathbf{l} \quad \text{and} \quad \delta_2 = \delta_2^{(0)} - \frac{e}{\hbar} \int_2 \mathbf{A} \cdot d\mathbf{l}$$



The intensity  $I(\phi)$  is given by

$$\begin{aligned} I(\phi) &= |a_1 + a_2|^2 = |a_1|^2 + |a_2|^2 + 2|a_1 a_2| \cos(\delta_1 - \delta_2) \\ &= I_1 + I_2 + 2\sqrt{I_1 I_2} \cos(\delta_1 - \delta_2) \end{aligned}$$

The phase difference  $\Delta\delta(\phi) = \delta_1 - \delta_2$  is modulated by the magnetic flux  $\phi$  :

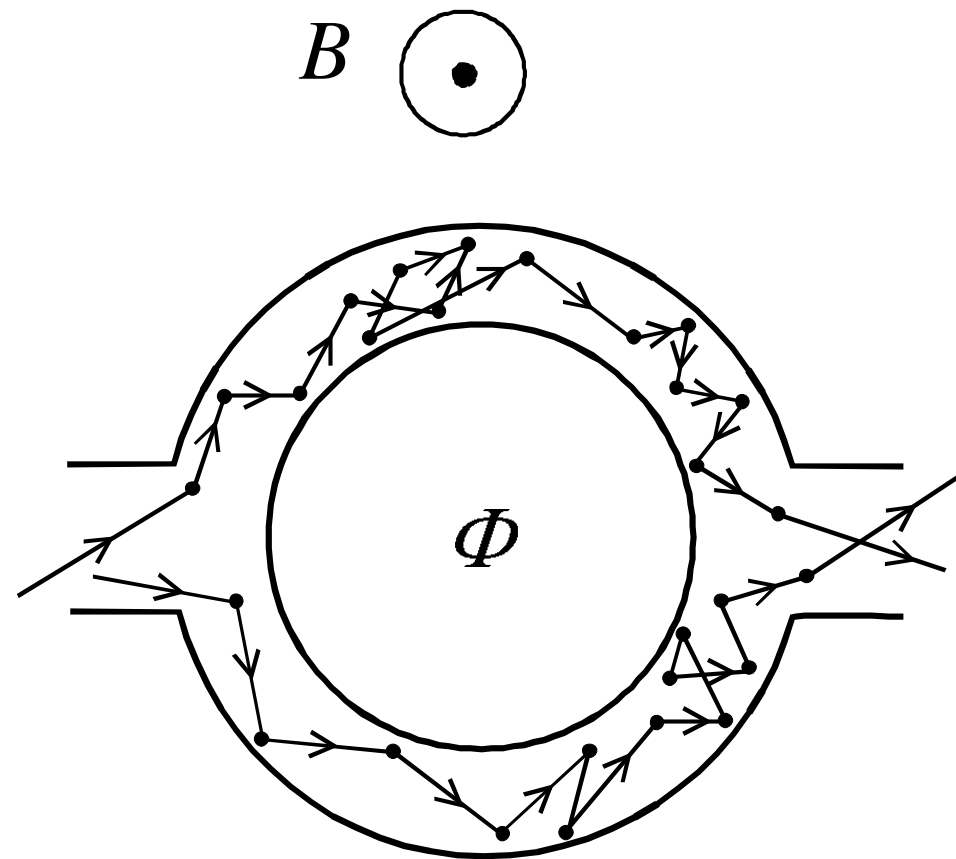
$$\Delta\delta(\phi) = \delta_1^{(0)} - \delta_2^{(0)} + \frac{e}{\hbar} \oint \mathbf{A} \cdot d\mathbf{l} = \Delta\delta^{(0)} + 2\pi \frac{\phi}{\phi_0}$$

where  $\phi_0 = h/e$  is the quantum of magnetic flux.

There is a continuous change of the state of interference:

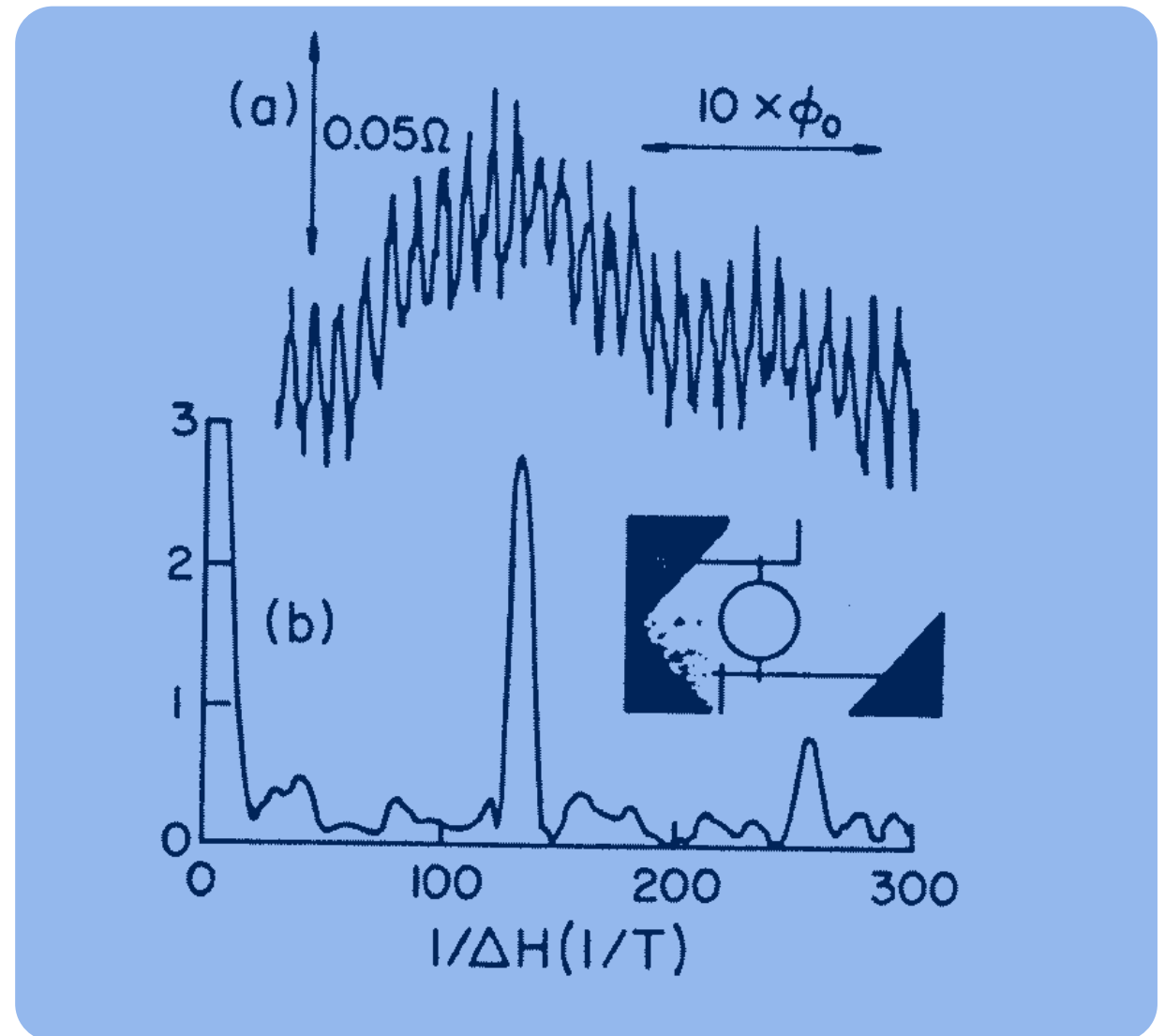
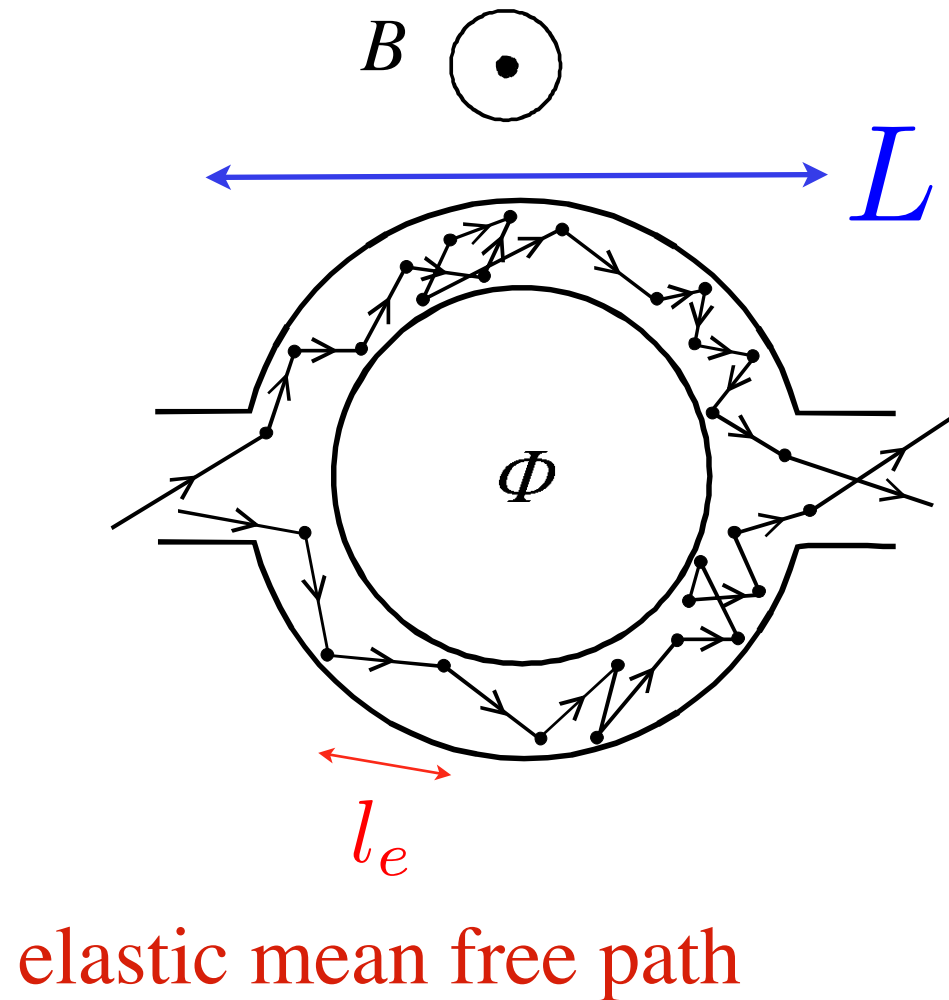
Aharonov-Bohm effect (1959)

- Aharonov-Bohm effect in disordered metals





**Implementation in metals** : the conductance  $G(\phi)$  is the analog of the intensity.



$$G(\phi) = G_0 + \delta G \cos(\Delta\delta^{(0)} + 2\pi \frac{\phi}{\phi_0})$$

*Webb et al. 1985*

Phase coherent effects subsist in **disordered** metals.  
**Reconsider the Drude theory.**

# Phase coherence and effect of disorder

The *Webb* experiment has been realized on a ring of size  $L \simeq 1\mu$ .  
For a macroscopic normal metal, coherent effects are washed out.



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*Quantum coherence*: gas of quantum particles in a finite volume

**Quantum states** of the gas are **coherent superposition** of single particle states and they extend over the total volume (*ex. superconductivity, superfluidity, free electron gas*).



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*Quantum coherence*: gas of quantum particles in a finite volume

**Quantum states** of the gas are **coherent superposition** of single particle states and they extend over the total volume (*ex. superconductivity, superfluidity, free electron gas*).

For the electron gas, coherence disappears at non zero temperature so that we can use a **classical description** of transport and thermodynamics

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Chaotic dynamical systems (large recurrence times, Feynman chain)

Vanishing of quantum coherence results from the existence of **incoherent** and **irreversible** processes associated to the coupling of electrons to their surrounding (additional degrees of freedom) :

Coupling to a bath of excitations : thermal excitations of the lattice (phonons)

Chaotic dynamical systems (large recurrence times, Feynman chain)

Impurities with internal degrees of freedom (magnetic impurities)

Electron-electron interactions,....



*The understanding of decoherence is difficult.*

It is a great challenge in quantum mesoscopic physics.

The phase coherence length  $L_\phi$  accounts in a generic way for decoherence processes.

The observation of coherent effects requires

$$L \ll L_\phi$$

*Averaging over disorder ?*



# Averaging over disorder ?

Expect to wash up interference effects



# Average coherence and multiple scattering

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Phase coherence leads to interference effects for a *given realization of disorder*.

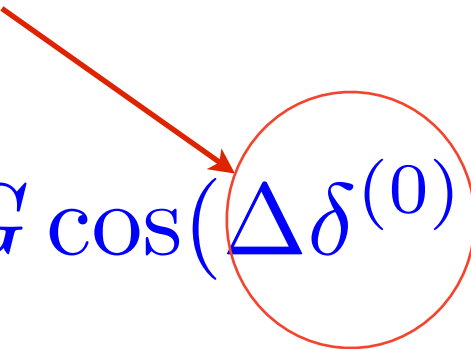
# Average coherence and multiple scattering

What is the role of **elastic** disorder ? Does it erase coherent effects ?

Phase coherence leads to interference effects for a *given realization of disorder*.

The *Webb* experiment corresponds to a fixed configuration of disorder.

Averaging over disorder  $\Rightarrow$  vanishing of the Aharonov-Bohm effect

$$G(\phi) = G_0 + \delta G \cos(\Delta\delta^{(0)} + 2\pi \frac{\phi}{\phi_0})$$


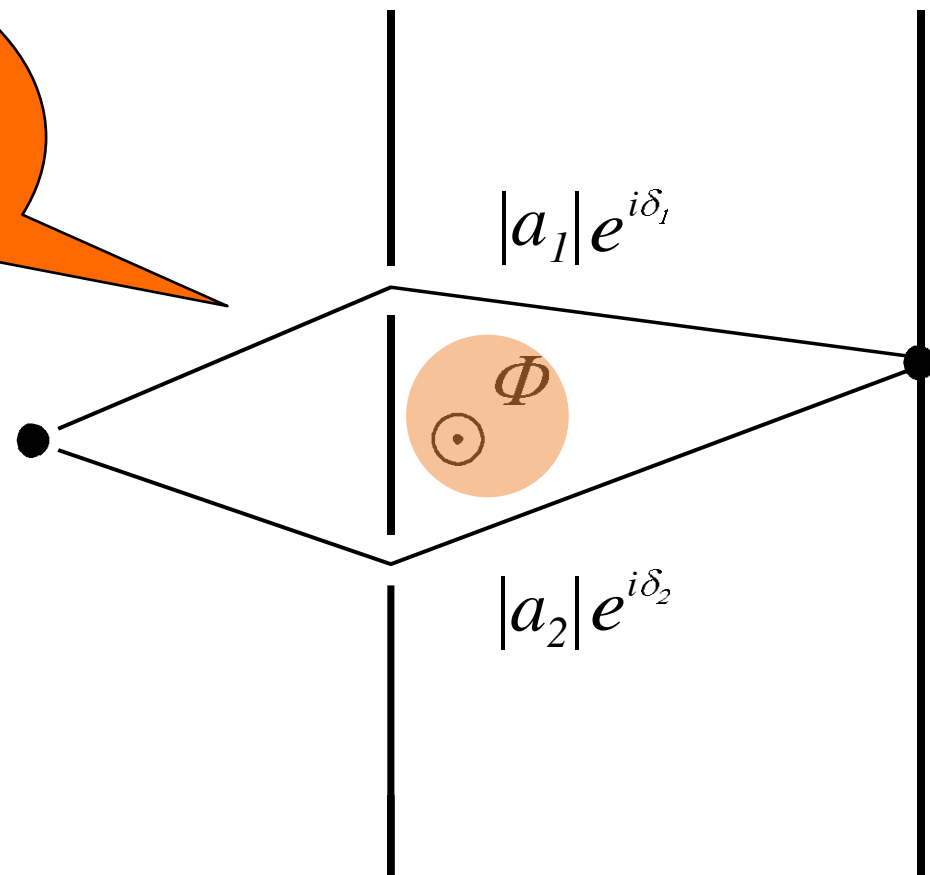
$$\Rightarrow \langle G(\phi) \rangle = G_0$$

Disorder seems to erase coherent effects....



# The setup : Young slits

A reminder

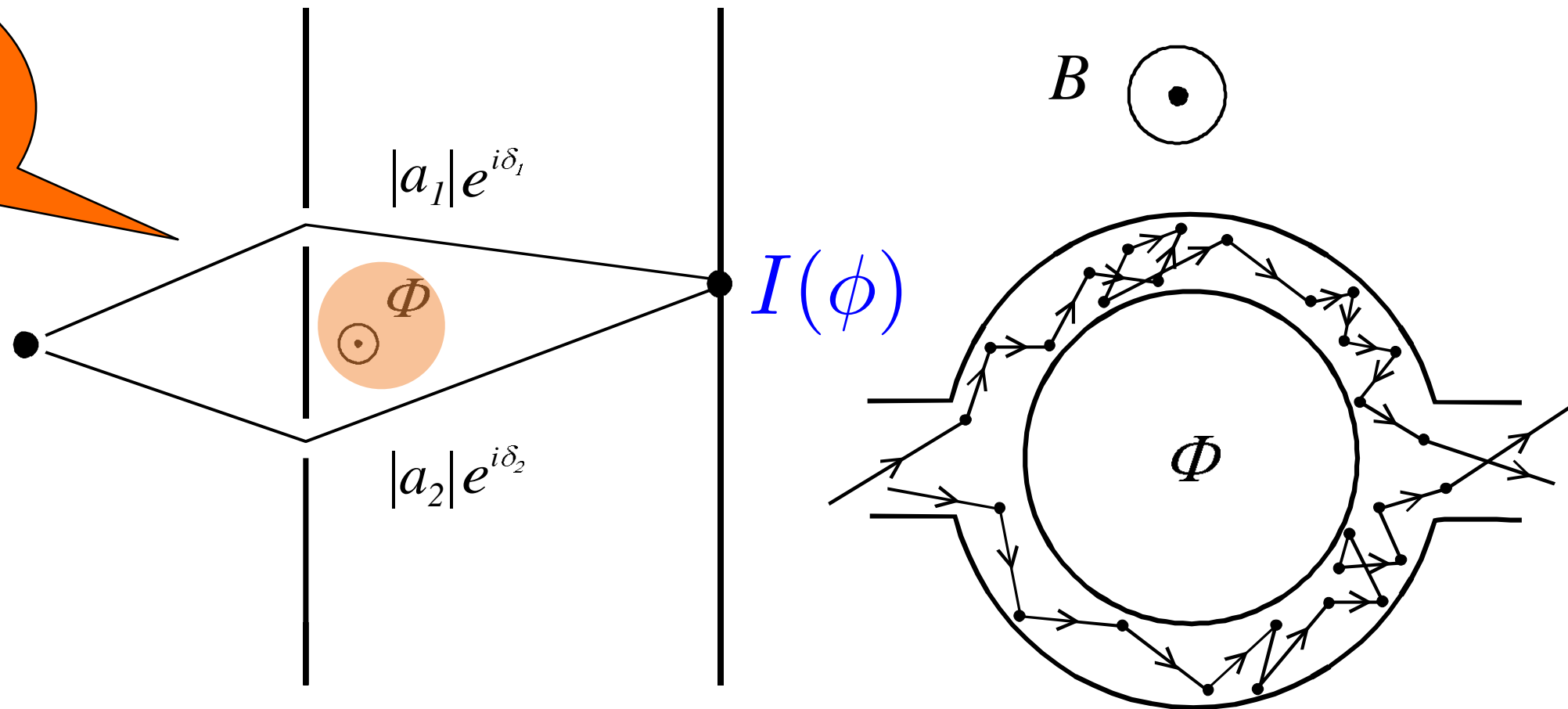


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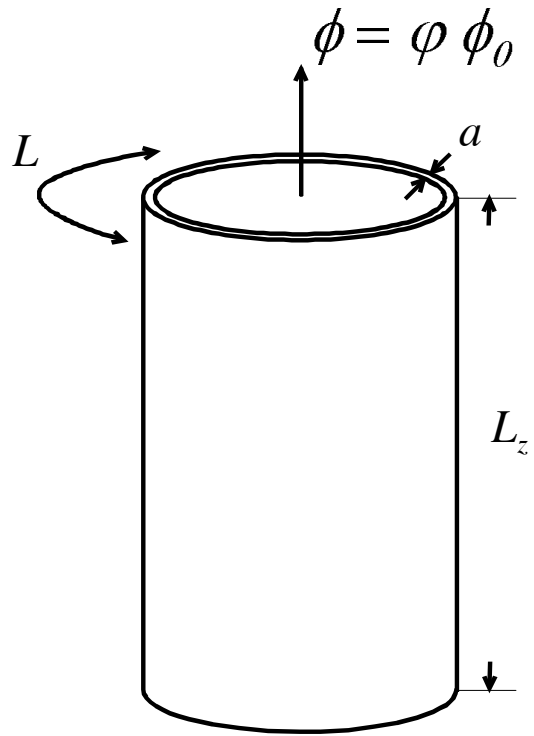
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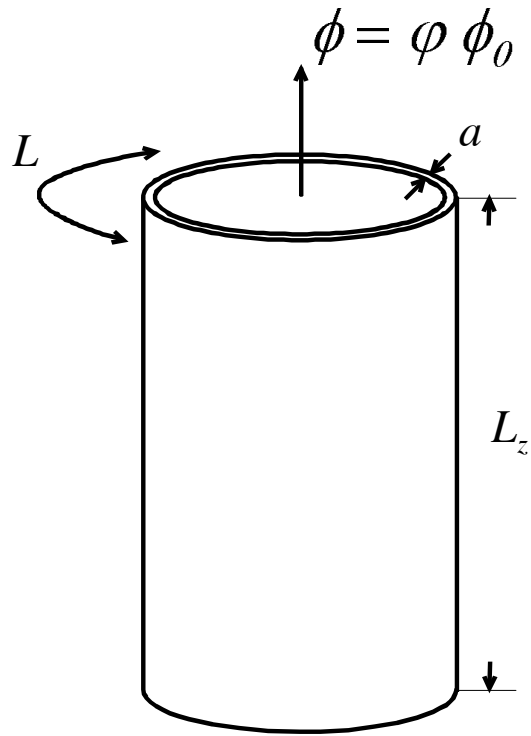
# The Sharvin<sup>2</sup> experiment



Experiment analogous to that of *Webb* but performed on a hollow cylinder of **height larger than  $L_\phi$**  pierced by a Aharonov-Bohm flux. **Ensemble of rings identical to those of *Webb* but incoherent between themselves.**



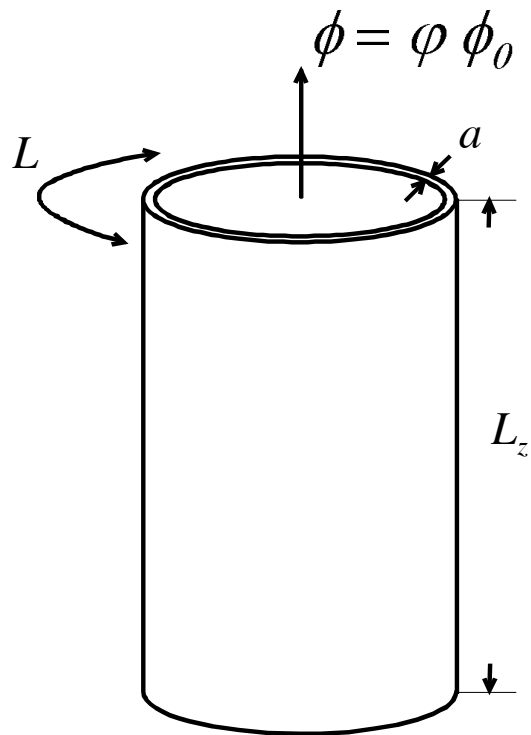
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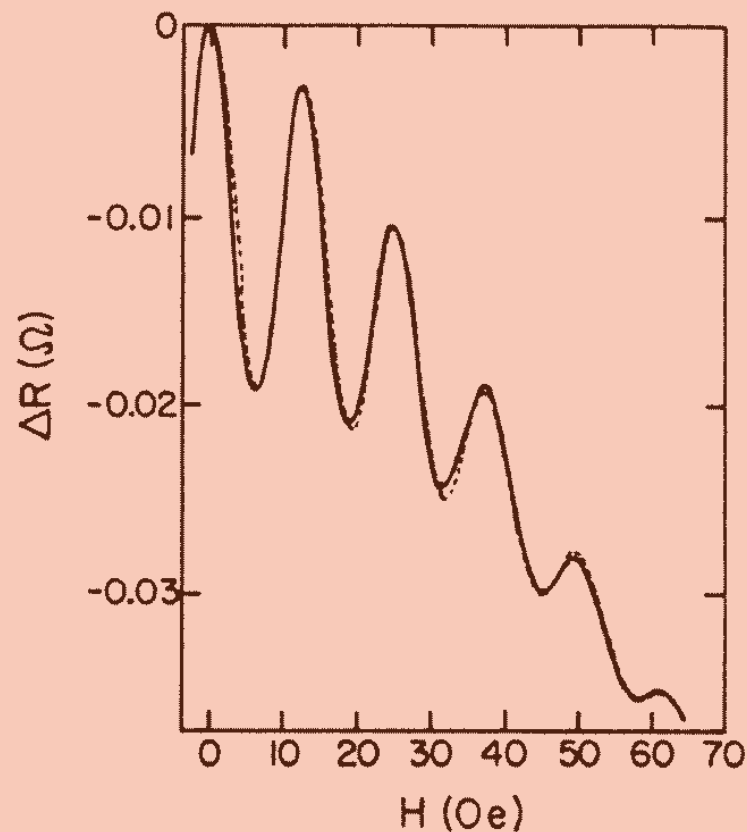
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The signal modulated at  $\phi_0$  *disappears*

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The signal modulated at  $\phi_0$  *disappears* but, instead, it appears a **new contribution** modulated at  $\phi_0/2$

**After all, disorder does not seem to erase coherent effects, but to modify them....**

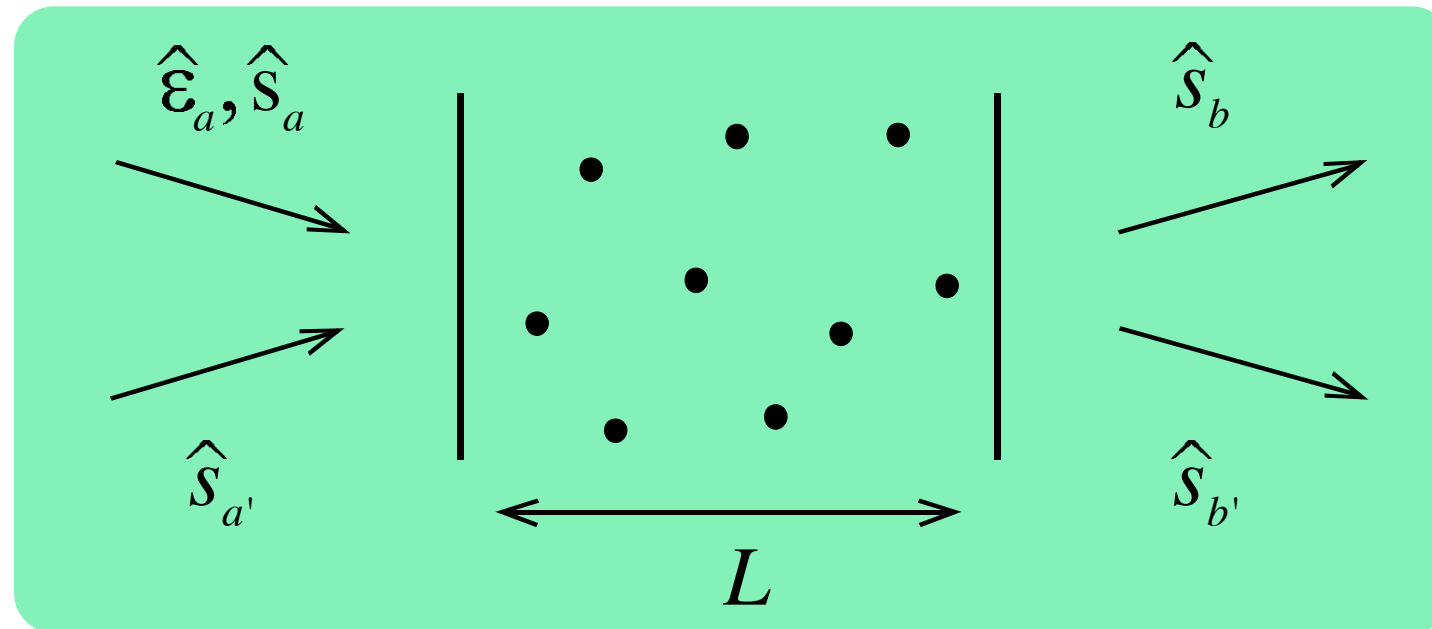
# Some “canonical” mesoscopic effects

Coherent backscattering in  
optics

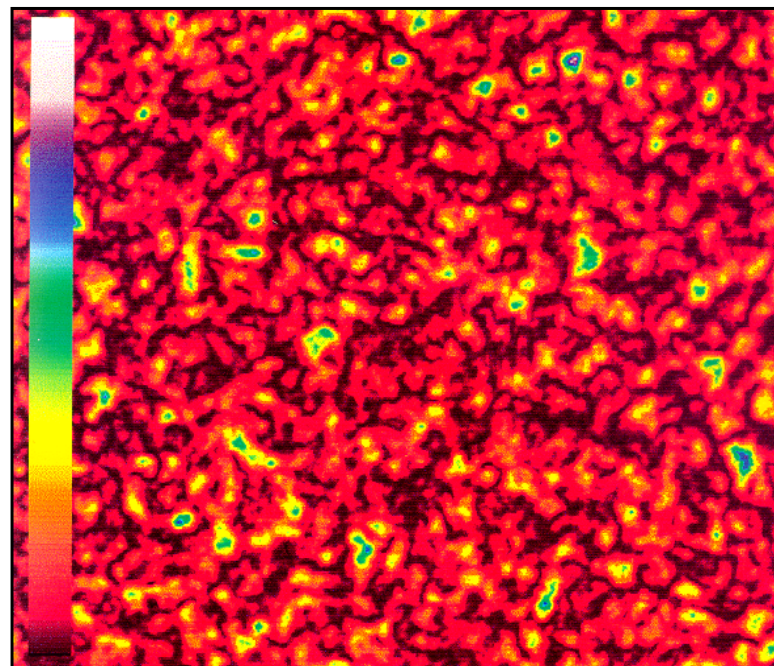


## An analogous problem: *Speckle patterns in optics*

Consider the elastic multiple scattering of light transmitted through a fixed disorder configuration.



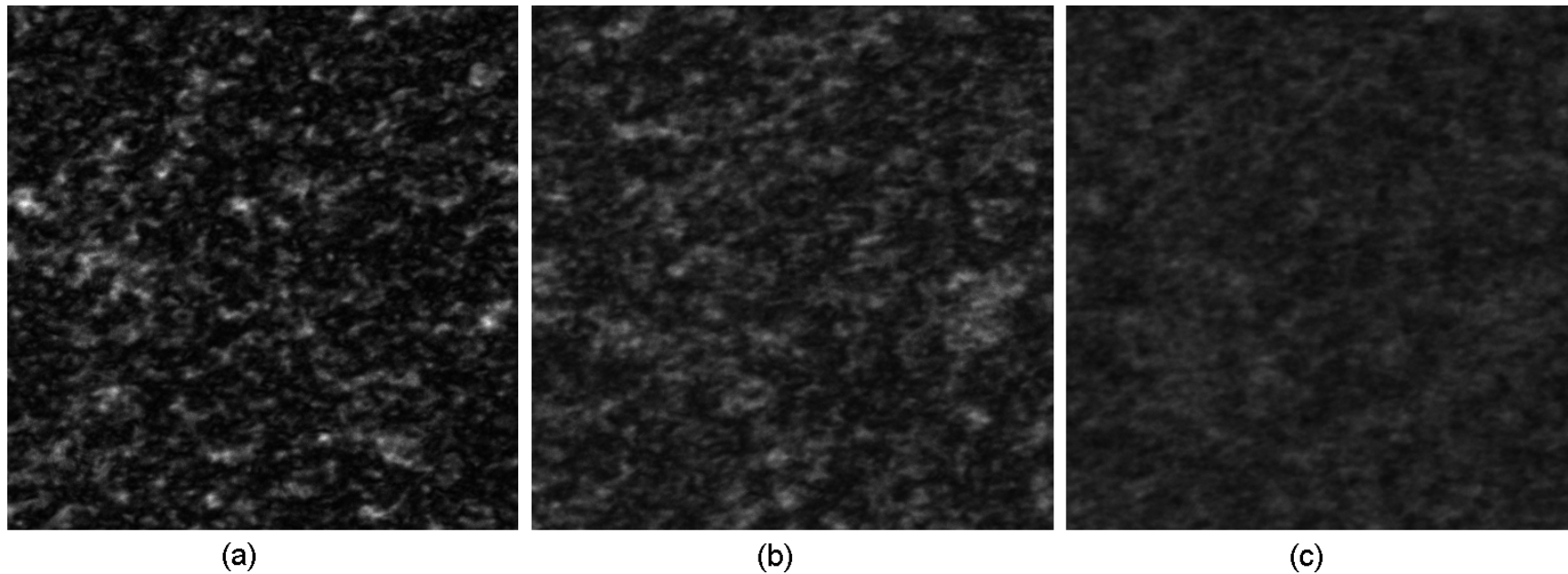
Outgoing light builds a **speckle pattern** *i.e.*, an **interference** picture:





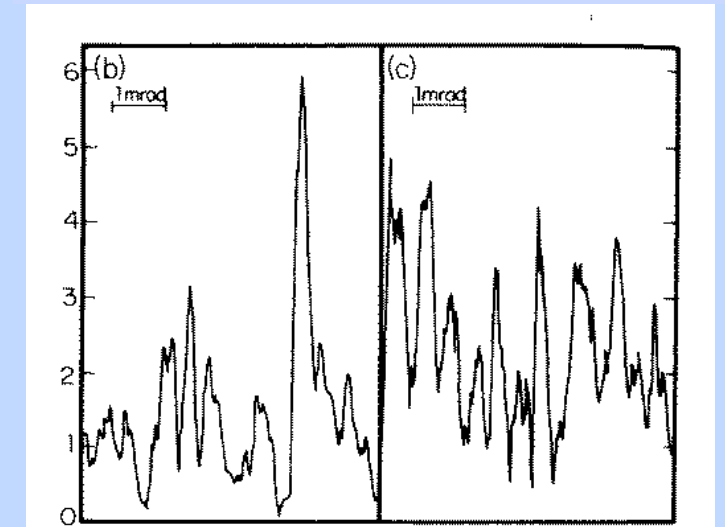
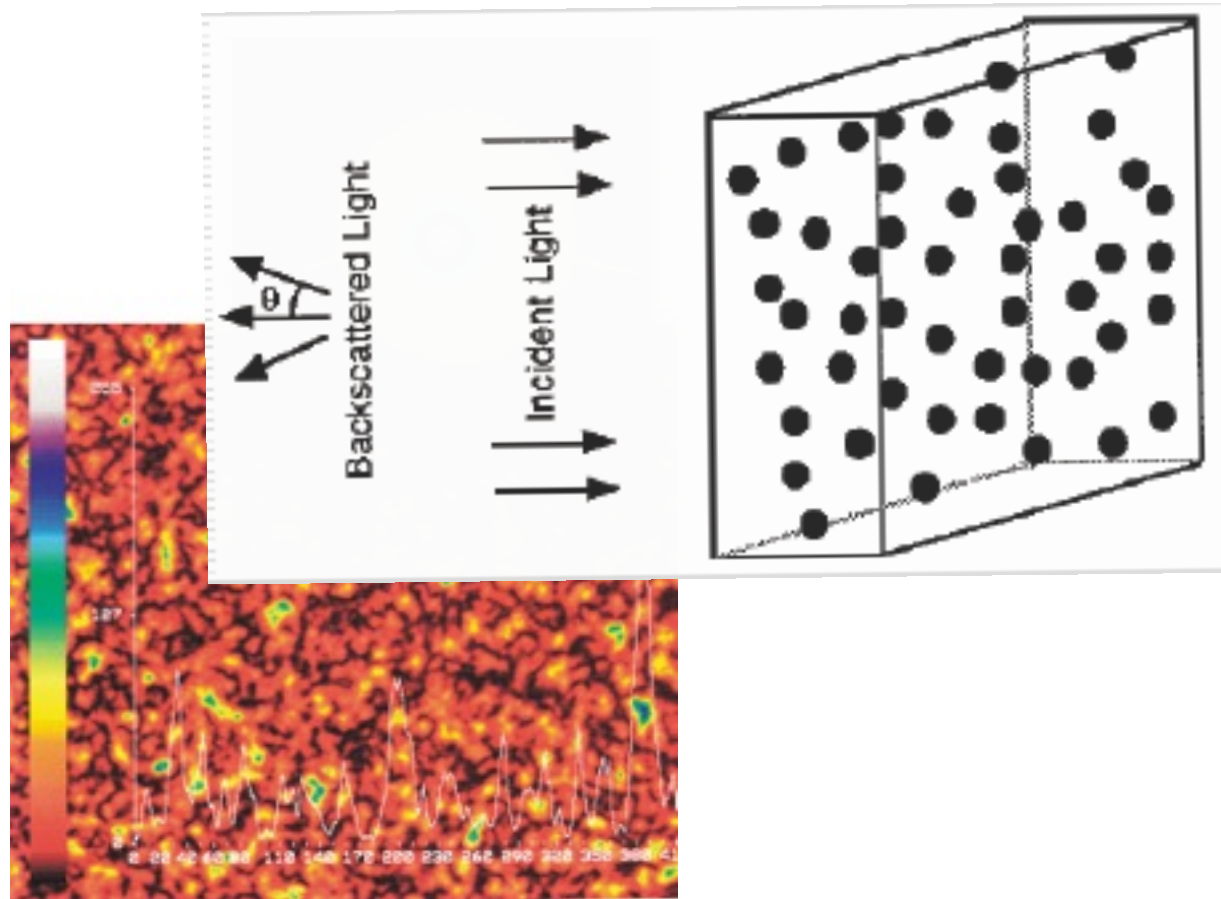
# Averaging over disorder erases the speckle pattern:

Integration over the motion of the scatterers leads to self-averaging

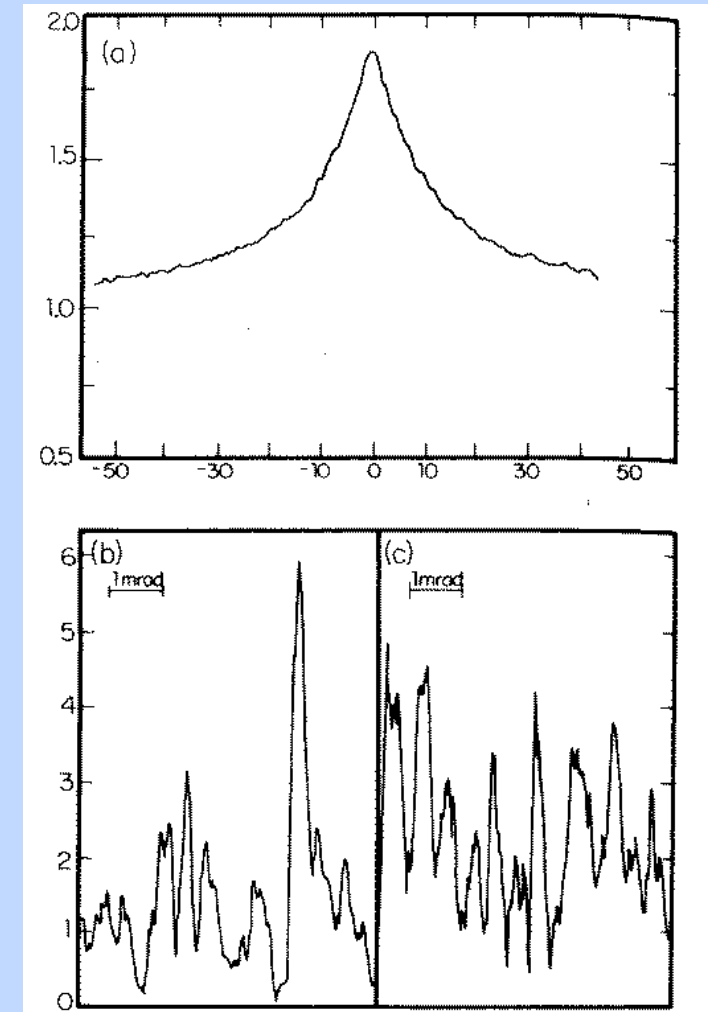
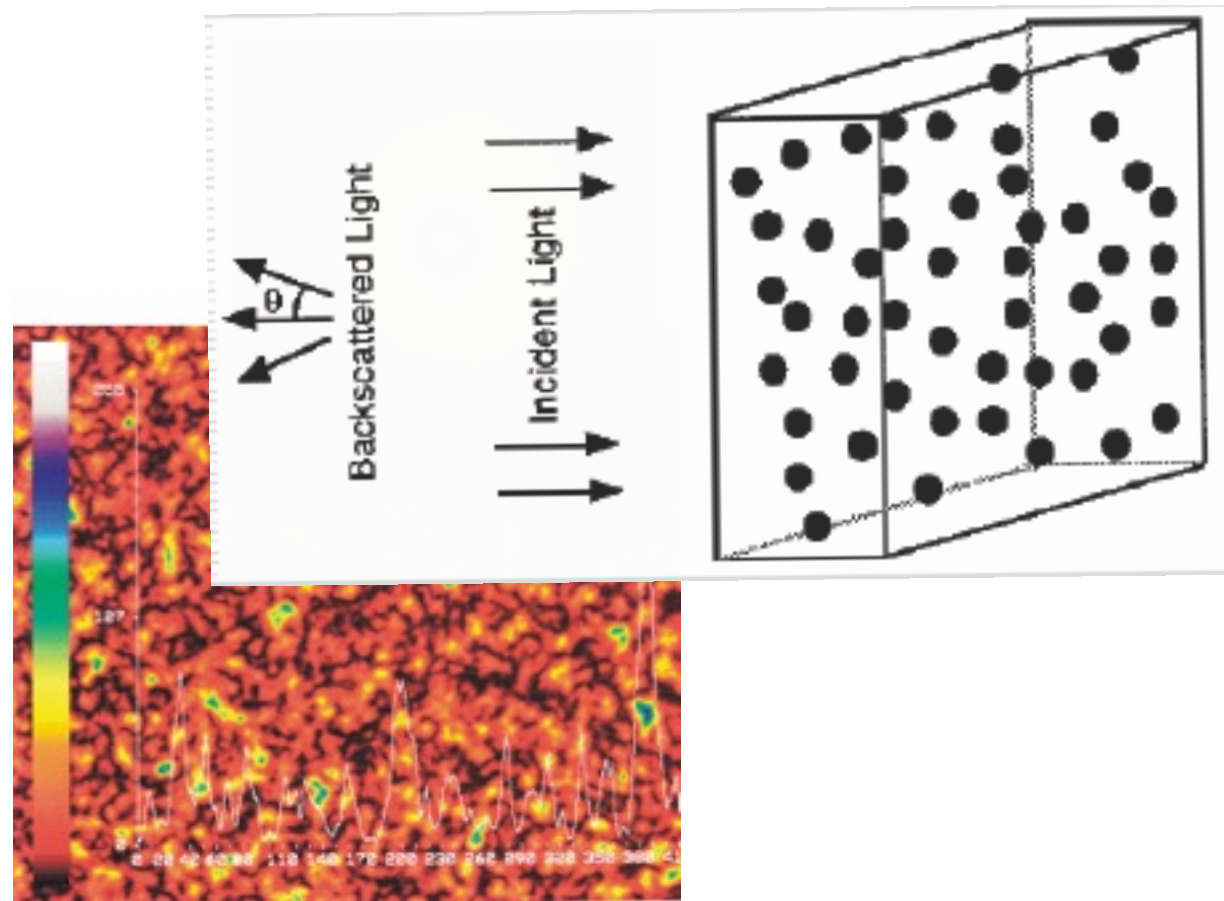


Time averaging

Does it erase all interferences ?



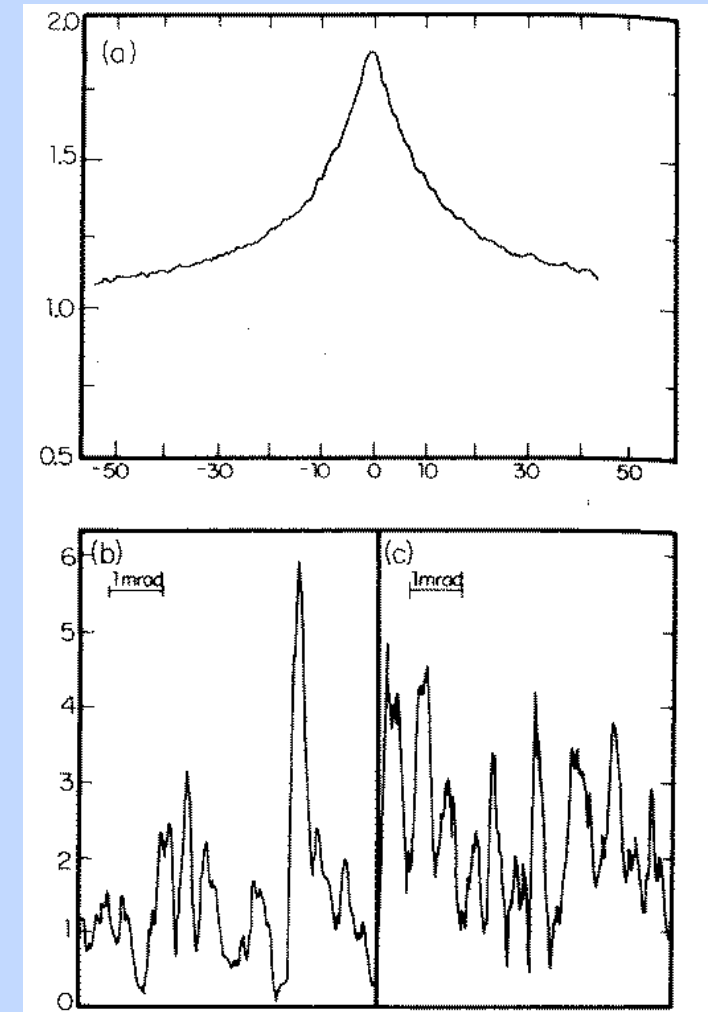
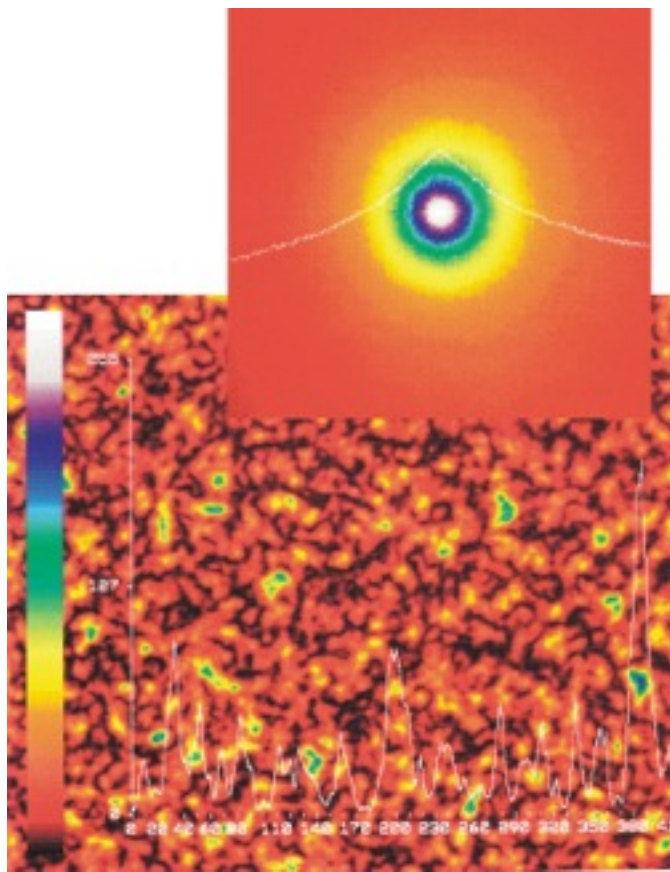
# Does it erase all interferences ?



Averaging over disorder does not produce incoherent intensity only, but also an angular dependent part, the coherent backscattering, which is a coherence effect. We may conclude:

Elastic disorder is **not related** to decoherence : **disorder does not destroy phase coherence and does not introduce irreversibility.**

# Does it erase all interferences ?

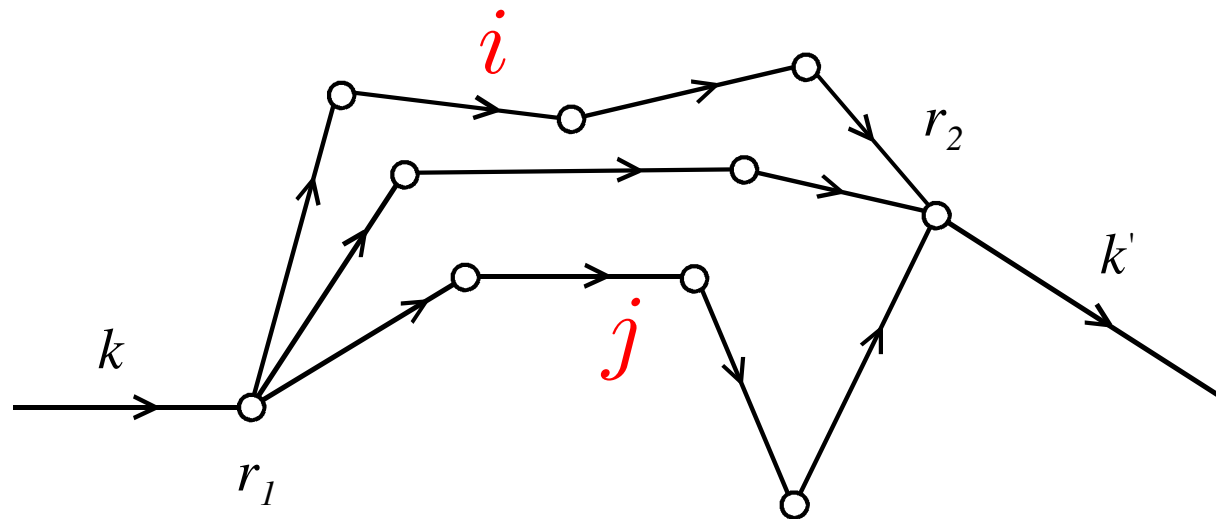


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Elastic disorder is **not related** to decoherence : **disorder does not destroy phase coherence and does not introduce irreversibility.**



# How to understand average coherent effects ?



Complex amplitude  $A(\mathbf{k}, \mathbf{k}')$  associated to the multiple scattering of a wave (electron or photon) incident with a wave vector  $\mathbf{k}$  and outgoing with  $\mathbf{k}'$

$$A(\mathbf{k}, \mathbf{k}') = \sum_{\mathbf{r}_1, \mathbf{r}_2} f(\mathbf{r}_1, \mathbf{r}_2) e^{i(\mathbf{k} \cdot \mathbf{r}_1 - \mathbf{k}' \cdot \mathbf{r}_2)}$$

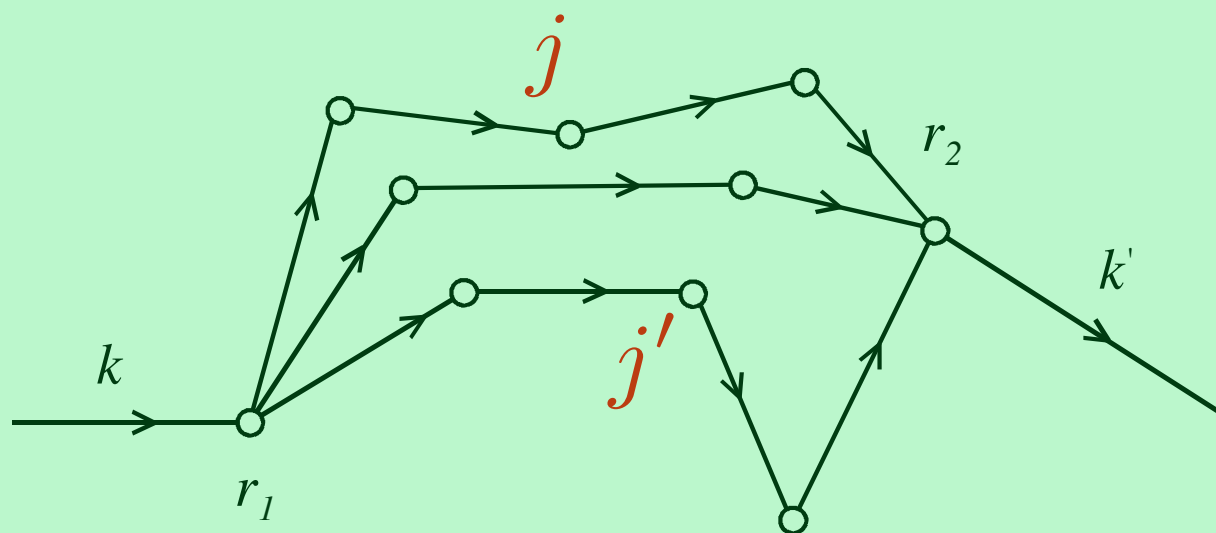
the complex amplitude  $f(\mathbf{r}_1, \mathbf{r}_2) = \sum_j |a_j| e^{i\delta_j}$  describes the propagation of the wave between  $\mathbf{r}_1$  and  $\mathbf{r}_2$ .

The corresponding intensity is

$$|A(\mathbf{k}, \mathbf{k}')|^2 = \sum_{\mathbf{r}_1, \mathbf{r}_2} \sum_{\mathbf{r}_3, \mathbf{r}_4} f(\mathbf{r}_1, \mathbf{r}_2) f^*(\mathbf{r}_3, \mathbf{r}_4) e^{i(\mathbf{k} \cdot \mathbf{r}_1 - \mathbf{k}' \cdot \mathbf{r}_2)} e^{-i(\mathbf{k} \cdot \mathbf{r}_3 - \mathbf{k}' \cdot \mathbf{r}_4)}$$

with

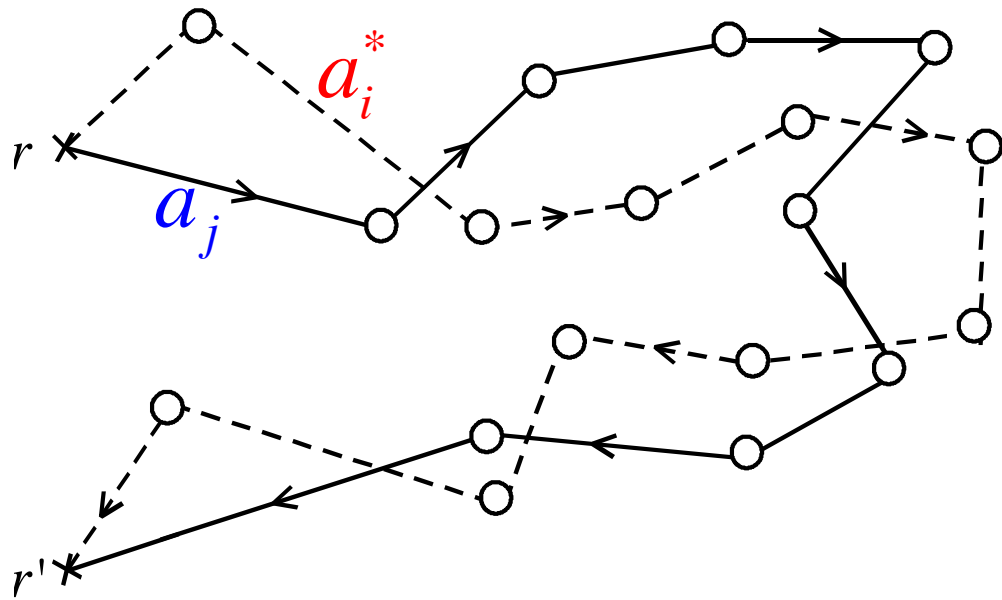
$$f(\mathbf{r}_1, \mathbf{r}_2) f^*(\mathbf{r}_3, \mathbf{r}_4) = \sum_{j, j'} a_j(\mathbf{r}_1, \mathbf{r}_2) a_{j'}^*(\mathbf{r}_3, \mathbf{r}_4) = \sum_{j, j'} |a_j| |a_{j'}| e^{i(\delta_j - \delta_{j'})}$$



On average over disorder, most contributions to  $f f^*$  disappear since the dephasing  $\delta_j - \delta_{j'} \gg 1$

The only remaining contributions to the intensity correspond to terms with **zero dephasing**, *i.e.*, to **identical trajectories**.

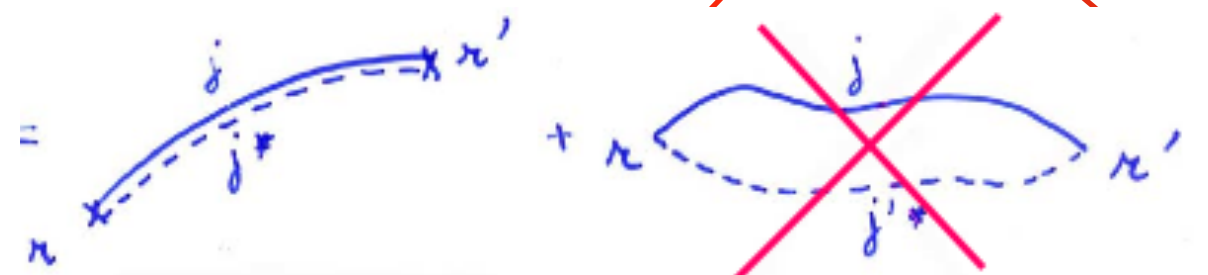
# *Quantum probability* for propagation between two points



$$P(r, r') = \sum_{i,j} \overline{a_i^*(r, r') a_j(r, r')}$$

$$P(r, r') = \sum_j \overline{|a_j(r, r')|^2} + \sum_{i \neq j} \overline{a_i^*(r, r') a_j(r, r')}$$

$$a_i^* a_j = |a_i| |a_j| e^{i(\delta_i - \delta_j)}$$

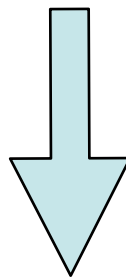
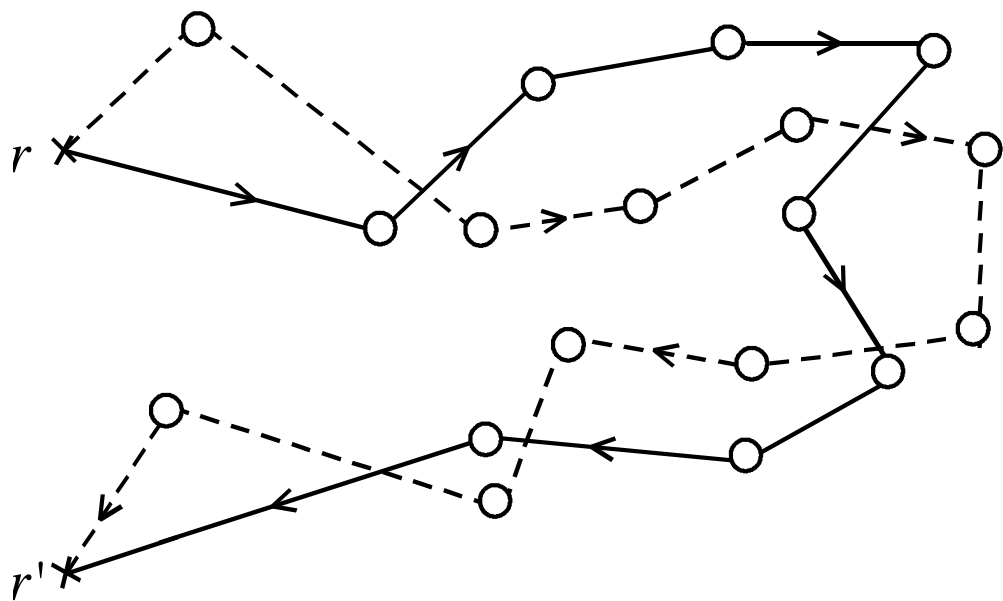


$$\delta_i - \delta_j \gg 1$$

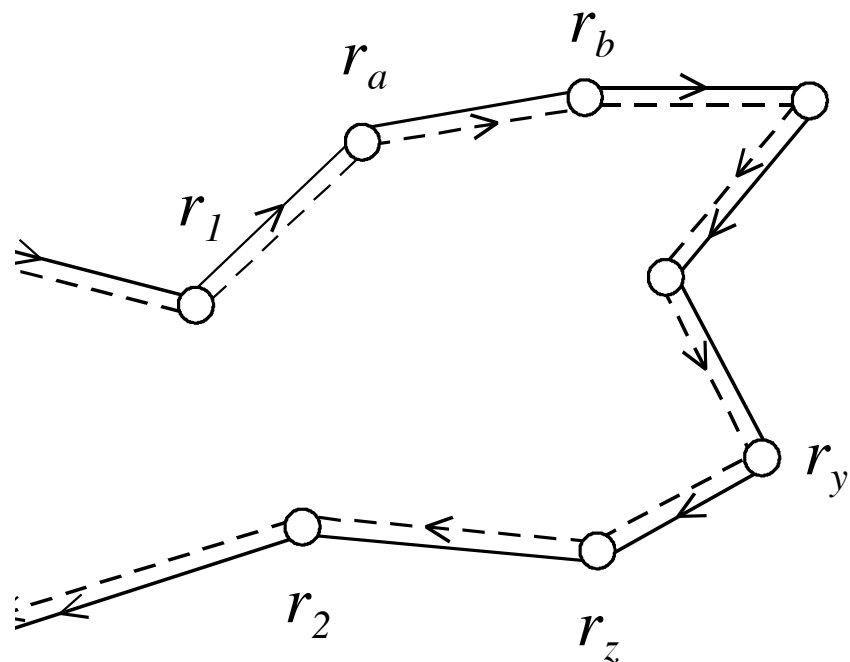
*Incoherent propagation !*

vanishes on average

# Some useful design

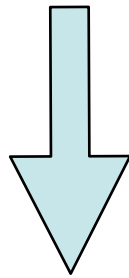
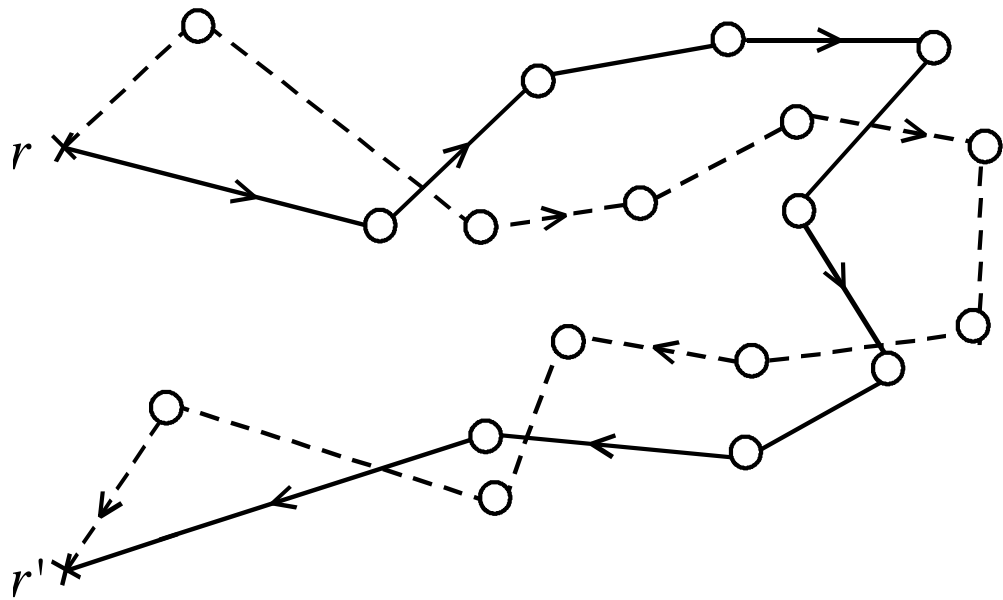


Disorder  
averaging

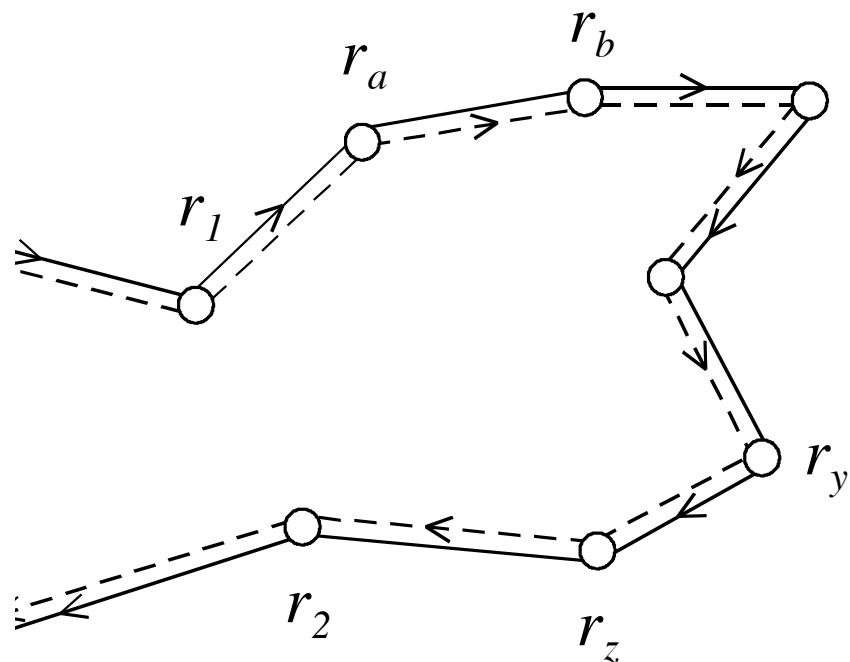




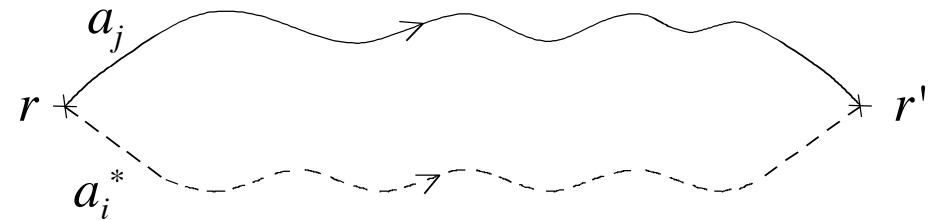
# Some useful design



Disorder  
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(a)



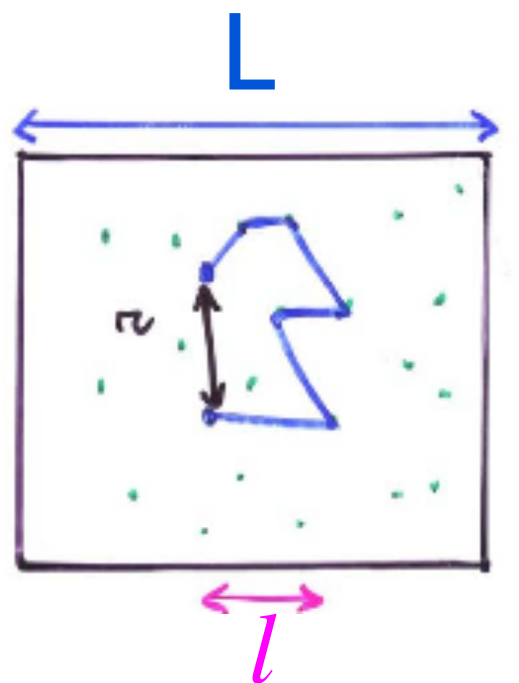
(b)



To a good approximation, the **incoherent contribution** obeys a classical **diffusion equation**

$$\left( \frac{\partial}{\partial t} - D\Delta \right) P(r, r', t) = \delta(r - r')\delta(t) \quad \Leftrightarrow \quad (-i\omega + Dq^2) P(q, \omega) = 1$$

*Incoherent electrons diffuse* in the conductor with a *diffusion coefficient  $D$*  (*Drude theory*)

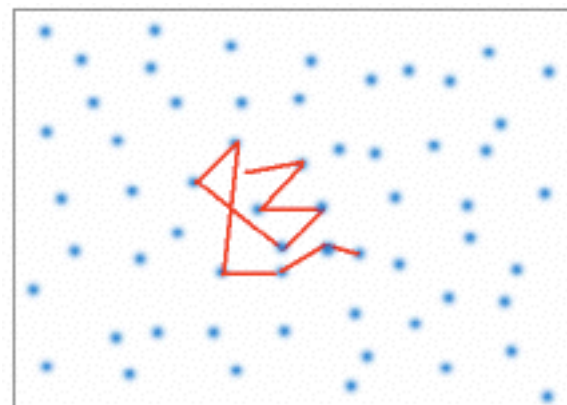


$$l \ll L$$

$$\langle r^2 \rangle = 2d Dt$$

space dimensionality

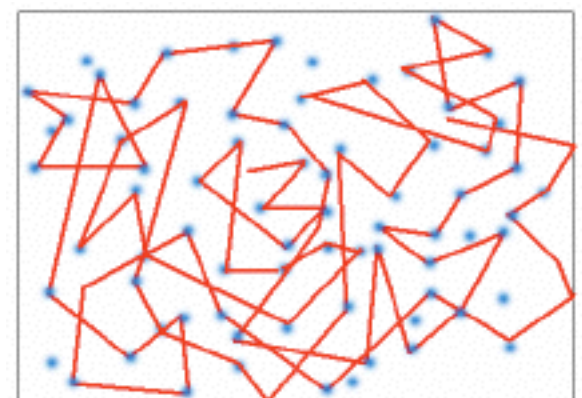
$$t \ll \tau_D$$



Thouless time

$$L^2 = D\tau_D$$

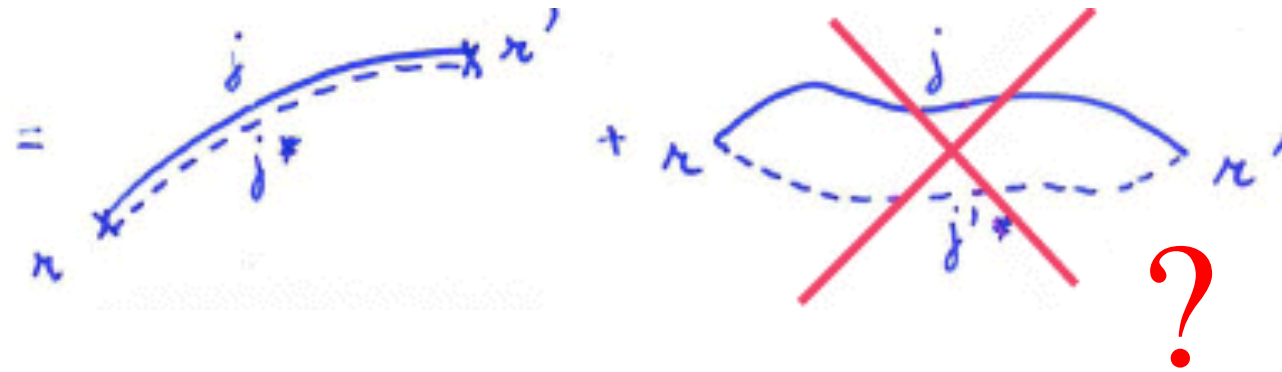
$$t \gg \tau_D$$



# Beyond incoherent diffusion (qualitative description)



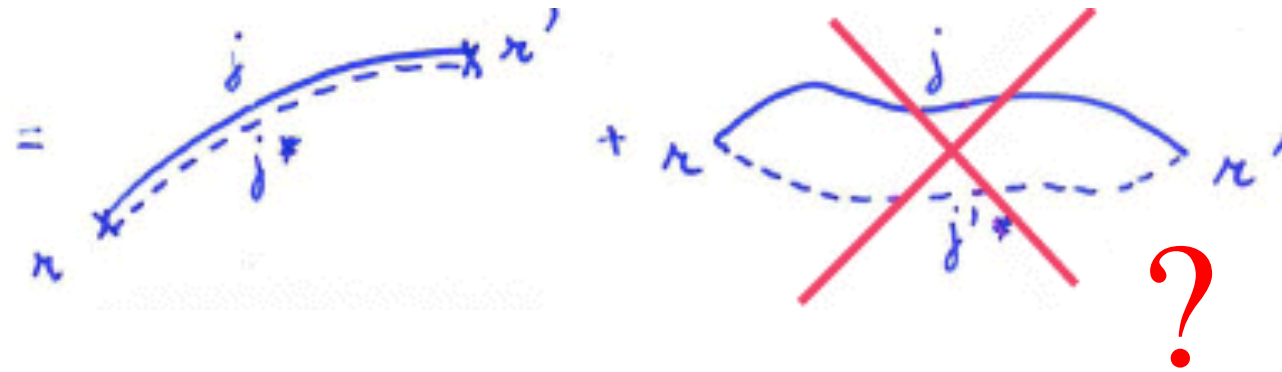
## Coherent effects



What is the first correction *i.e.*, with the *smallest phase shift* ?



# Coherent effects

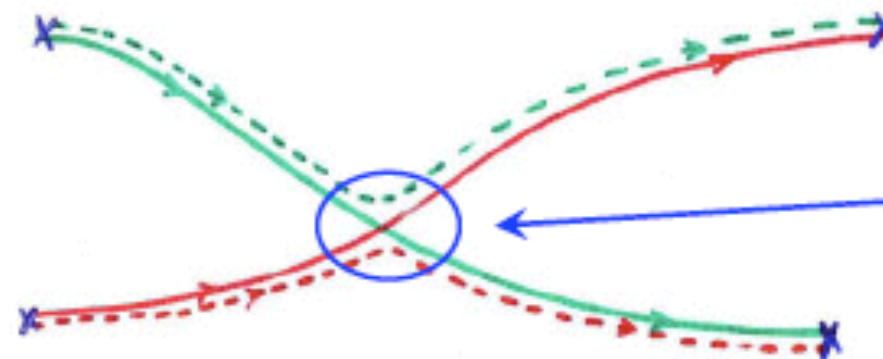


What is the first correction *i.e.*, with the  
*smallest phase shift* ?  
When amplitude paths cross

Example :



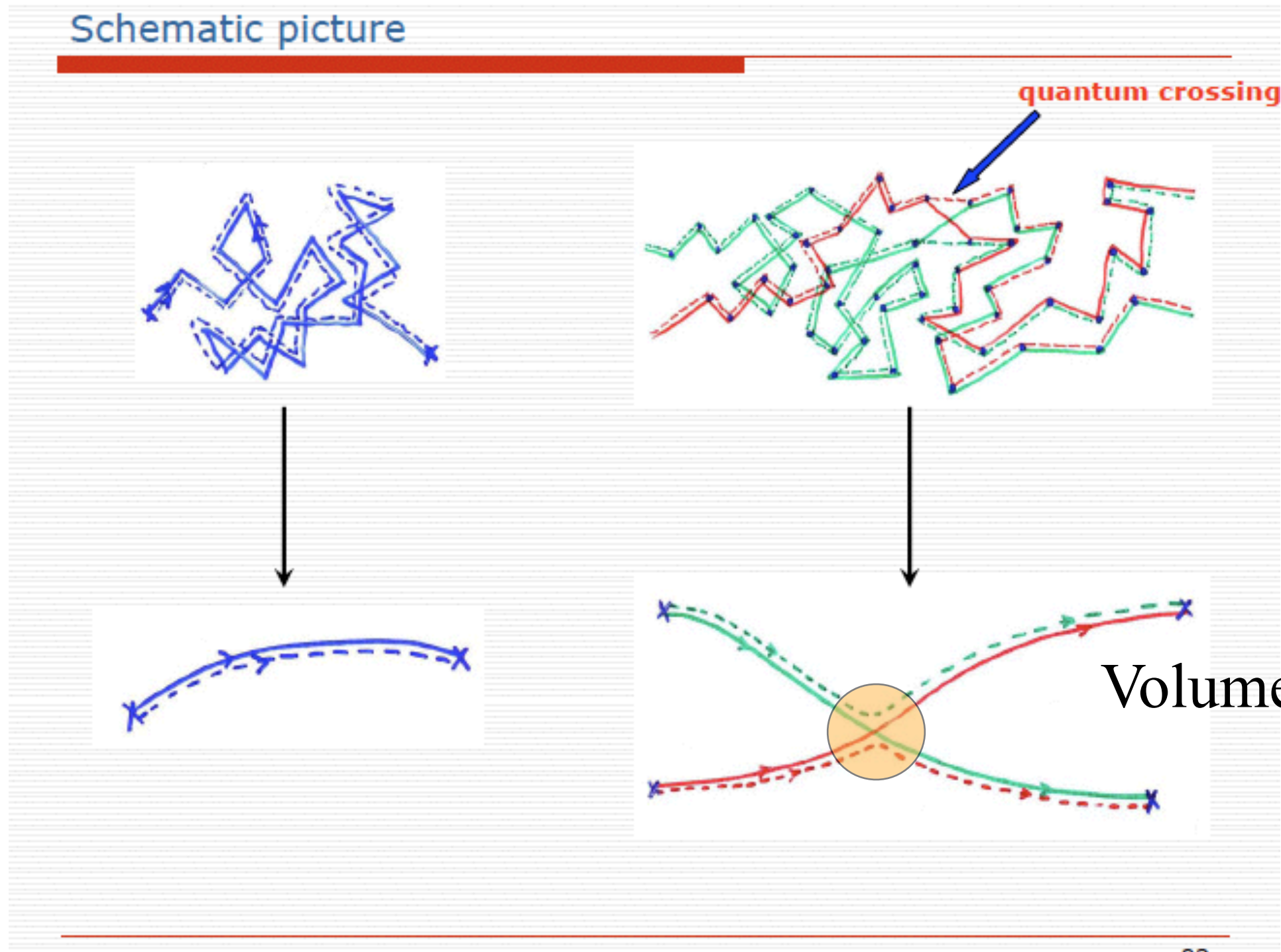
Classical diffusion



quantum  
crossing

Exchange of amplitudes

## Schematic picture



Quantum crossings decrease the diffusion coefficient  $D$  :  
weak localization

$\lambda$  : Fermi wavelength

Occurrence of a *quantum crossing* after a time  $t$  for a electron diffusing in a volume  $L^d$

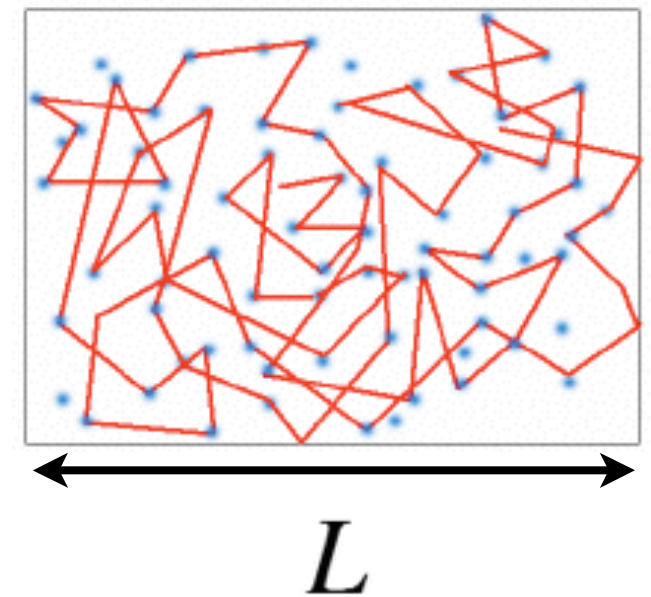
$$p_{\times}(t) = \frac{\lambda_F^{d-1} v_F t}{L^d}$$

$v_F$  : Fermi velocity

The time spent by a diffusing electron is  $\tau_D = L^2/D$  so that

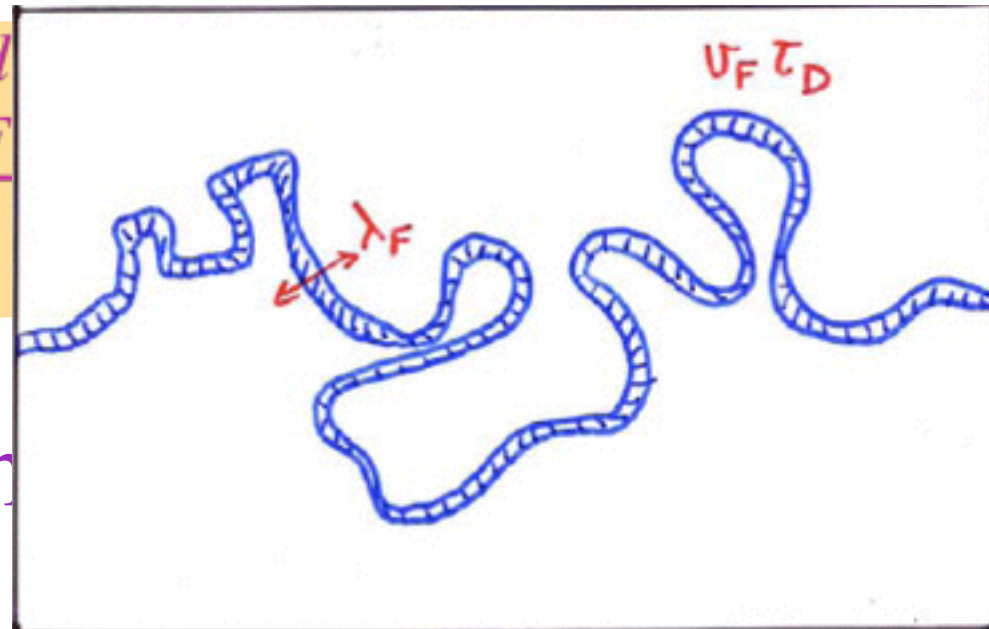
$$p_{\times}(\tau_D) = \frac{\lambda_F^{d-1} v_F \tau_D}{L^d} \equiv \frac{1}{g}$$

$$g = \frac{D}{v_F \lambda_F^{d-1}} L^{d-2}$$



Occurrence of a *quantum crossing* after a time  $t$  for a electron diffusing in a volume  $L^d$

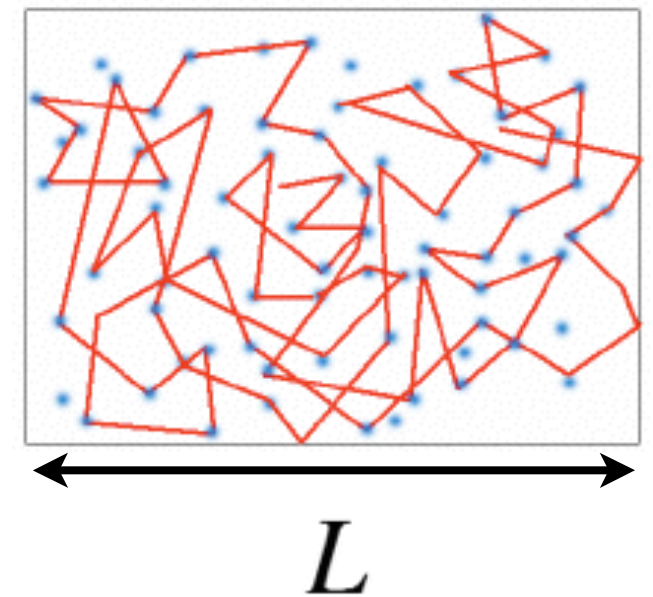
$$p_{\times}(t) = \frac{\lambda_F^d}{V}$$



mi velocity

The time spent by a diffusing electron in a volume  $L^d$  is  $t \sim L^d / (v_F \lambda_F^d)$  so that

$$g \sim \frac{V}{\lambda_F^{d-1} v_F \tau_D}$$





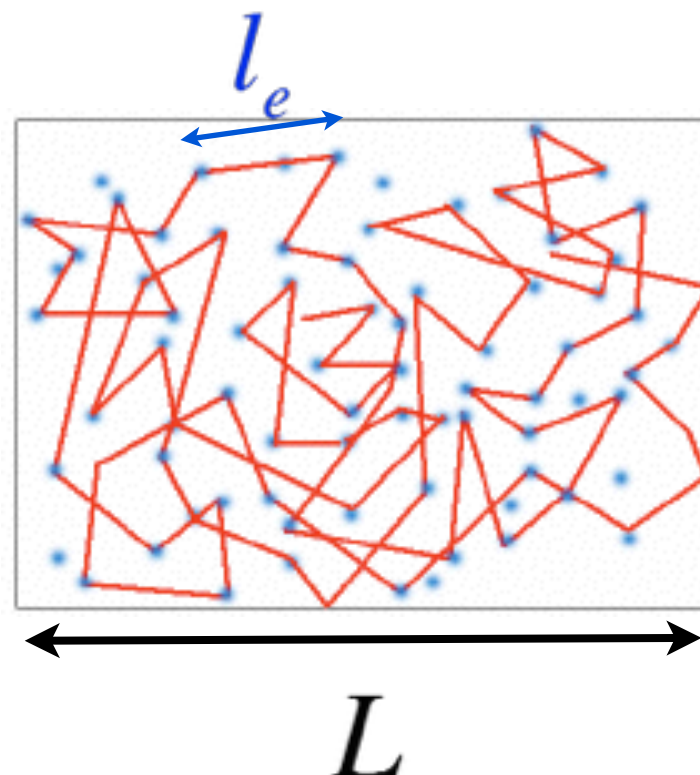
$$g = \frac{D}{v_F \lambda_F^{d-1}} L^{d-2}$$

*Physical meaning of this parameter ?*

# Electrical conductance of a metal

A metal can be modeled as a quantum gas of electrons scattered by an elastic disorder  $l_e$ .

Classically, the conductance of a cubic sample of size  $L^d$  is given by Ohm's law:  $G = \sigma L^{d-2}$  where  $\sigma$  is the conductivity.



## Conductance = ratio of two volumes

---

Conductivity (Drude)  $\sigma = \frac{ne^2\tau_e}{m} = e^2 D \rho$   $D = \frac{v_F l_e}{3}$

Conductance (Ohm's law)  $G = \sigma \frac{S}{L}$   $S = W^{d-1}$

Dimensionless conductance  $G = g \frac{e^2}{h}$

$$g \sim \frac{D}{v_F \lambda^{d-1}} \frac{W^{d-1}}{L} = \frac{D}{v_F \lambda^{d-1}} \frac{V}{L^2} \sim \frac{V}{\lambda^{d-1} v_F \tau_D}$$

A direct consequence: quantum corrections to electrical transport

Classical transport :  $G_{cl} = g \times \frac{e^2}{h}$  with  $g \gg 1$



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What does it mean ?

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# A direct consequence: quantum corrections to electrical transport

Classical transport :  $G_{cl} = g \times \frac{e^2}{h}$  with  $g \gg 1$

Quantum

Independent of the microscopic  
(and often unknown) disorder -  
Depends only on the geometry

so that  $\Delta G \simeq \frac{e^2}{h}$  is universal



# A direct consequence: quantum corrections to electrical transport

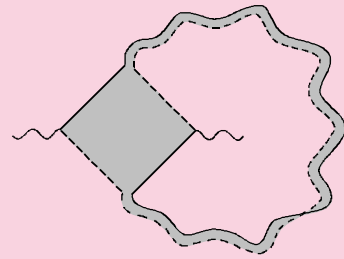
Classical  $G_{cl} = g \times \frac{e^2}{h}$  with  $g \gg 1$

Not that simple !  
numbers... Need to sum up Feynman diagrams.

Quantum corrections

so that  $\Delta G = \# \frac{e^2}{h}$  is universal

# A direct consequence: quantum corrections to electrical



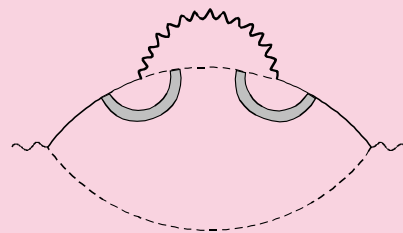
$$\text{Shaded Diamond} = \text{Diamond with Left Triangle Shaded} + \text{Diamond with Right Triangle Shaded} + \text{Diamond with Diagonal Line} + \text{Diamond with Diagonal Line}$$

$$\Delta\sigma^* = \text{Diagram 1} + \text{Diagram 2} + \text{Diagram 3}$$

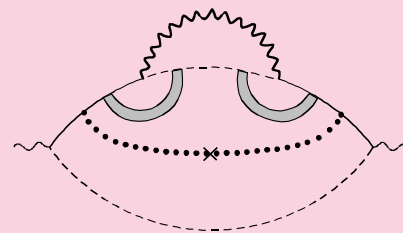
Diagram 1: A bubble diagram with a shaded vertical rectangle labeled  $\Gamma'$  inside. Diagram 2: A bubble diagram with a shaded vertical rectangle labeled  $\Gamma'$  inside, and a dashed line with a cross at the top. Diagram 3: A bubble diagram with a shaded vertical rectangle labeled  $\Gamma'$  inside, and a dashed line with a cross at the bottom.

c)

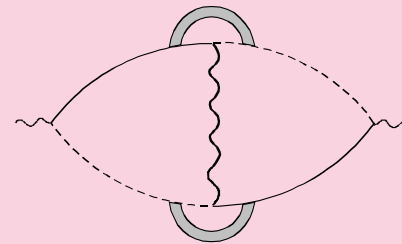
$$\text{Triangle with wavy line} = \text{Wavy line with dashed line and cross}$$



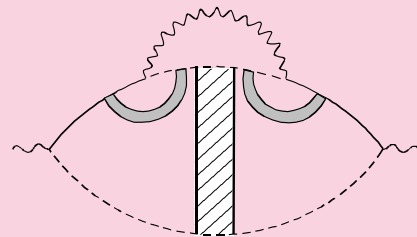
(a)



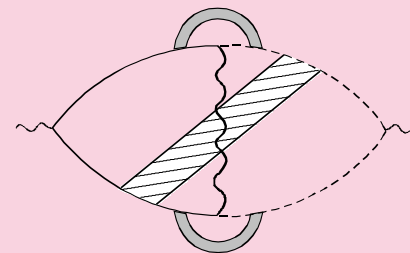
(b)



(c)



(d)

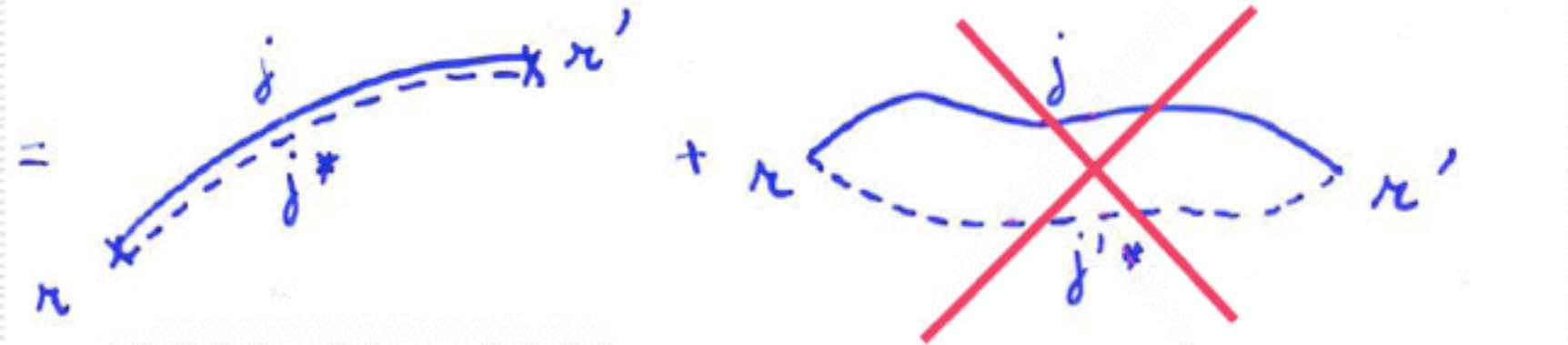


(e)

# Summary

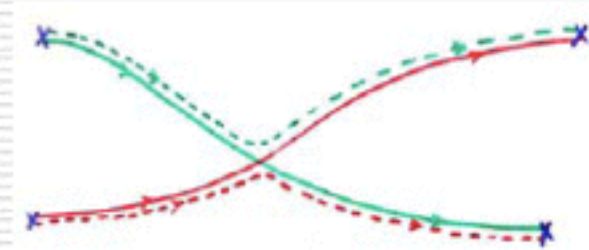
conductance  $\sim$  transmission  $\sim$  probability

$$P_{cl}(r, r')$$



classical diffusion

quantum corrections



quantum crossing  $\rightarrow$   $1/g$  correction

$$p_x(\tau_D) \sim \frac{\lambda_F^{d-1} v_F \tau_D}{V} \sim \frac{1}{g}$$

classical transport  $\propto g \frac{e^2}{h}$

quantum effects  $\propto \frac{e^2}{h}$



# Complexity of a quantum mesoscopic system



Elastic disorder does not break phase coherence  
and it does not introduce irreversibility



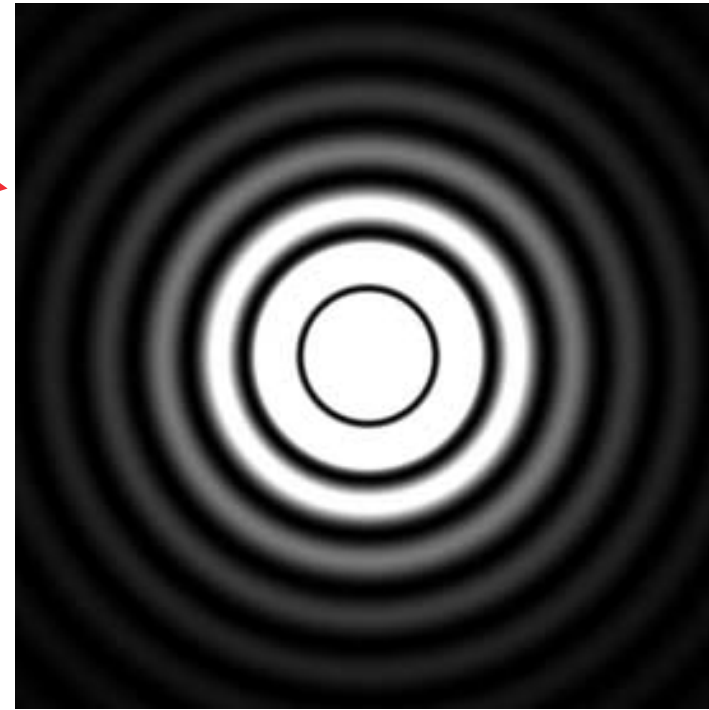
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Disorder introduces randomness and  
complexity:

All symmetries are lost, there are no good  
quantum numbers.

## Example: speckle patterns in optics

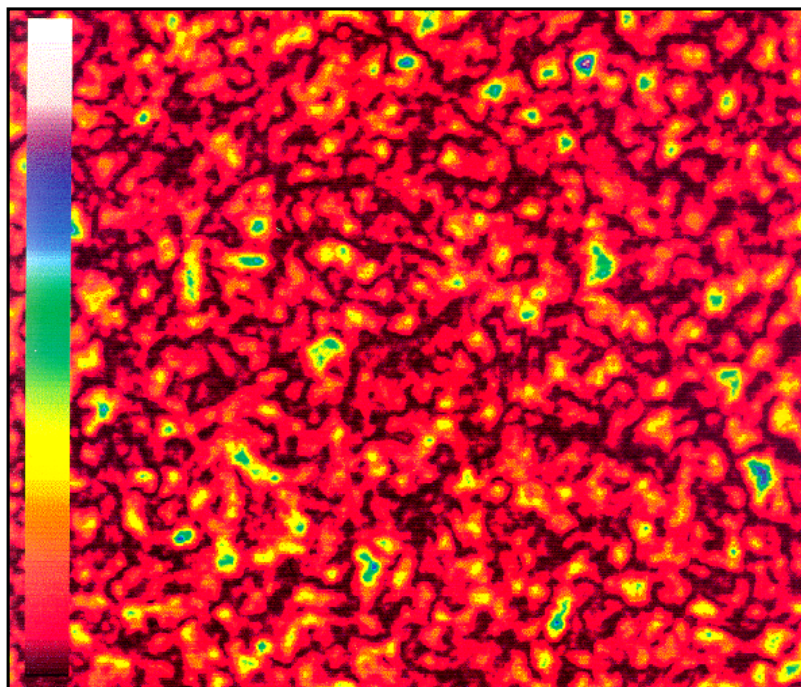
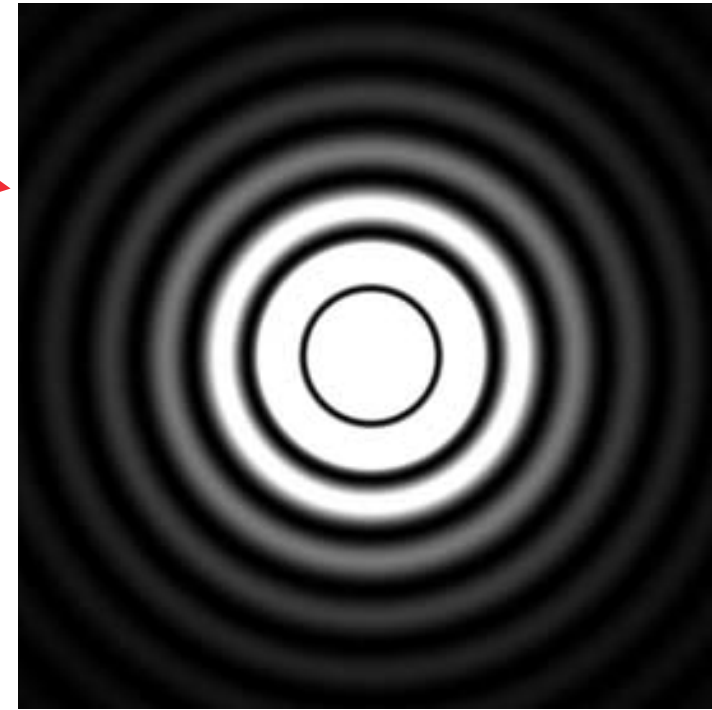
Diffraction  
through a circular  
aperture: order in  
interference





## Example: speckle patterns in optics

Diffraction  
through a circular  
aperture: order in  
interference



Transmission of  
light through a  
disordered  
suspension:  
complex system



# Mesoscopic quantum systems

- Most (all ?) quantum systems are complex
- Complexity (randomness) and decoherence are separate and independent notions.
- Complexity  
quantum numbers)
- Decoherence  
coherence  $L \gg L_\varphi$

A mesoscopic quantum system is a coherent complex quantum system with  $L \leq L_\varphi$

# Phase coherence and self-averaging: universal fluctuations.

Classical limit :  $L \gg L_\varphi$

The system is a collection of  $N = (L/L_\varphi)^d \gg 1$   
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The system is a collection of  $N = (L/L_\varphi)^d \gg 1$  statistically independent subsystems.

A macroscopic observable defined in each subsystem takes independent random values in each of the  $N$  pieces.

**Law of large numbers:** any macroscopic observable is equal with probability one to its average value.



The system performs an  
average over realizations of  
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For  $L \ll L_\varphi$ , we expect deviations from self-averaging which reflect the underlying quantum coherence.

# Quantum conductance fluctuations

Classical self-averaging limit :  $\frac{\delta G}{\overline{G}} = \frac{1}{N} = \left(\frac{L_\varphi}{L}\right)^{d/2}$

where  $\delta G = \sqrt{\overline{G^2} - \overline{G}^2}$  and  $\overline{G} = \sigma L^{d-2}$

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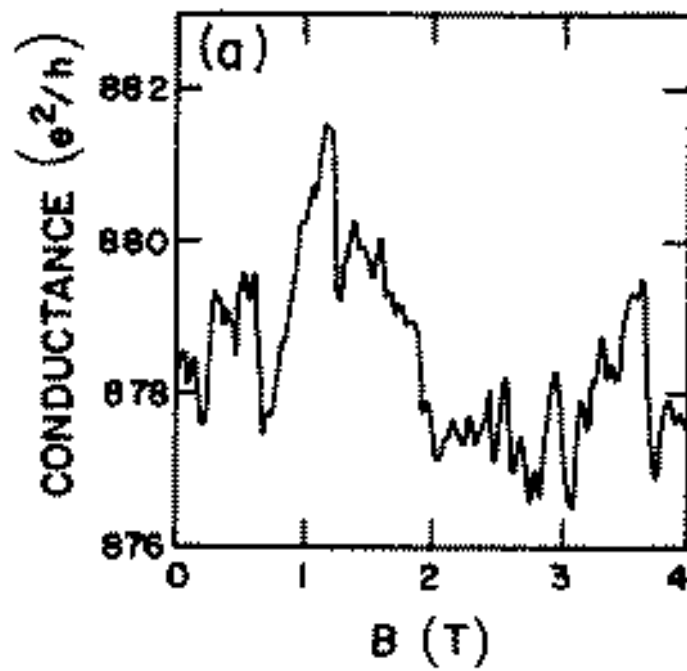
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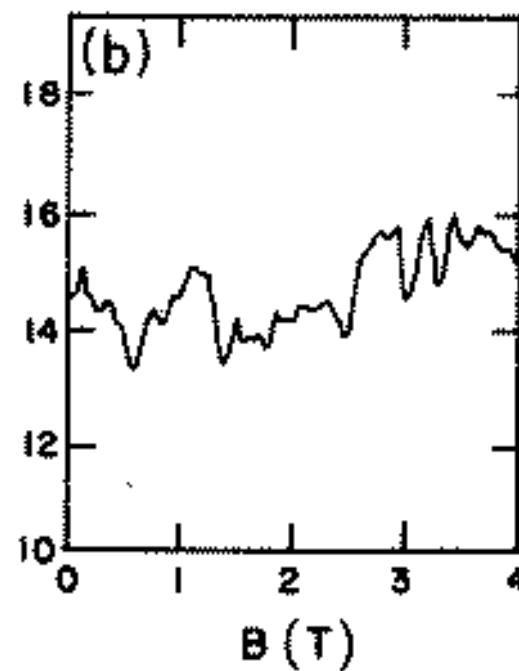
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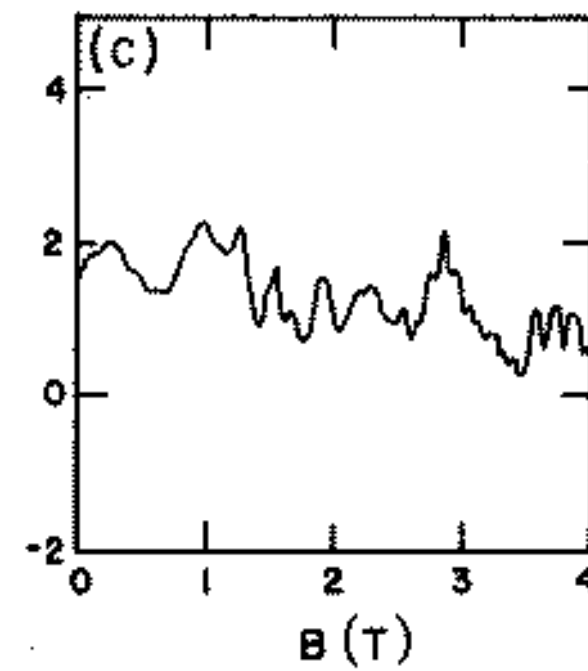
*In the mesoscopic limit, the electrical conductance is not self-averaging.*



Gold ring



Si-MOSFET



NUMERICS ON  
THE ANDERSON MODEL

# Summary : key ideas and concepts in quantum mesoscopic physics

1. Classical diffusion

2.

3.

crossings) are propagated over long distances by means of classical diffusion.

3. Complexity  
from decoherence.