

Physique mesoscopique des electrons et des photons -

Structures fractales et quasi-periodiques

ERIC AKKERMANS
PHYSICS-TECHNION



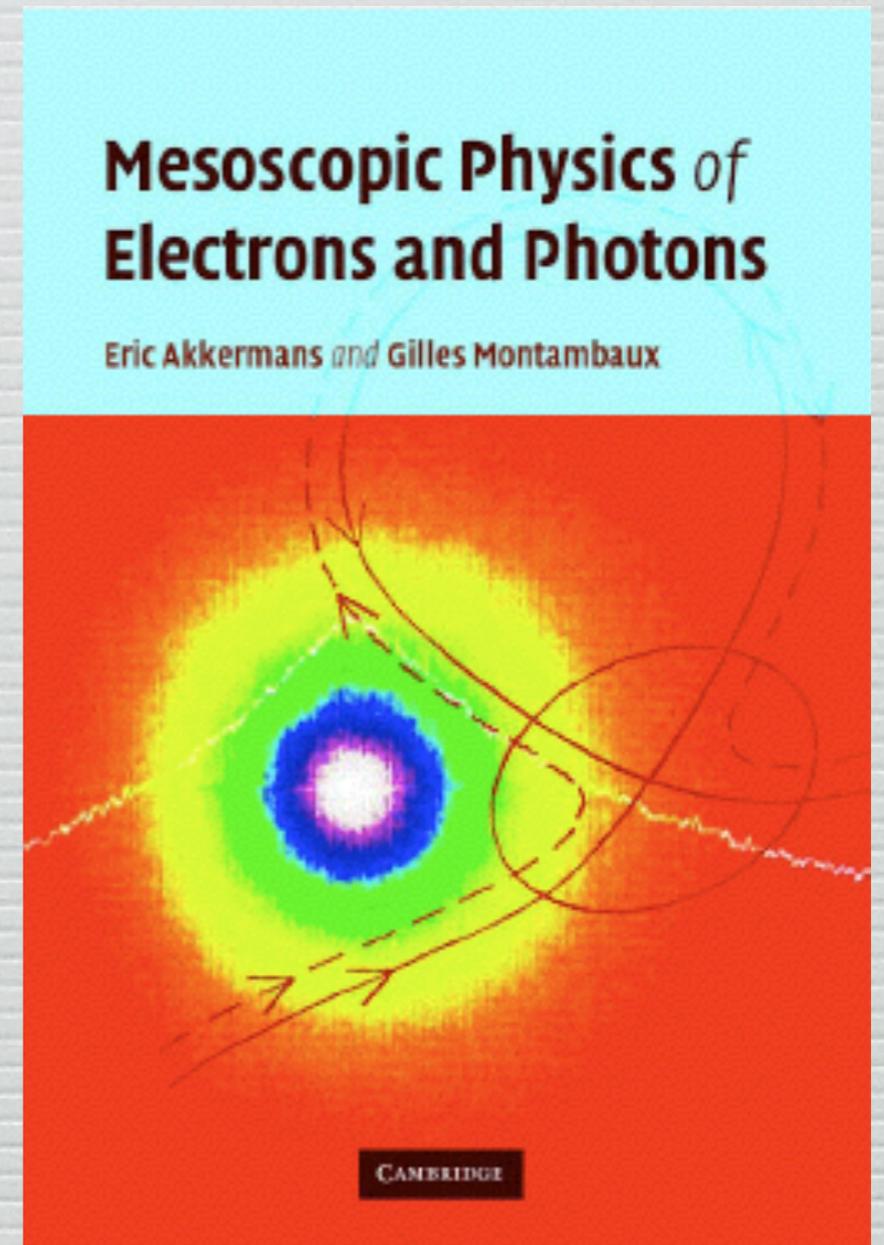
Aux frontieres de la physique mesoscopique,
Mont Orford Quebec, Canada,

Part 1

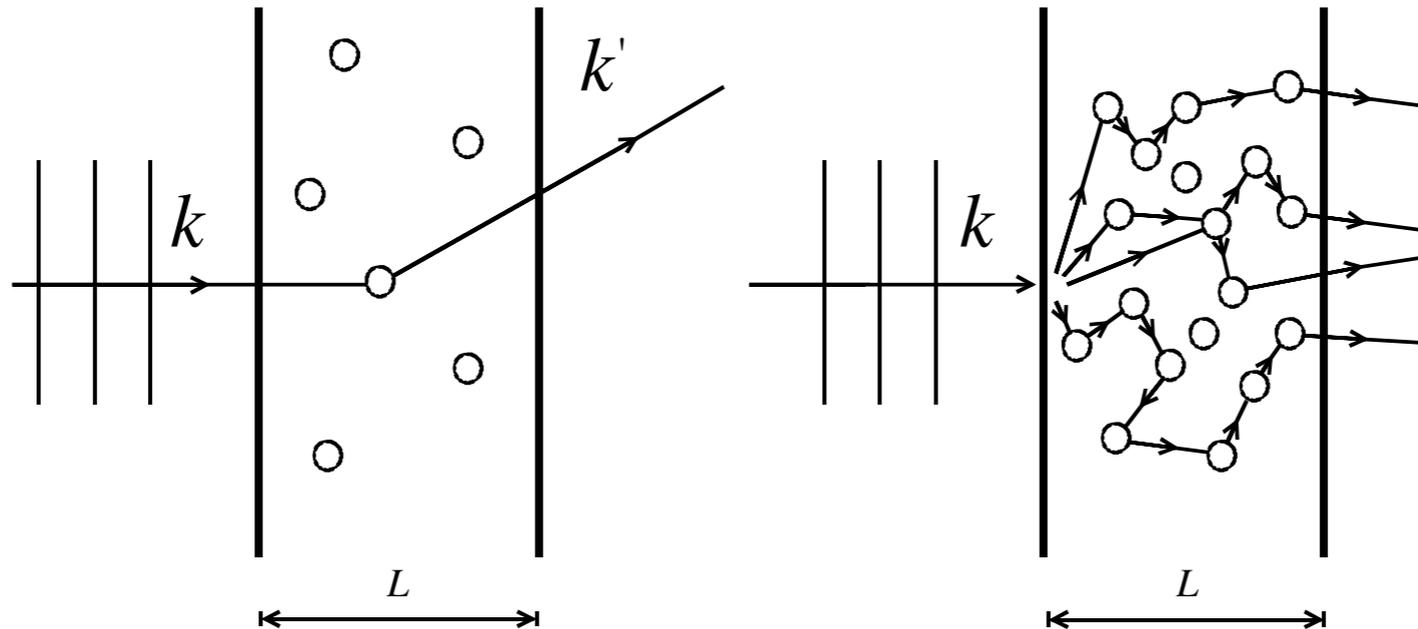
What is quantum in mesoscopic physics ? Transport and interferences

(E.A., G. Montambaux)

more details in:



Multiple scattering of electrons



2 characteristic lengths:

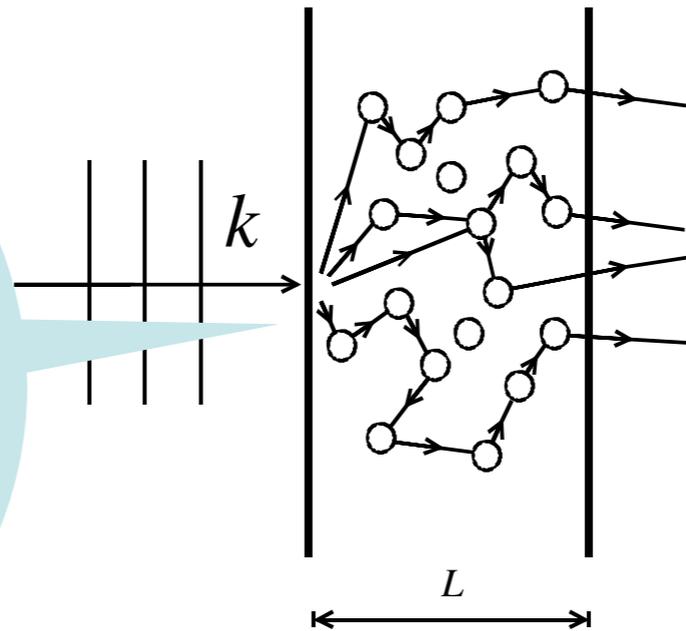
Wavelength: $\lambda_F = k_F^{-1}$

Elastic mean free path: l (Disorder - Origin ?)

Weak disorder $\lambda_F \ll l$: independent scattering events

Multiple scattering of electrons

We shall be interested only by this limit



2 characteristic lengths:

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Some “canonical”
mesoscopic effects

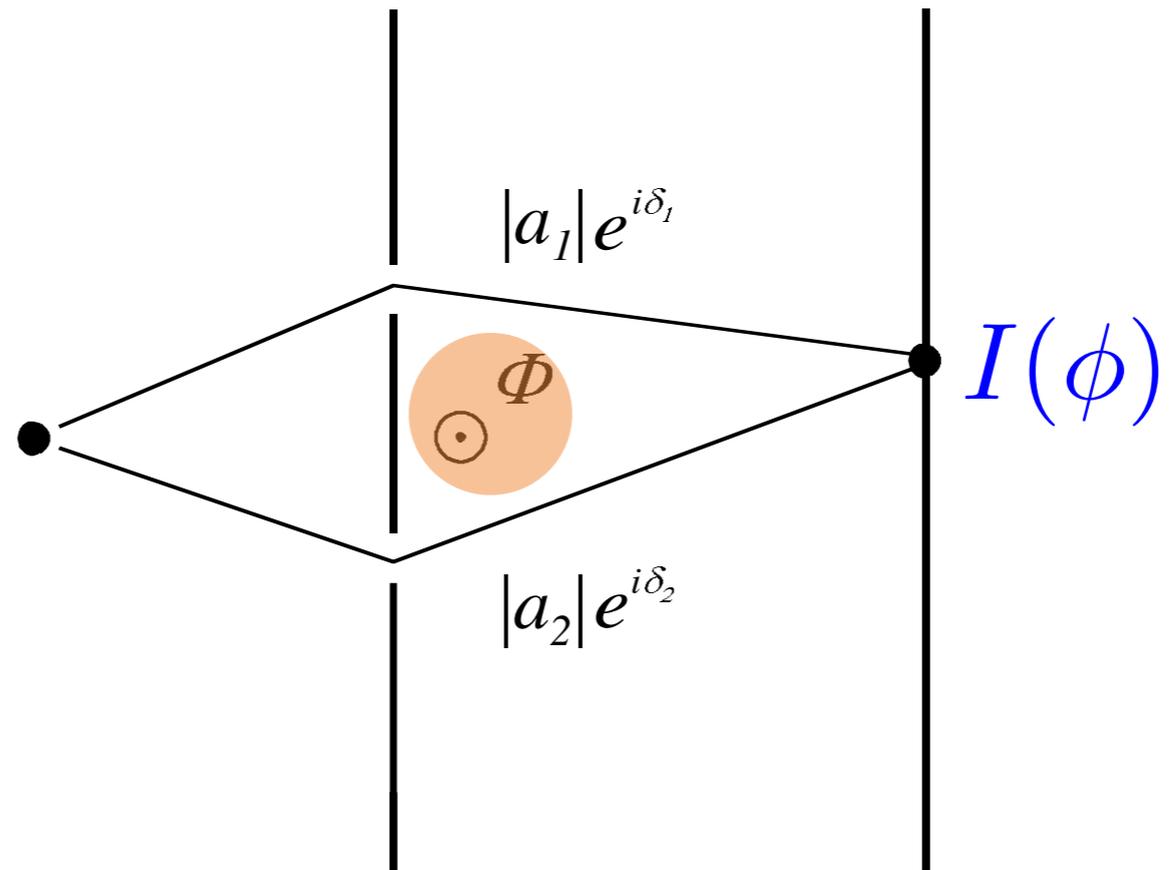
Some “canonical” mesoscopic effects

The Aharonov-Bohm effect

Aharonov-Bohm (1959)

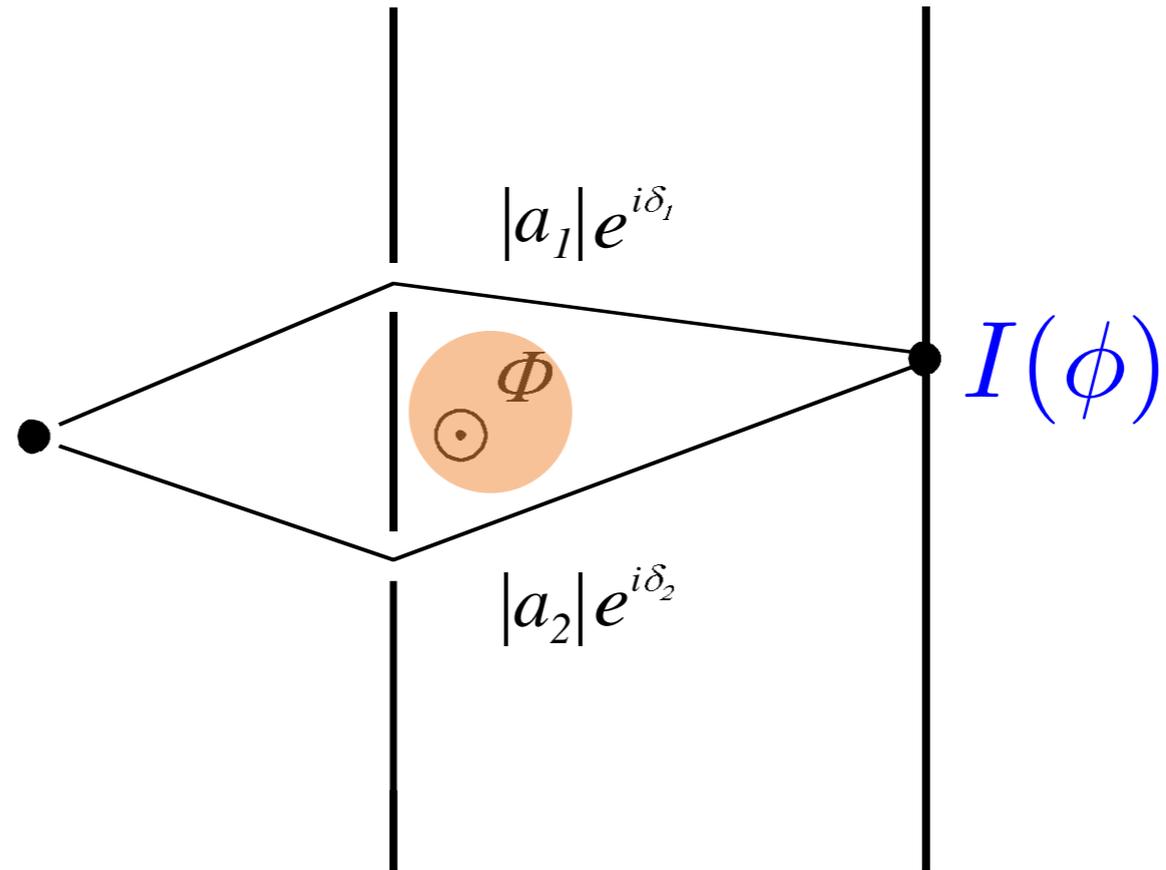
The setup : Young slits

No magnetic field on
the electrons : **no**
Lorentz force and no
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No magnetic field on the electrons : **no Lorentz force** and no orbital motion.



The quantum amplitudes $a_{1,2} = |a_{1,2}|e^{i\delta_{1,2}}$ have phases:

$$\delta_1 = \delta_1^{(0)} - \frac{e}{\hbar} \int_1 \mathbf{A} \cdot d\mathbf{l} \quad \text{and} \quad \delta_2 = \delta_2^{(0)} - \frac{e}{\hbar} \int_2 \mathbf{A} \cdot d\mathbf{l}$$

The intensity $I(\phi)$ is given by

$$I(\phi) = |a_1 + a_2|^2 = |a_1|^2 + |a_2|^2 + 2|a_1 a_2| \cos(\delta_1 - \delta_2)$$
$$= I_1 + I_2 + 2\sqrt{I_1 I_2} \cos(\delta_1 - \delta_2)$$

The phase difference $\Delta\delta(\phi) = \delta_1 - \delta_2$ is modulated by the magnetic flux ϕ :

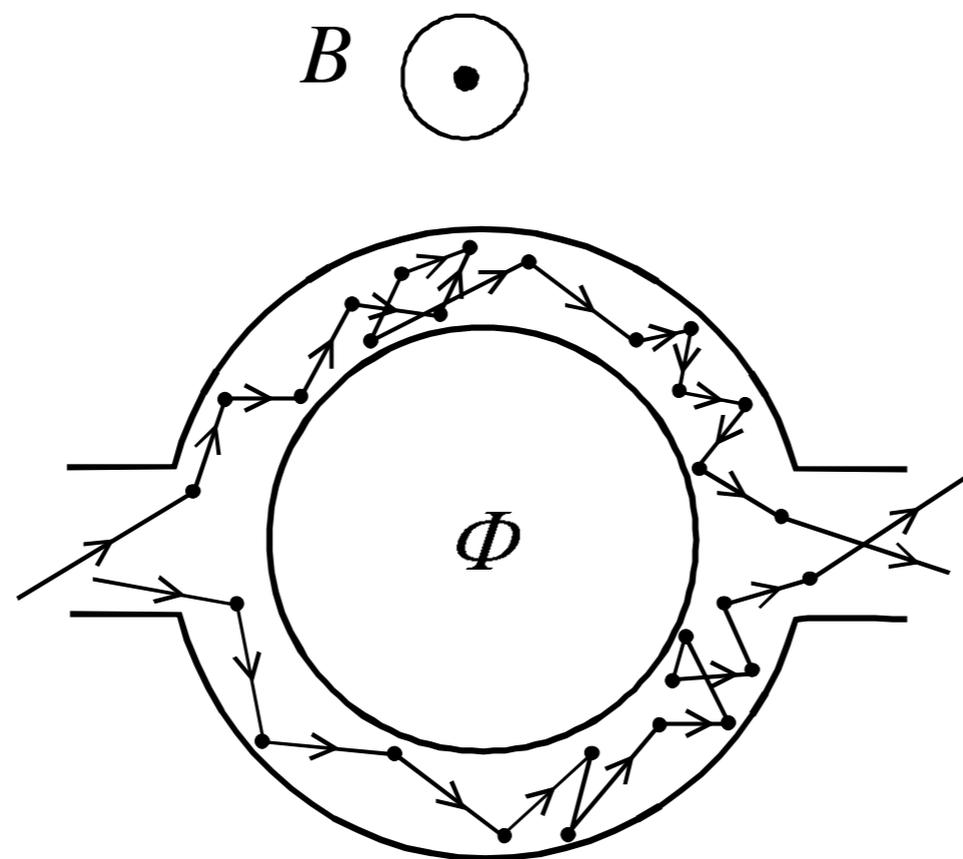
$$\Delta\delta(\phi) = \delta_1^{(0)} - \delta_2^{(0)} + \frac{e}{\hbar} \oint \mathbf{A} \cdot d\mathbf{l} = \Delta\delta^{(0)} + 2\pi \frac{\phi}{\phi_0}$$

where $\phi_0 = h/e$ is the quantum of magnetic flux.

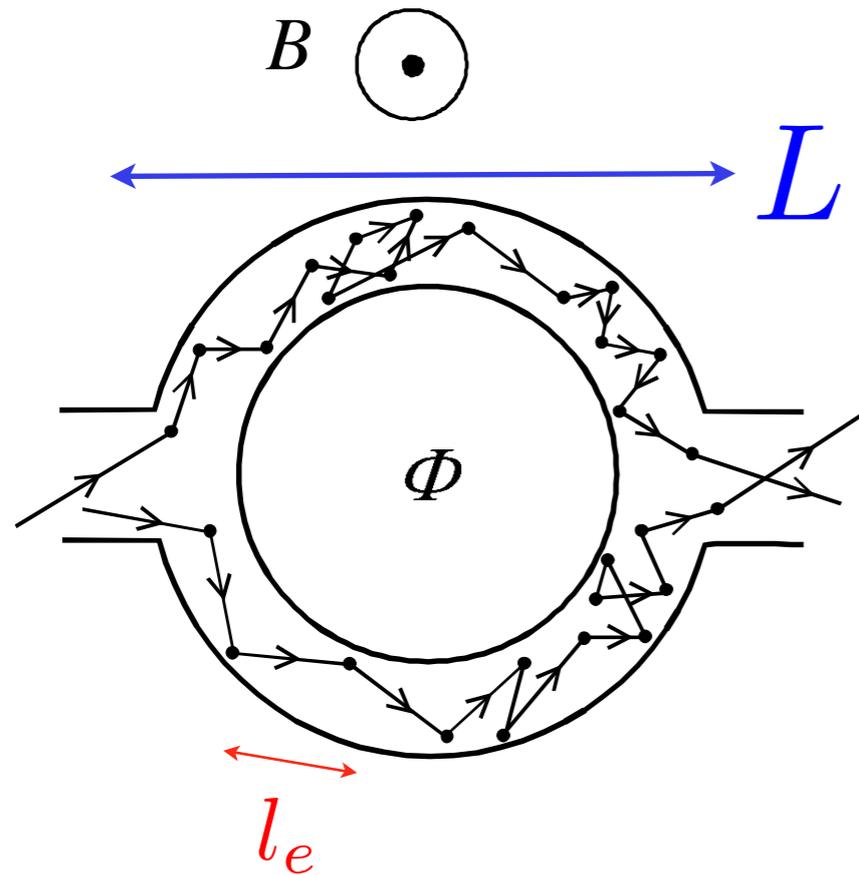
There is a continuous change of the state of interference:

Aharonov-Bohm effect (1959)

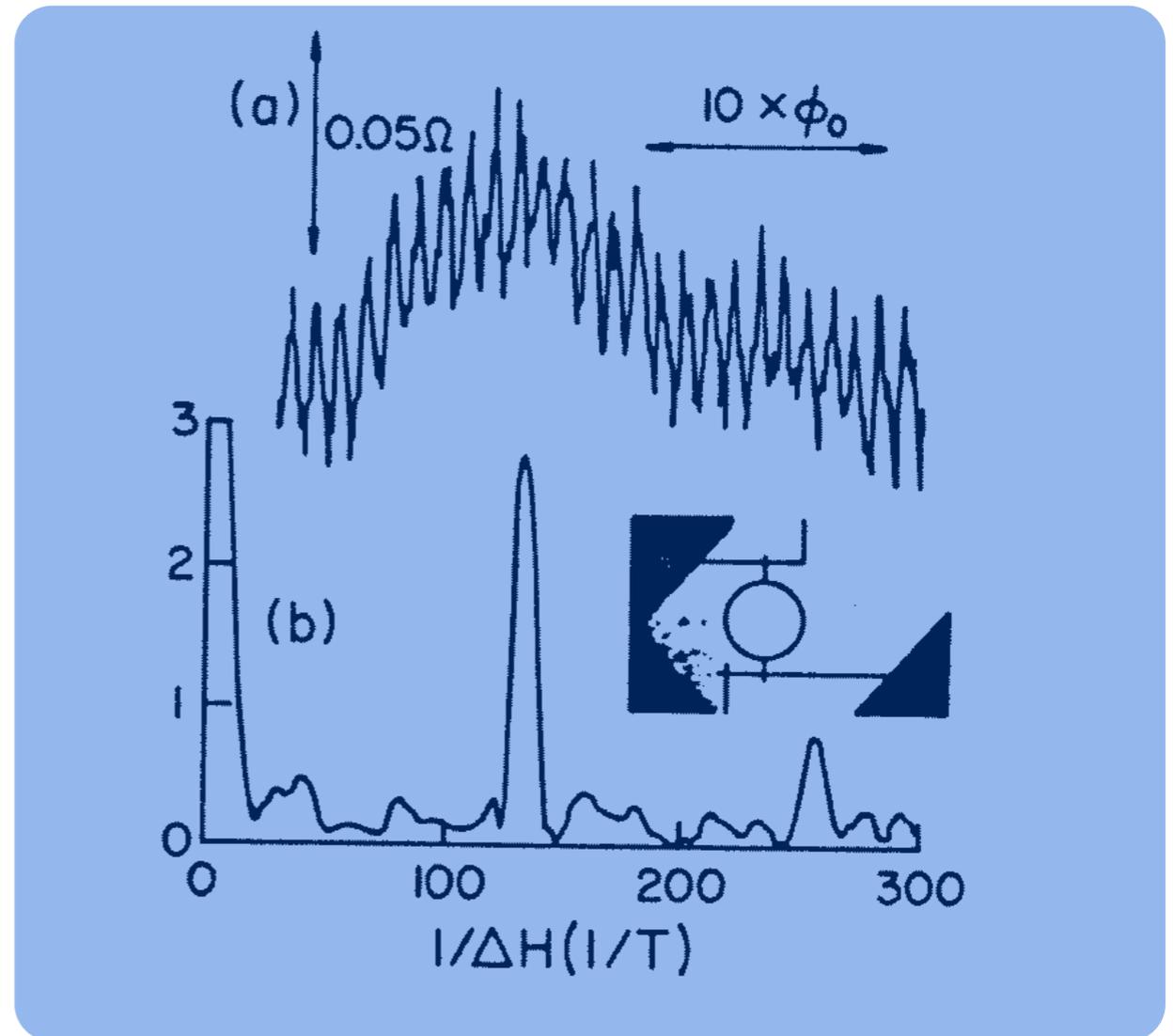
- Aharonov-Bohm effect in disordered metals



Implementation in metals : the conductance $G(\phi)$ is the analog of the intensity.



elastic mean free path



$$G(\phi) = G_0 + \delta G \cos(\Delta\delta^{(0)} + 2\pi \frac{\phi}{\phi_0})$$

Webb et al. 1985

Phase coherent effects subsist in **disordered** metals.

Reconsider the Drude theory.

Phase coherence and effect of disorder

The *Webb* experiment has been realized on a ring of size $L \simeq 1\mu$.
For a macroscopic normal metal, coherent effects are washed out.

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Quantum coherence: gas of quantum particles in a finite volume

Quantum states of the gas are **coherent superposition** of single particle states and they extend over the total volume (*ex. superconductivity, superfluidity, free electron gas*).

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Quantum coherence: gas of quantum particles in a finite volume

Quantum states of the gas are **coherent superposition** of single particle states and they extend over the total volume (*ex. superconductivity, superfluidity, free electron gas*).

For the electron gas, coherence disappears at non zero temperature so that we can use a **classical description** of transport and thermodynamics

Vanishing of quantum coherence results from the existence of **incoherent** and **irreversible** processes associated to the coupling of electrons to their surrounding (additional degrees of freedom) :

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Coupling to a bath of excitations : thermal excitations of the lattice (phonons)

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Chaotic dynamical systems (large recurrence times, Feynman chain)

Vanishing of quantum coherence results from the existence of **incoherent** and **irreversible** processes associated to the coupling of electrons to their surrounding (additional degrees of freedom) :

Coupling to a bath of excitations : thermal excitations of the lattice (phonons)

Chaotic dynamical systems (large recurrence times, Feynman chain)

Impurities with internal degrees of freedom (magnetic impurities)

Electron-electron interactions,....

The understanding of decoherence is difficult.

It is a great challenge in quantum mesoscopic physics.

The phase coherence length L_ϕ accounts in a generic way for decoherence processes.

The observation of coherent effects requires

$$L \ll L_\phi$$

Averaging over disorder ?

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Expect to wash up interference effects

Average coherence and multiple scattering

What is the role of elastic disorder ? Does it erase coherent effects ?

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Phase coherence leads to interference effects for a *given realization of disorder*.

Average coherence and multiple scattering

What is the role of **elastic** disorder ? Does it erase coherent effects ?

Phase coherence leads to interference effects for a *given realization of disorder*.

The *Webb* experiment corresponds to a fixed configuration of disorder.

Averaging over disorder \Rightarrow vanishing of the Aharonov-Bohm effect

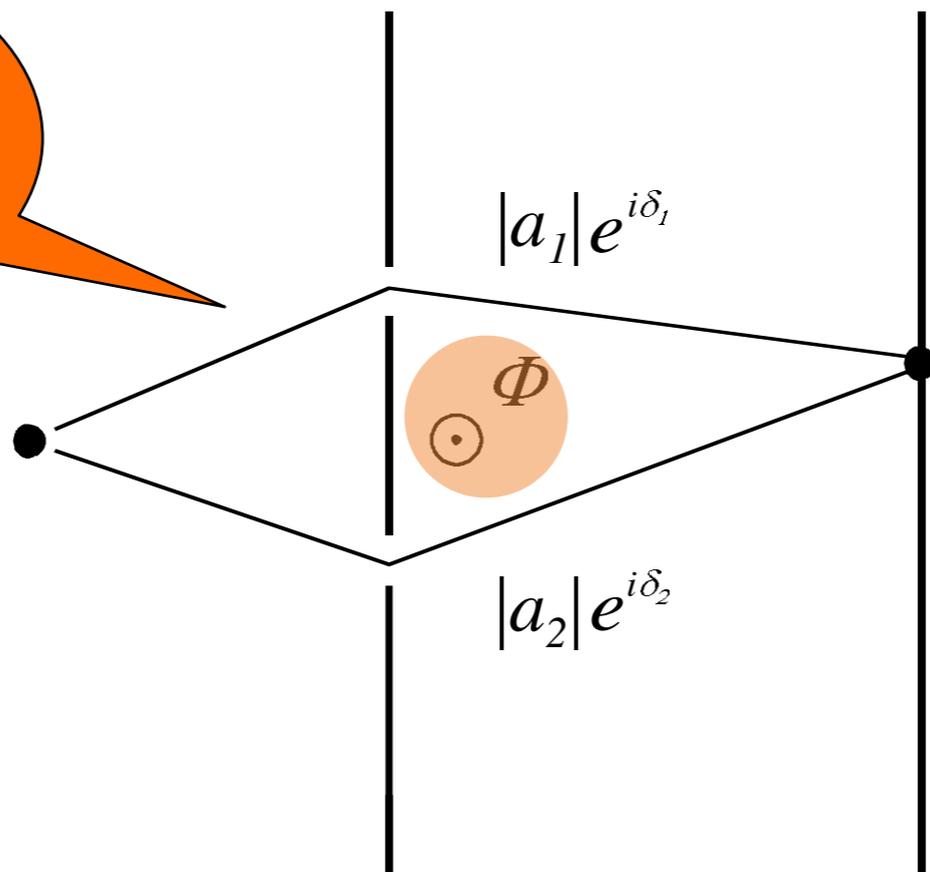
$$G(\phi) = G_0 + \delta G \cos\left(\Delta\delta^{(0)} + 2\pi \frac{\phi}{\phi_0}\right)$$

$$\Rightarrow \langle G(\phi) \rangle = G_0$$

Disorder seems to erase coherent effects....

The setup : Young slits

A reminder

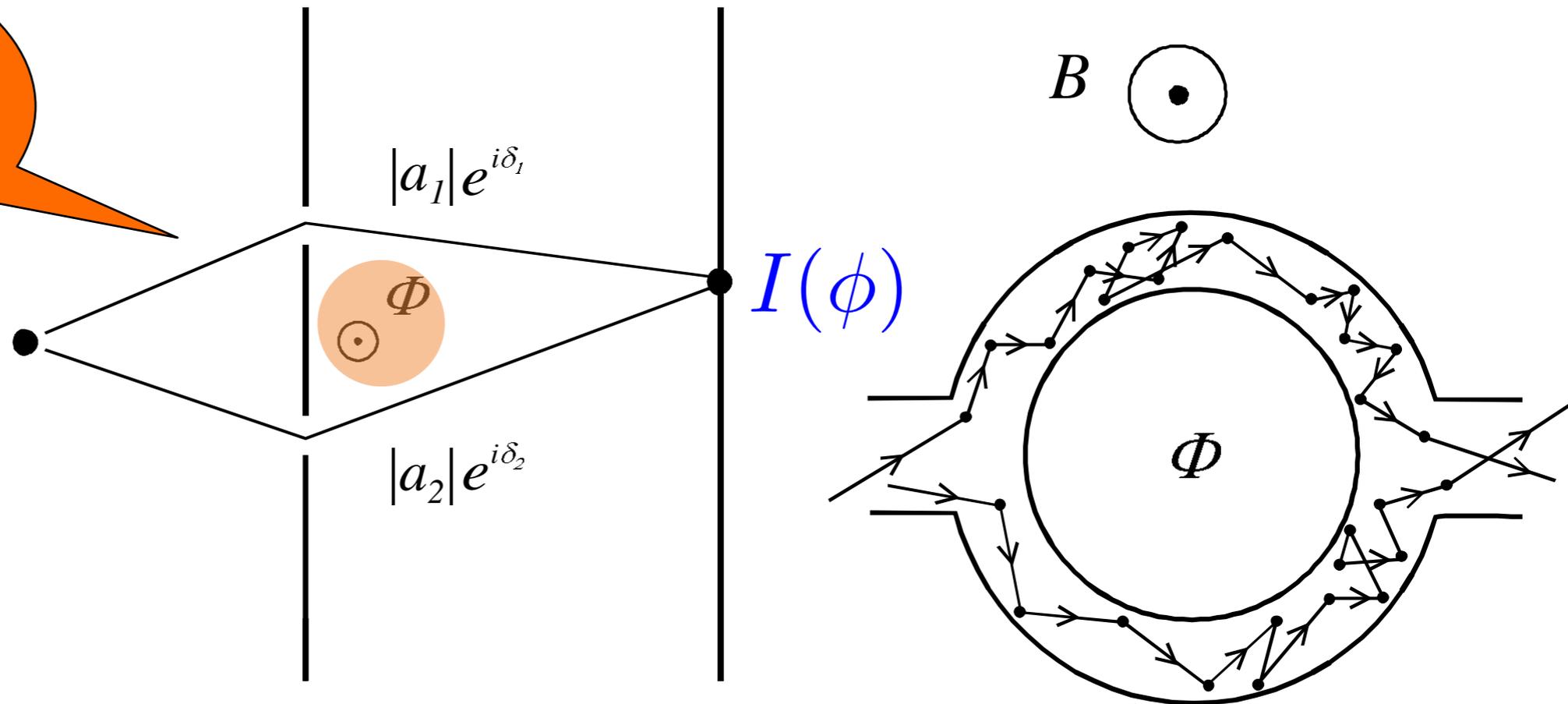


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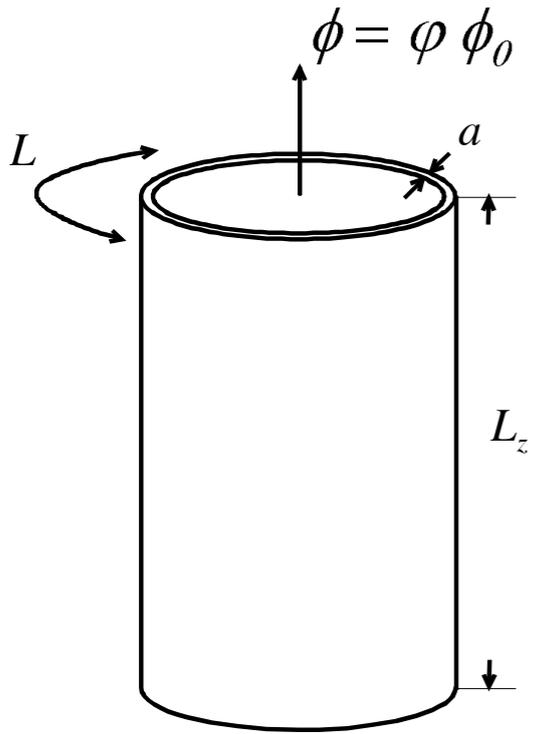
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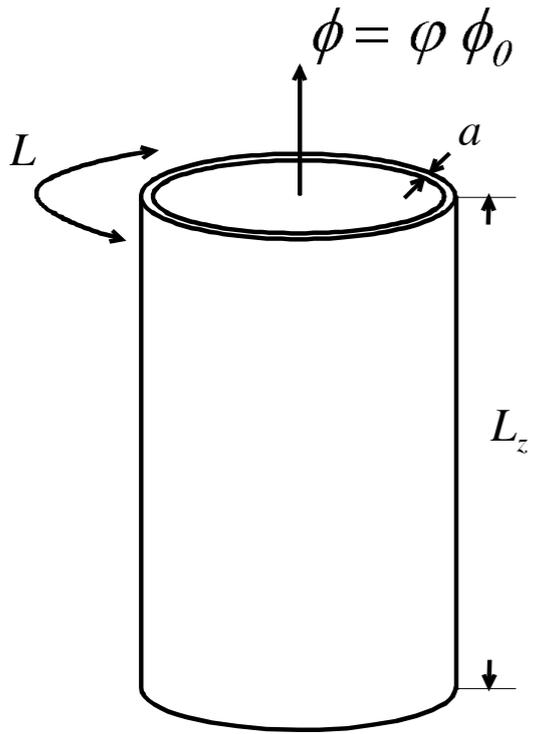
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The Sharvin² experiment



Experiment analogous to that of *Webb* but performed on a hollow cylinder of **height larger than L_ϕ** pierced by a Aharonov-Bohm flux. **Ensemble of rings identical to those of *Webb* but incoherent between themselves.**

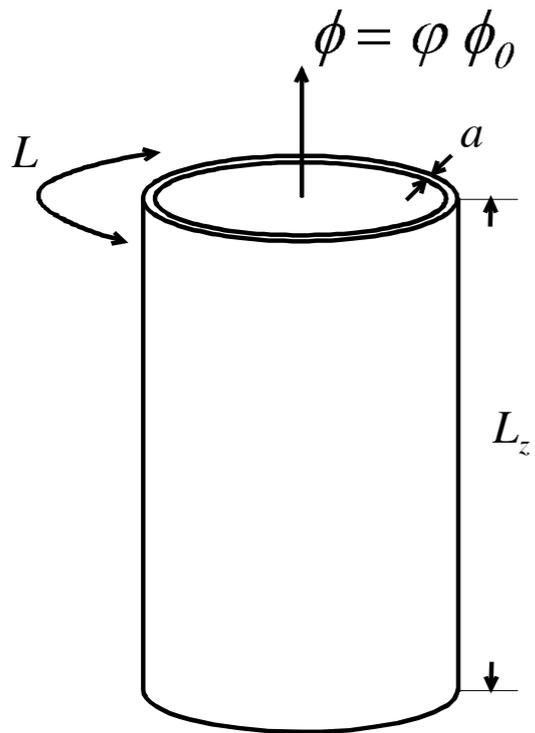
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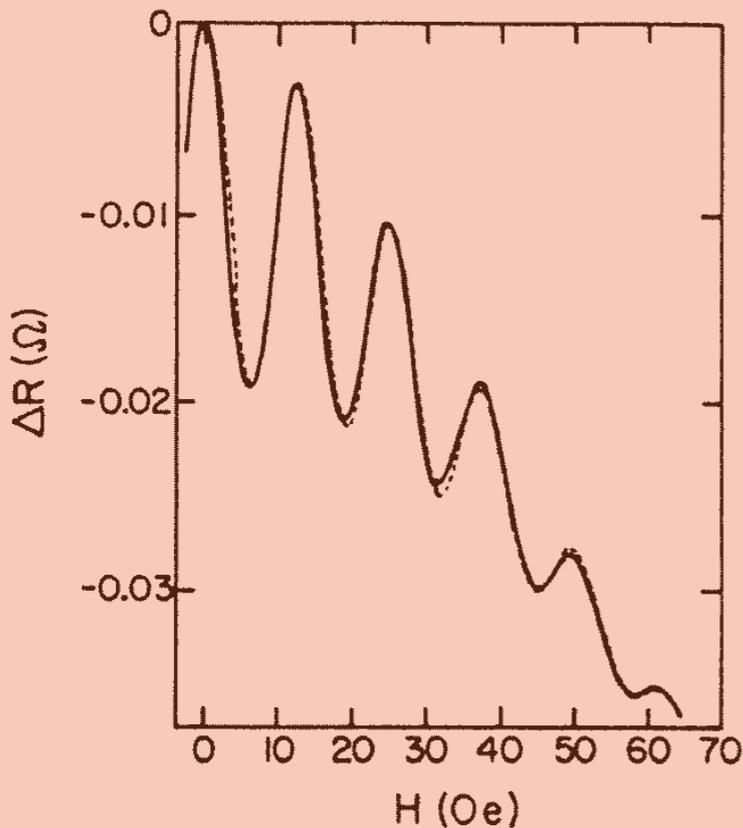
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The signal modulated at ϕ_0 *disappears*

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The signal modulated at ϕ_0 *disappears* but, instead, it appears a **new contribution** modulated at $\phi_0/2$

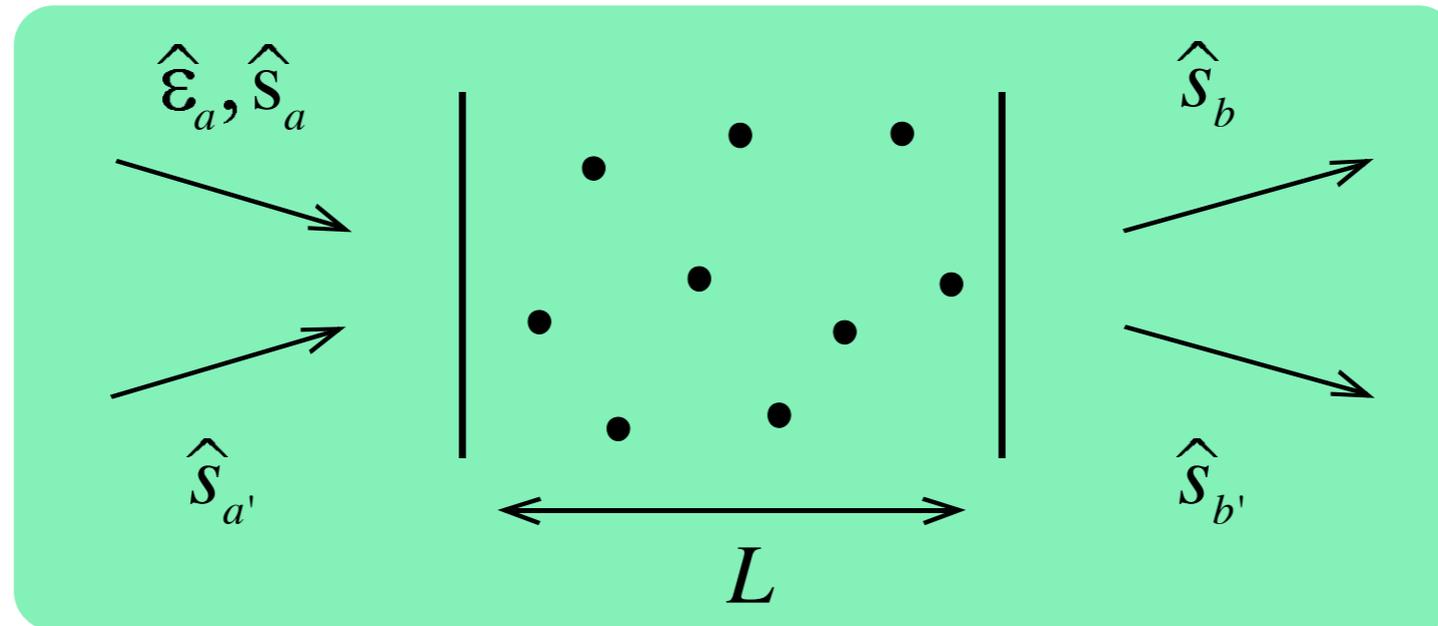
After all, disorder does not seem to erase coherent effects, but to modify them....

Some “canonical” mesoscopic effects

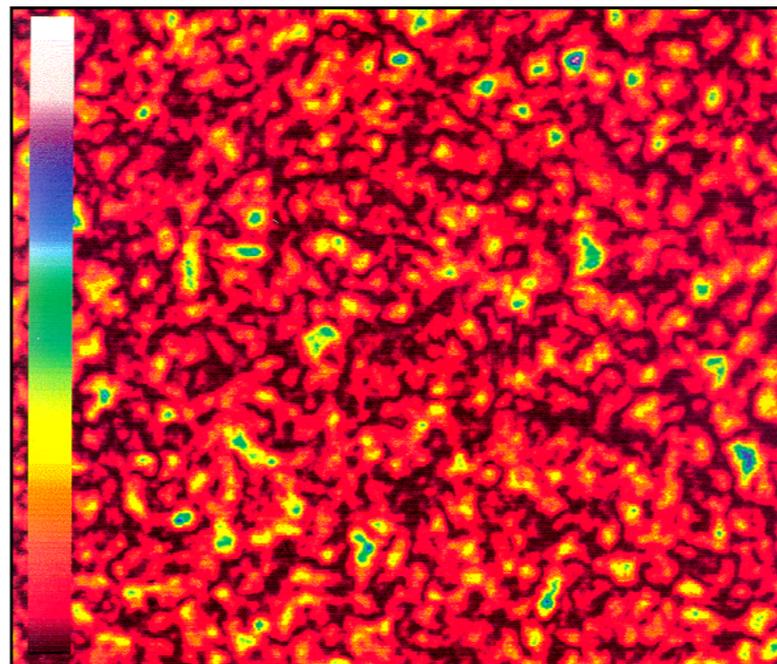
Coherent backscattering in
optics

An analogous problem: *Speckle patterns in optics*

Consider the elastic multiple scattering of light transmitted through a fixed disorder configuration.

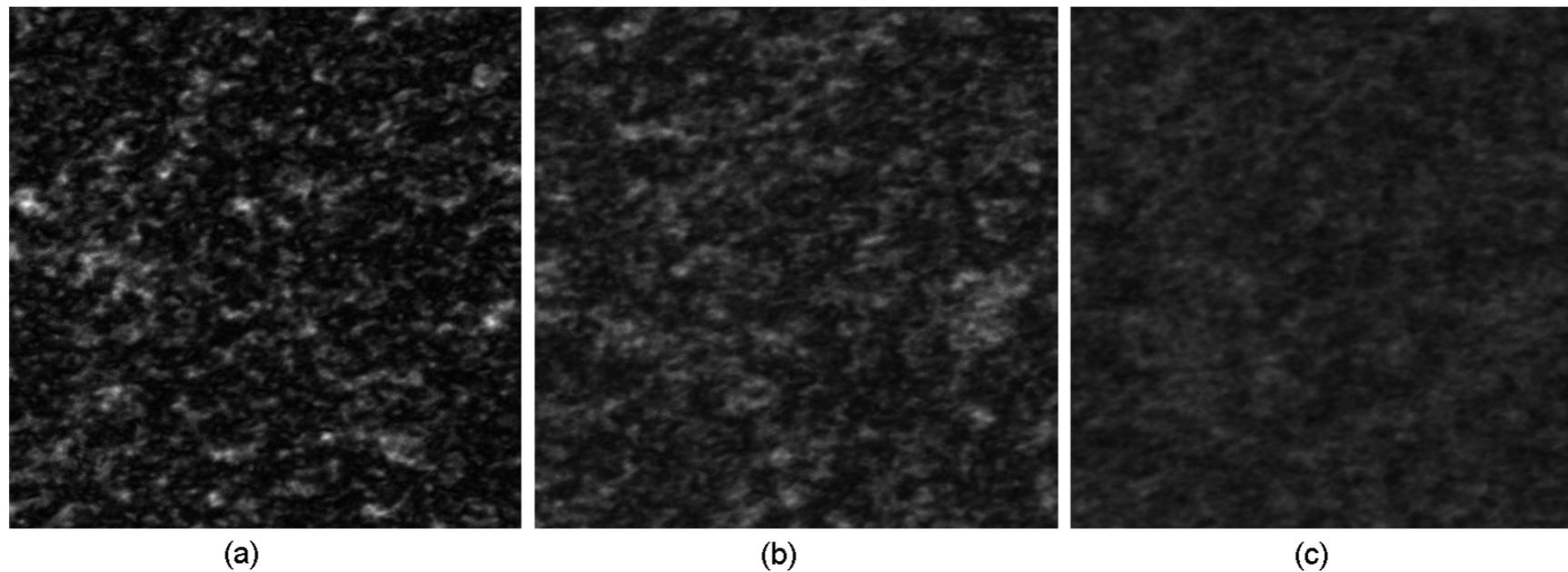


Outgoing light builds a **speckle pattern** *i.e.*, an **interference** picture:



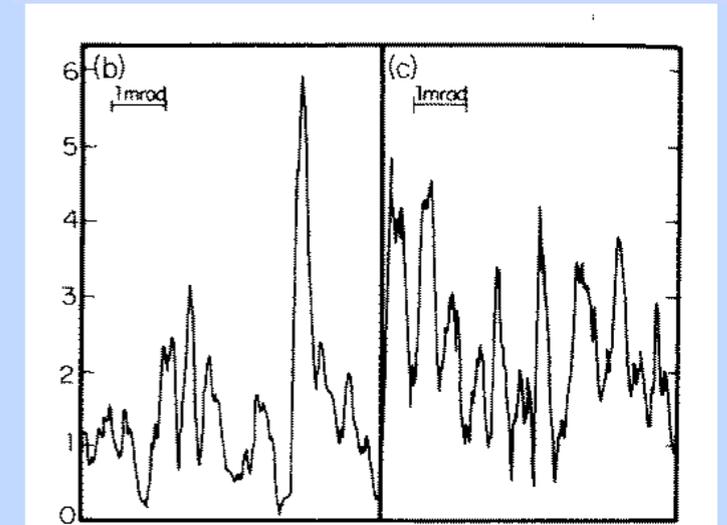
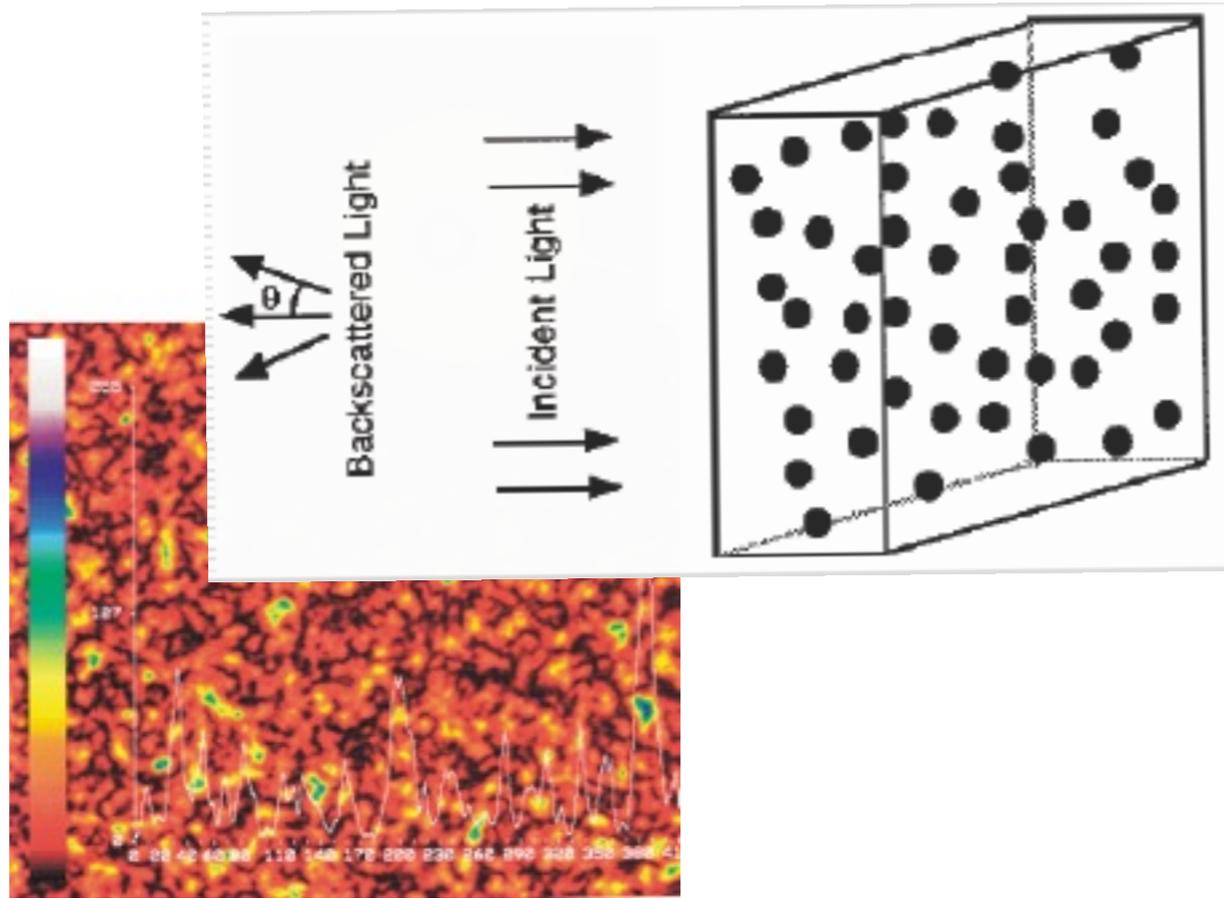
Averaging over disorder erases the speckle pattern:

Integration over the motion of the scatterers leads to self-averaging

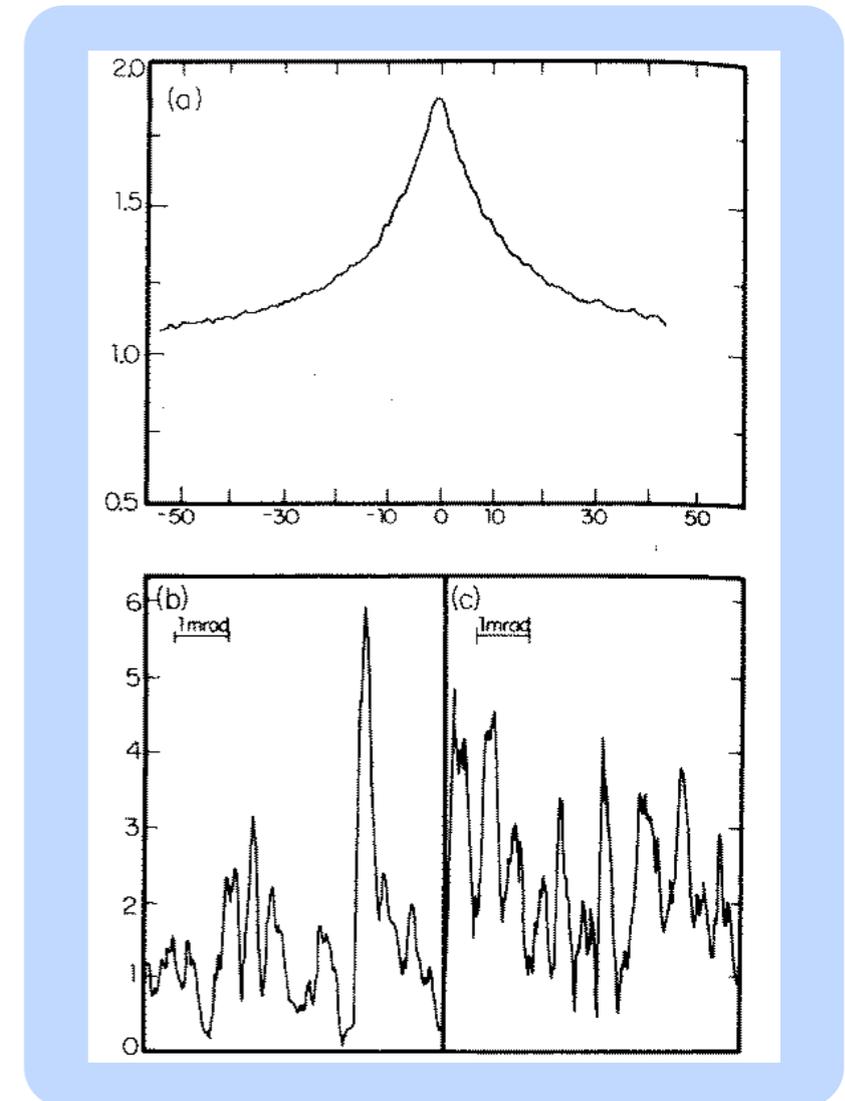
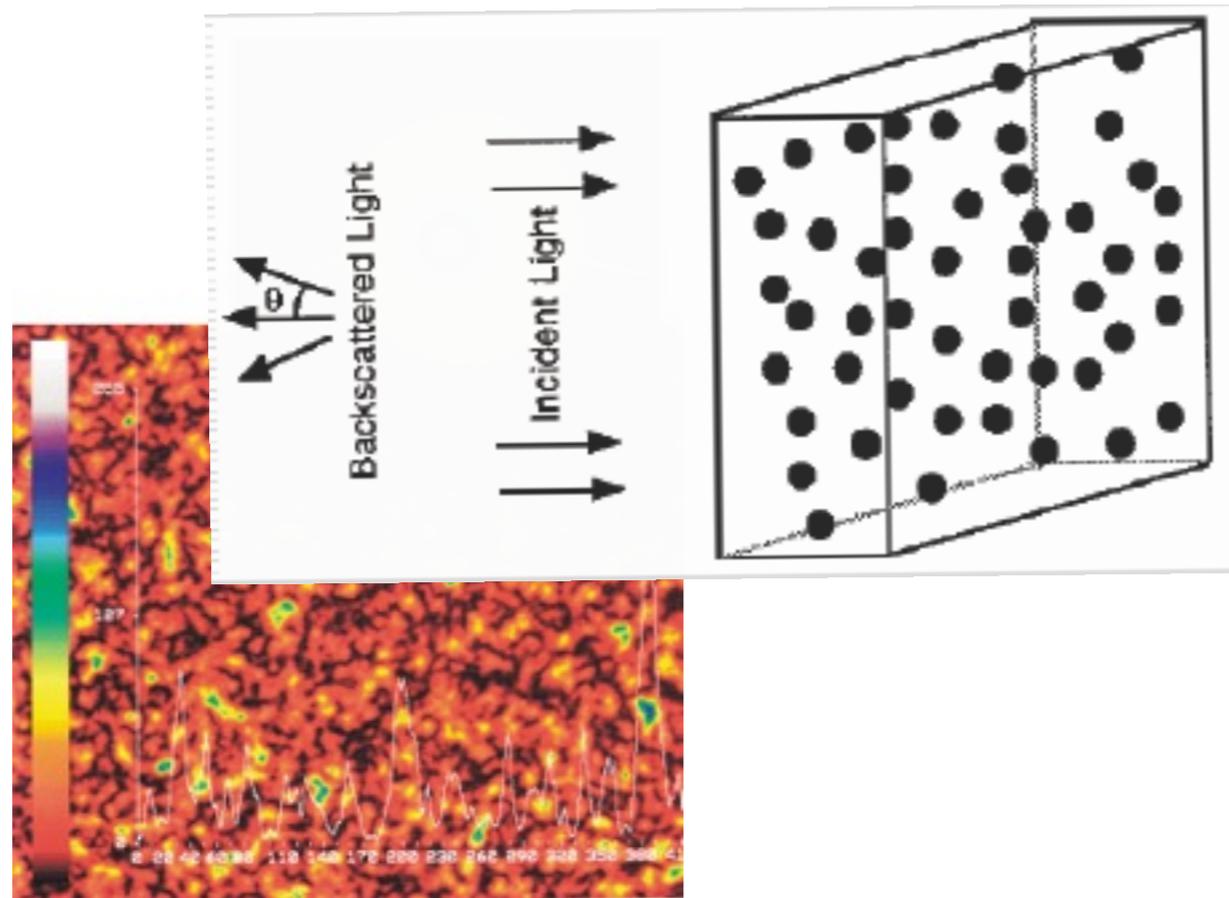


Time averaging

Does it erase all interferences ?



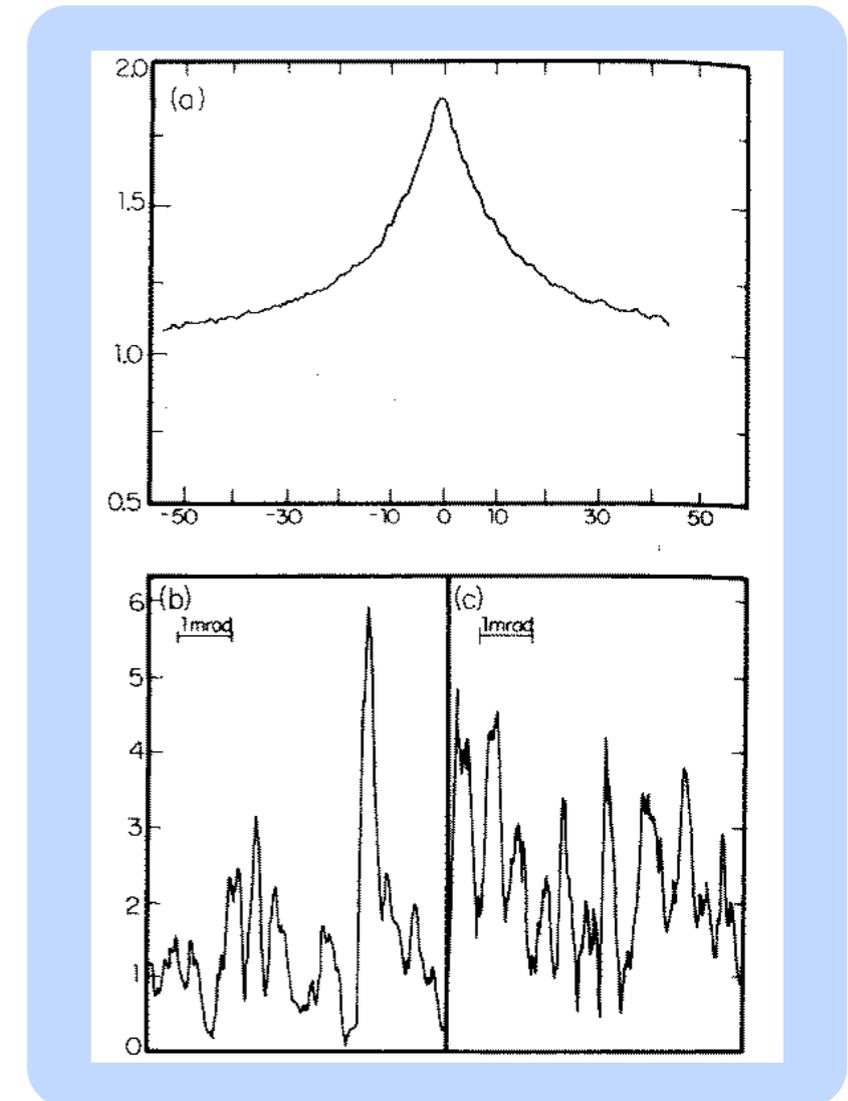
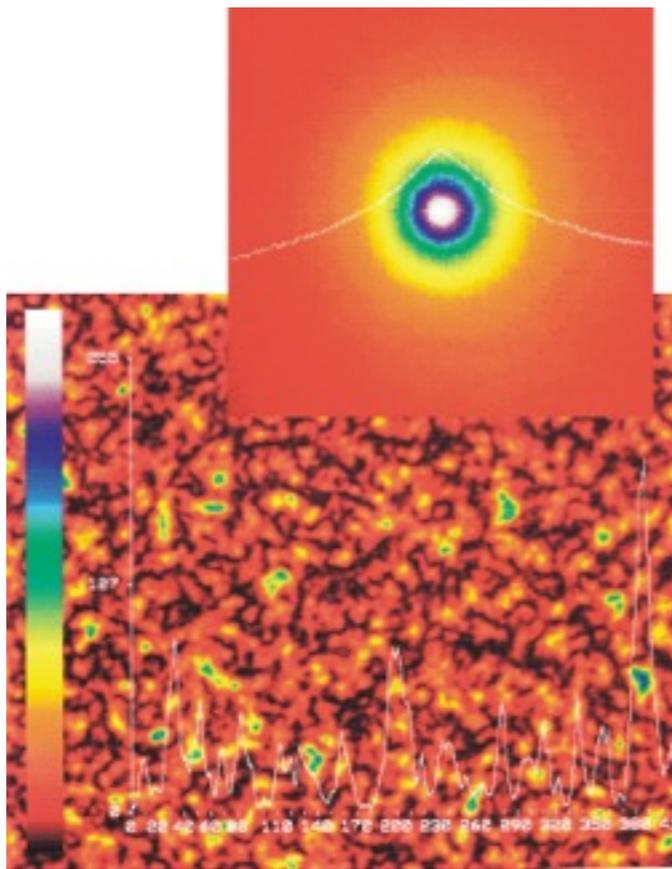
Does it erase all interferences ?



Averaging over disorder does not produce incoherent intensity only, but also an angular dependent part, the coherent backscattering, which is a coherence effect. We may conclude:

Elastic disorder is **not related** to decoherence : **disorder does not destroy phase coherence and does not introduce irreversibility.**

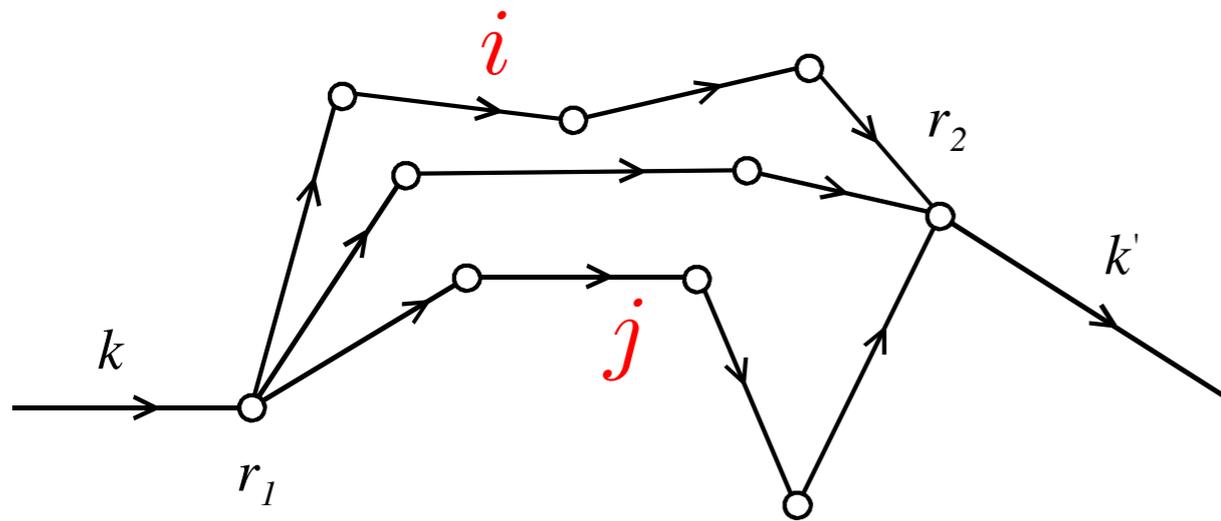
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How to understand average coherent effects ?



Complex amplitude $A(\mathbf{k}, \mathbf{k}')$ associated to the multiple scattering of a wave (electron or photon) incident with a wave vector \mathbf{k} and outgoing with \mathbf{k}'

$$A(\mathbf{k}, \mathbf{k}') = \sum_{\mathbf{r}_1, \mathbf{r}_2} f(\mathbf{r}_1, \mathbf{r}_2) e^{i(\mathbf{k} \cdot \mathbf{r}_1 - \mathbf{k}' \cdot \mathbf{r}_2)}$$

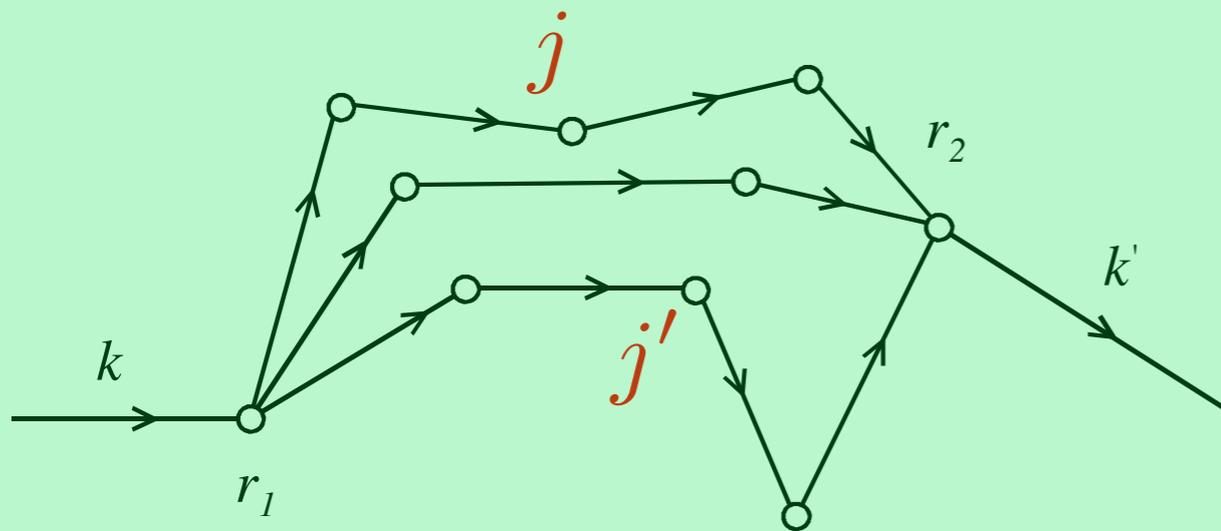
the complex amplitude $f(\mathbf{r}_1, \mathbf{r}_2) = \sum_j |a_j| e^{i\delta_j}$ describes the propagation of the wave between \mathbf{r}_1 and \mathbf{r}_2 .

The corresponding intensity is

$$|A(\mathbf{k}, \mathbf{k}')|^2 = \sum_{\mathbf{r}_1, \mathbf{r}_2} \sum_{\mathbf{r}_3, \mathbf{r}_4} f(\mathbf{r}_1, \mathbf{r}_2) f^*(\mathbf{r}_3, \mathbf{r}_4) e^{i(\mathbf{k} \cdot \mathbf{r}_1 - \mathbf{k}' \cdot \mathbf{r}_2)} e^{-i(\mathbf{k} \cdot \mathbf{r}_3 - \mathbf{k}' \cdot \mathbf{r}_4)}$$

with

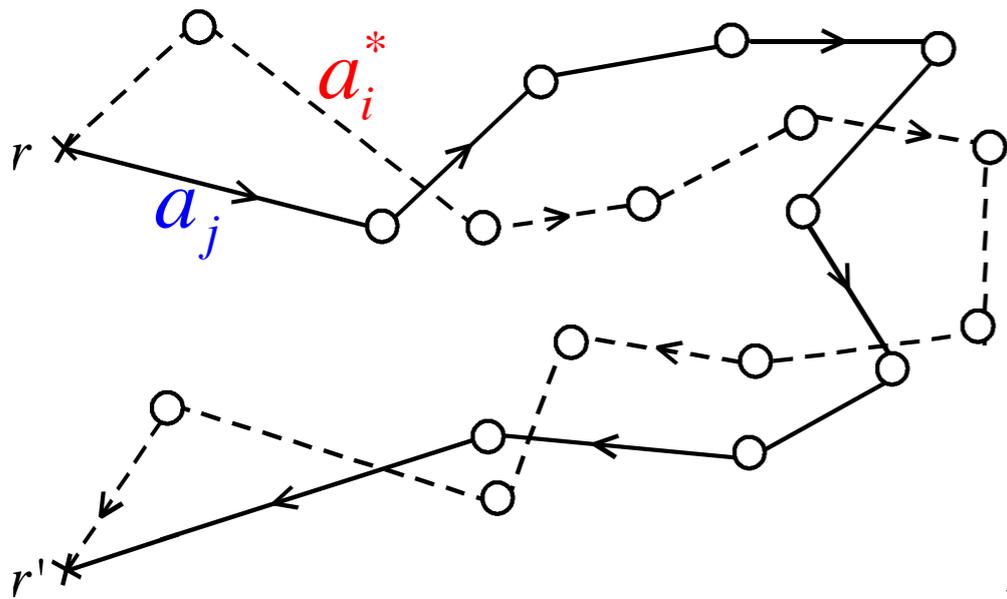
$$f(\mathbf{r}_1, \mathbf{r}_2) f^*(\mathbf{r}_3, \mathbf{r}_4) = \sum_{j, j'} a_j(\mathbf{r}_1, \mathbf{r}_2) a_{j'}^*(\mathbf{r}_3, \mathbf{r}_4) = \sum_{j, j'} |a_j| |a_{j'}| e^{i(\delta_j - \delta_{j'})}$$



On average over disorder, most contributions to $f f^*$ disappear since the dephasing $\delta_j - \delta_{j'} \gg 1$

The only remaining contributions to the intensity correspond to terms with **zero dephasing**, *i.e.*, to **identical trajectories**.

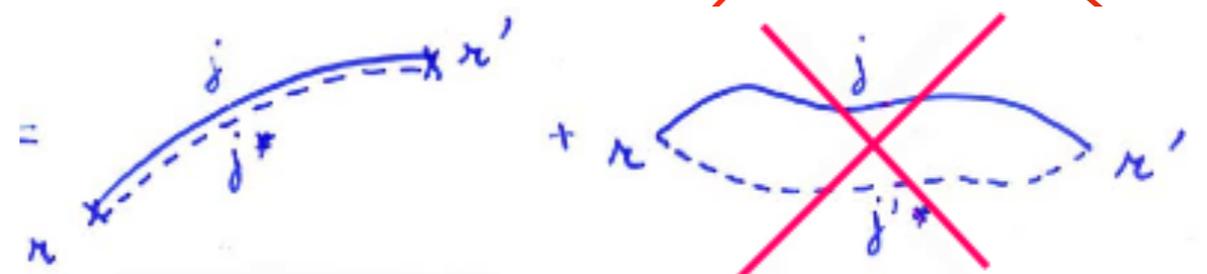
Quantum probability for propagation between two points



$$P(r, r') = \sum_{i,j} \overline{a_i^*(r, r') a_j(r, r')}$$

$$P(r, r') = \sum_j \overline{|a_j(r, r')|^2} + \sum_{i \neq j} \overline{a_i^*(r, r') a_j(r, r')}$$

$$a_i^* a_j = |a_i| |a_j| e^{i(\delta_i - \delta_j)}$$

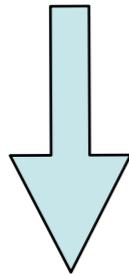
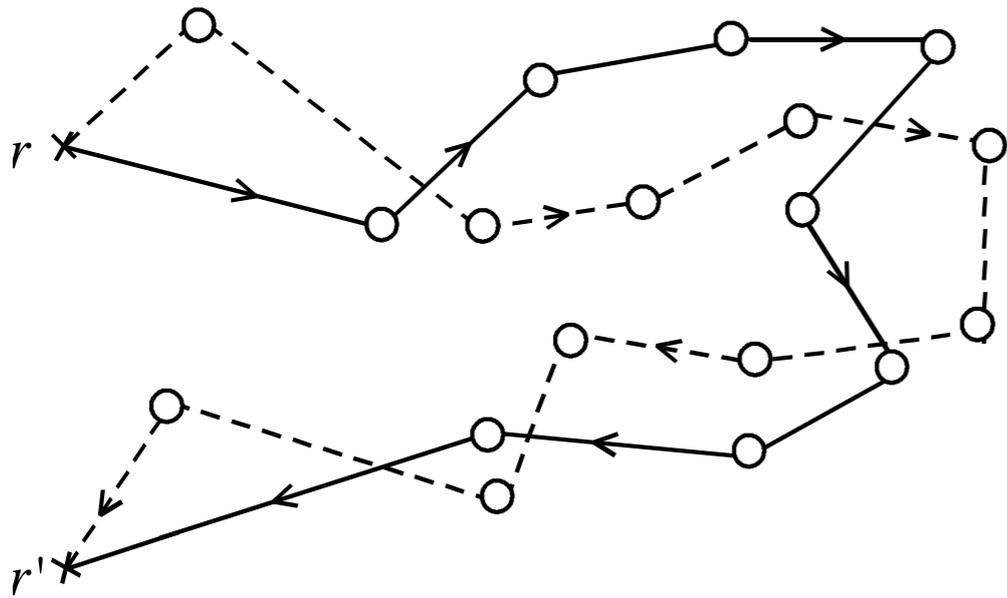


$$\delta_i - \delta_j \gg 1$$

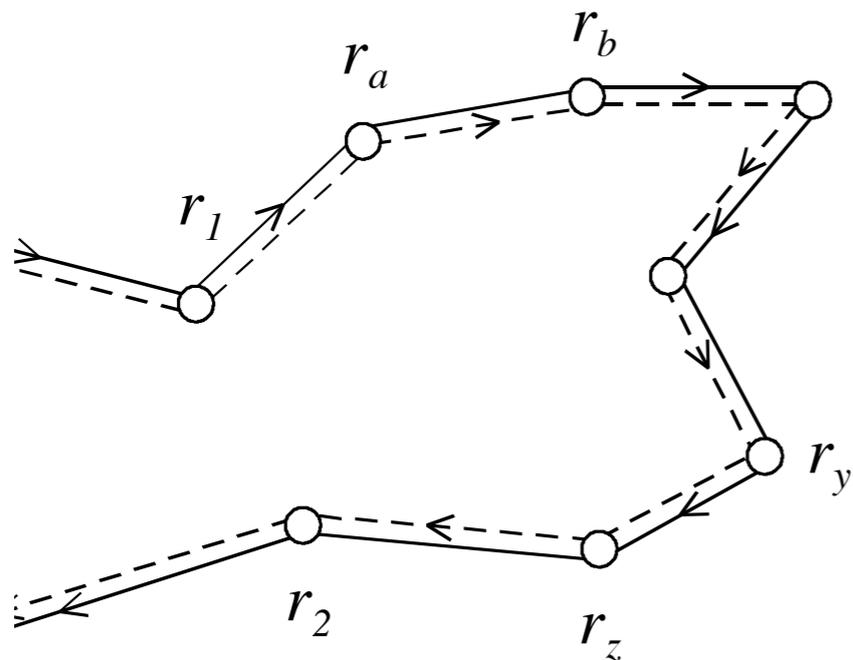
Incoherent propagation !

vanishes on average

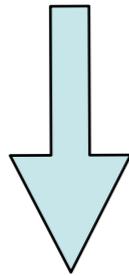
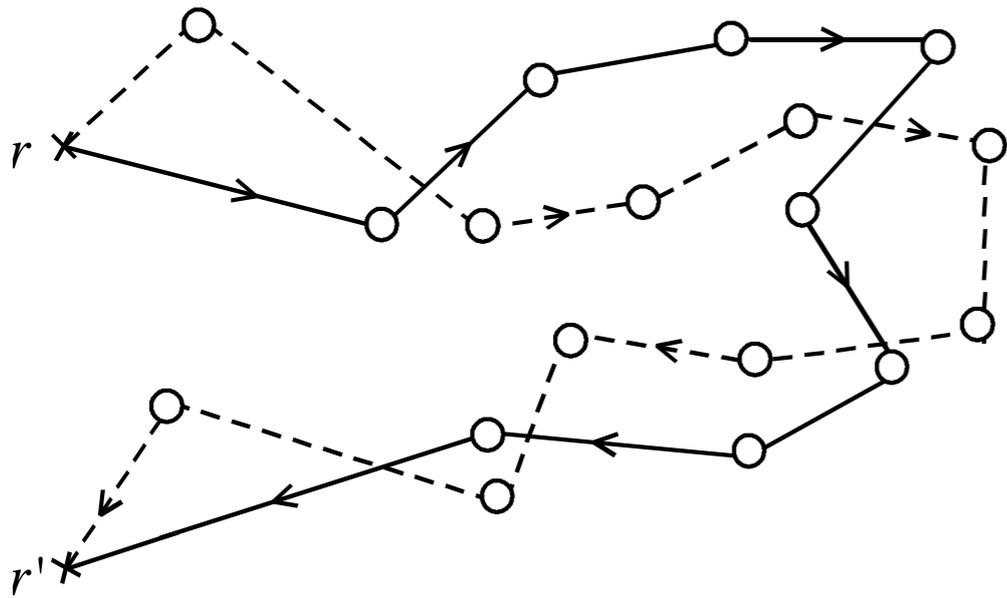
Some useful design



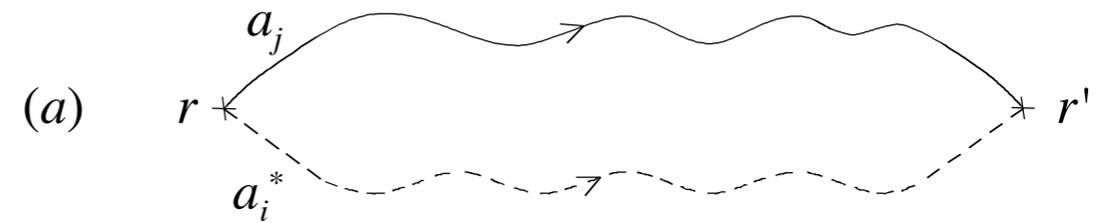
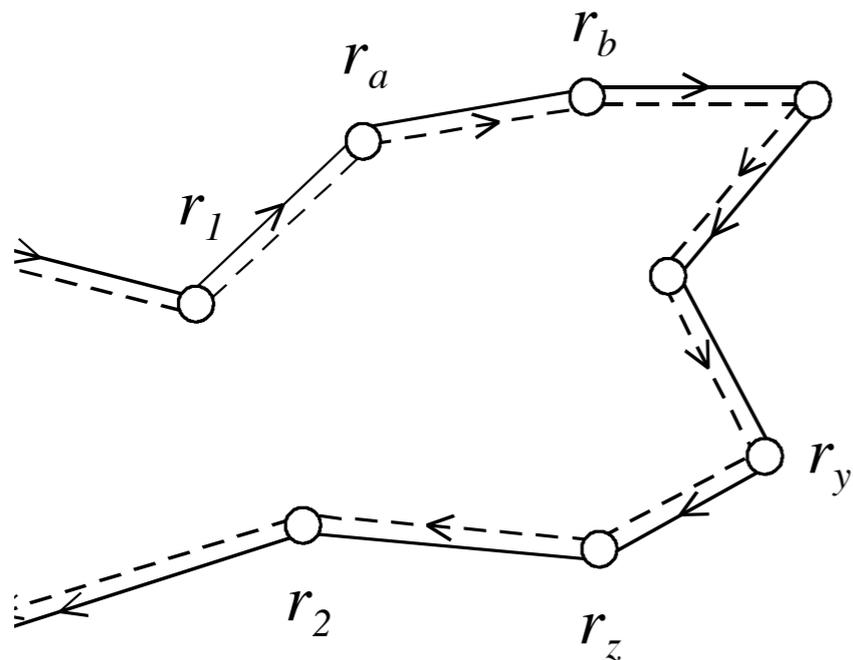
Disorder
averaging



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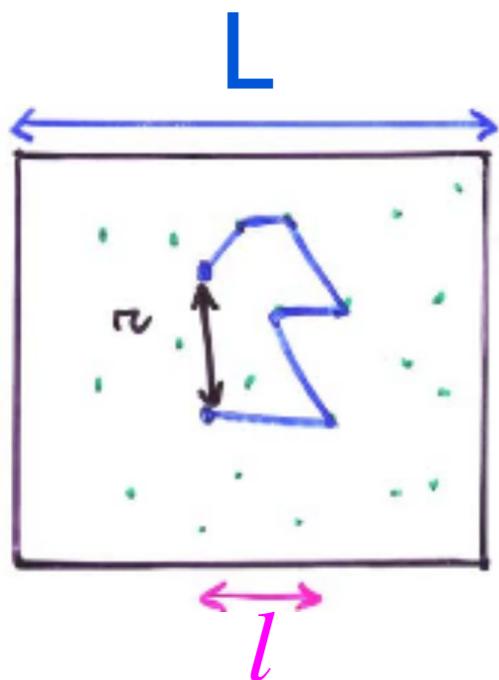
Disorder
averaging



To a good approximation, the **incoherent contribution** obeys a classical **diffusion equation**

$$\left(\frac{\partial}{\partial t} - D\Delta \right) P(r, r', t) = \delta(r - r')\delta(t) \Leftrightarrow (-i\omega + Dq^2)P(q, \omega) = 1$$

Incoherent electrons diffuse in the conductor with a *diffusion coefficient D* (*Drude theory*)

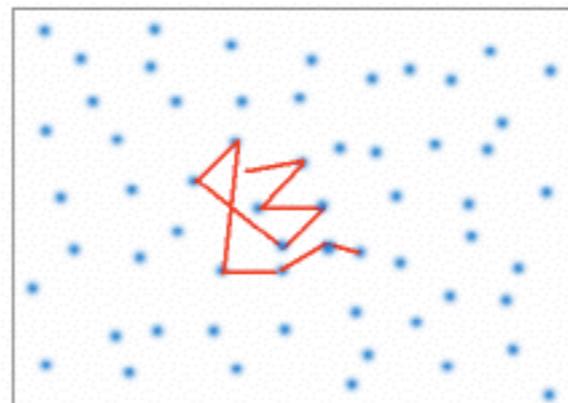


$$l \ll L$$

$$\langle r^2 \rangle = 2d Dt$$

space dimensionality

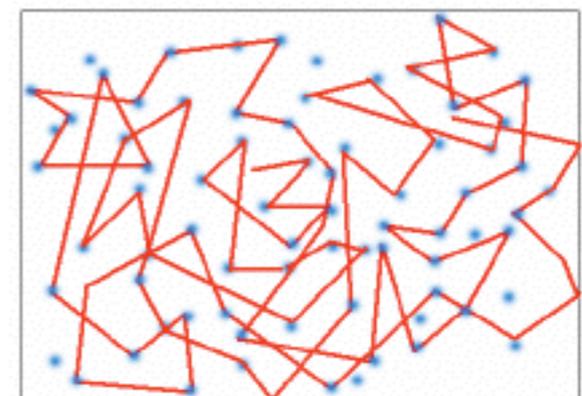
$$t \ll \tau_D$$



Thouless time

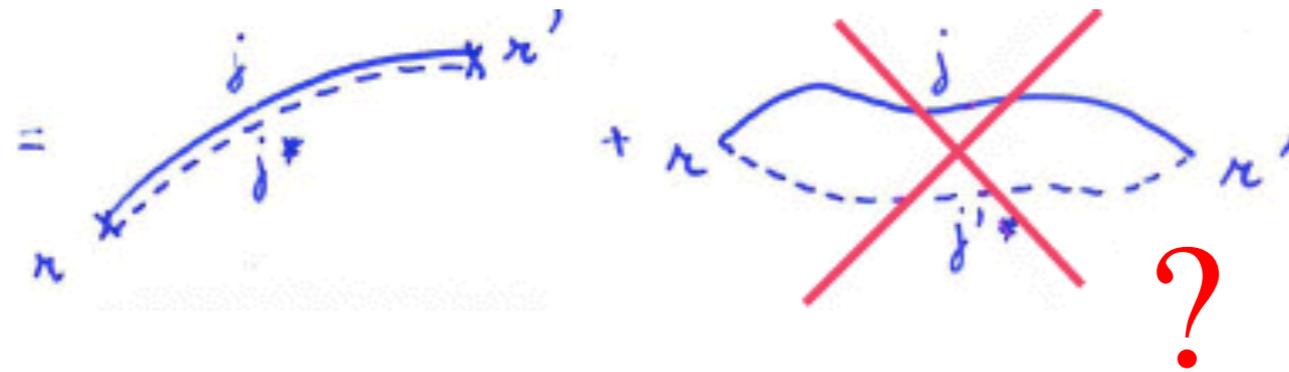
$$L^2 = D\tau_D$$

$$t \gg \tau_D$$



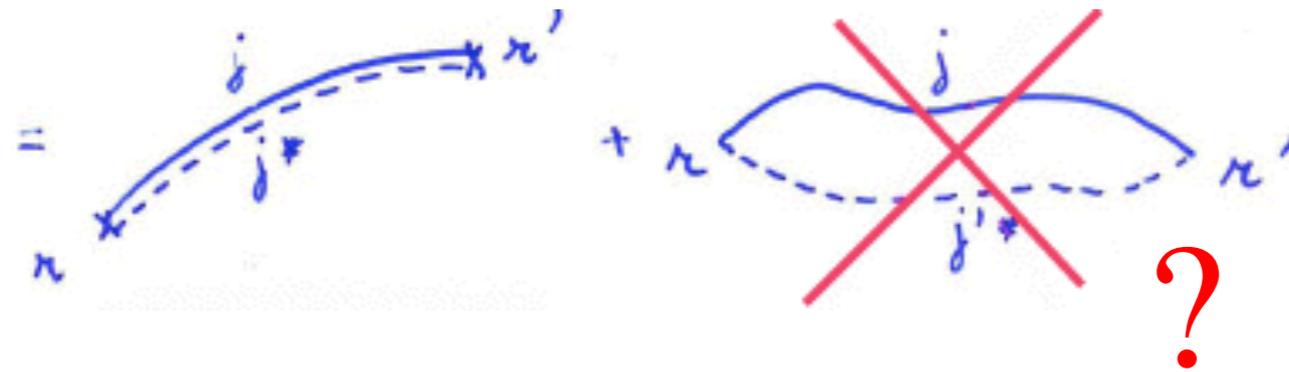
Beyond incoherent
diffusion
(qualitative description)

Coherent effects



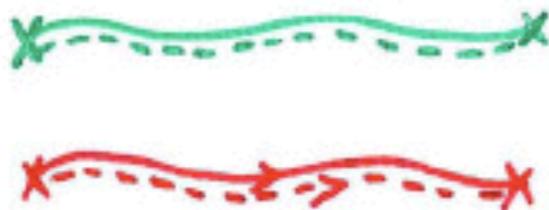
What is the first correction *i.e.*, with the *smallest phase shift* ?

Coherent effects

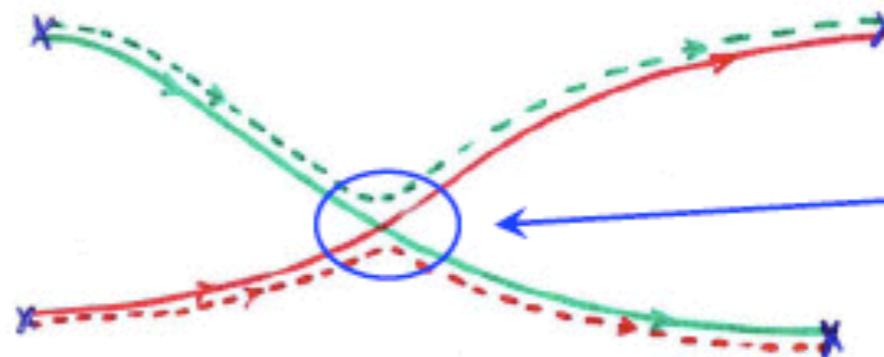


What is the first correction *i.e.*, with the *smallest phase shift* ?
When amplitude paths cross

Example :



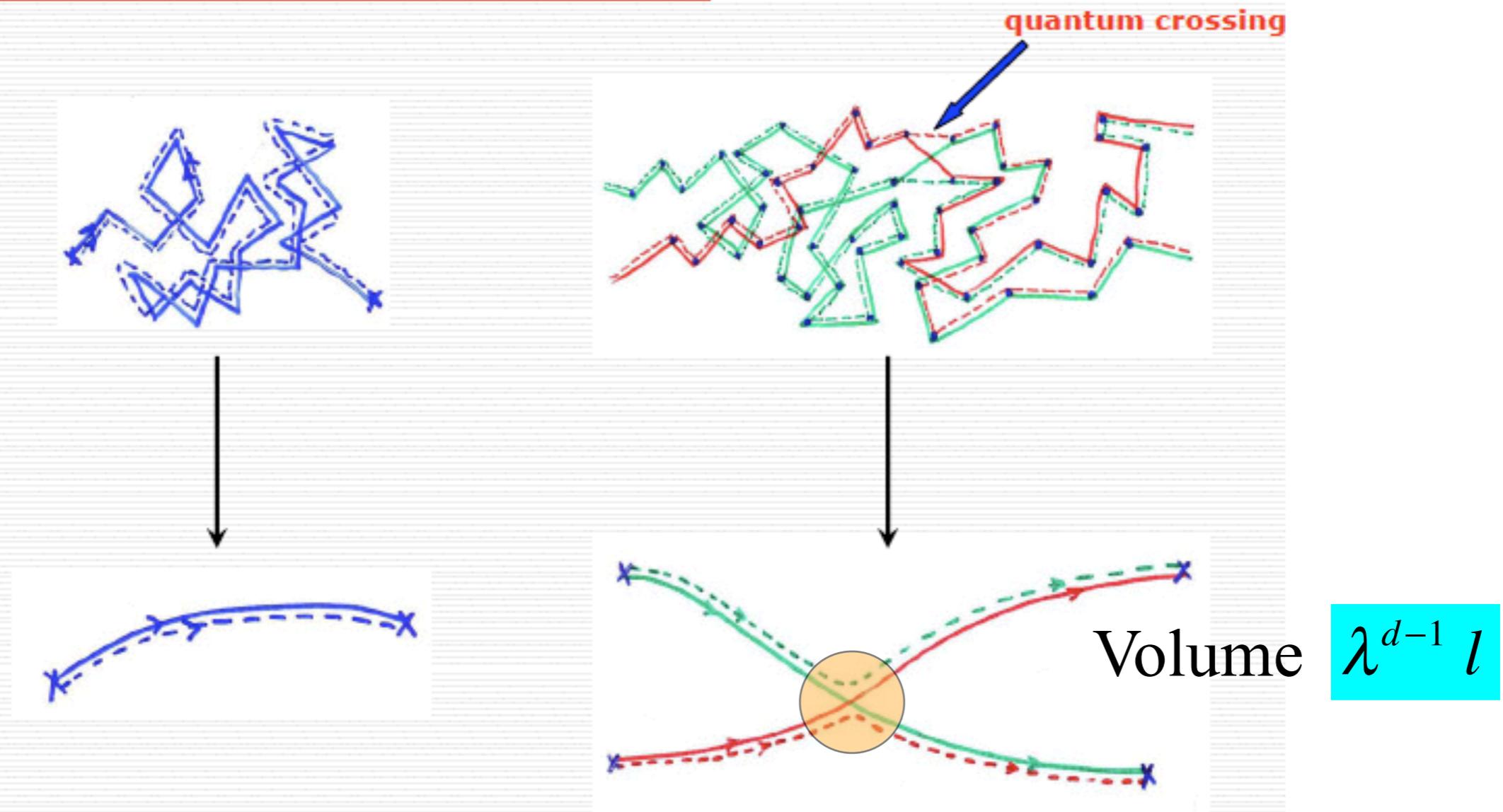
Classical diffusion



quantum crossing

Exchange of amplitudes

Schematic picture



Quantum crossings decrease the diffusion coefficient D :
weak localization

λ : Fermi wavelength

Occurrence of a *quantum crossing* after a time t for a electron diffusing in a volume L^d

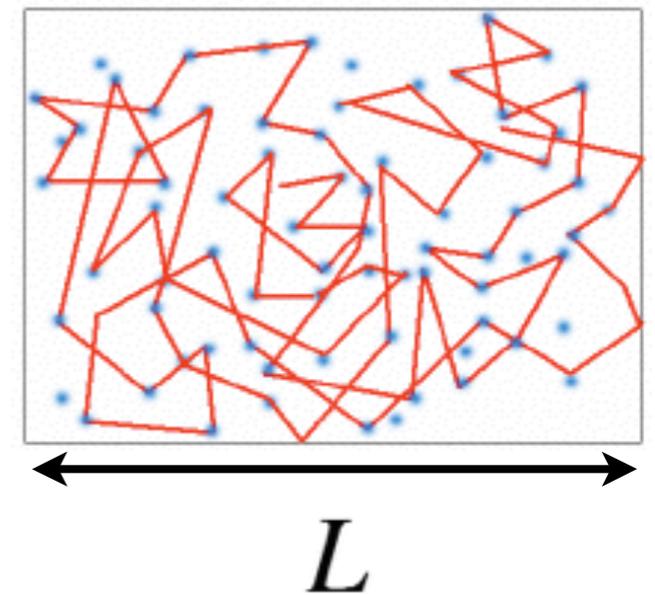
$$p_{\times}(t) = \frac{\lambda_F^{d-1} v_F t}{L^d}$$

v_F : Fermi velocity

The time spent by a diffusing electron is $\tau_D = L^2/D$ so that

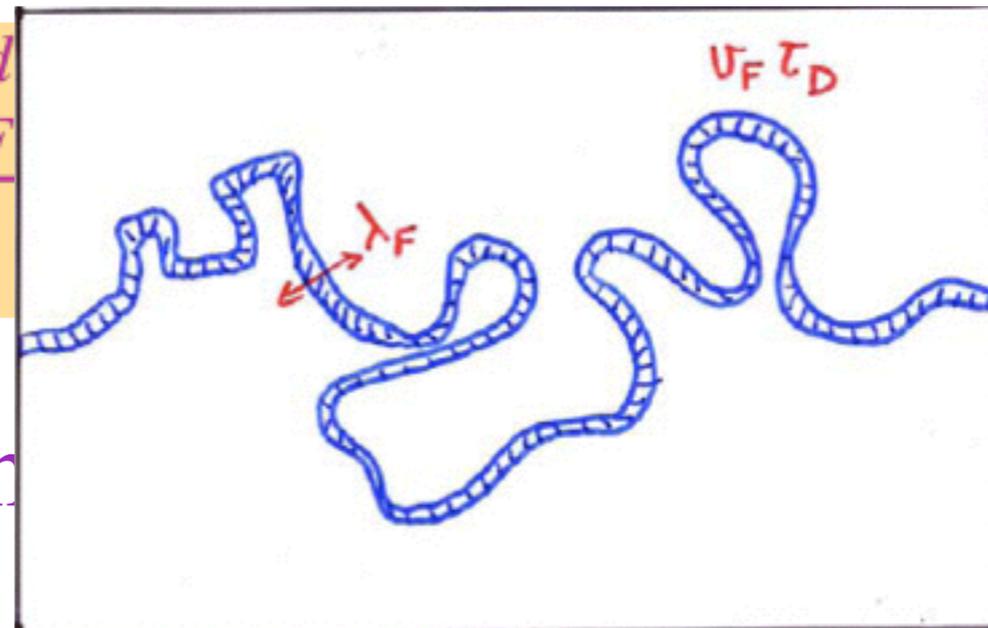
$$p_{\times}(\tau_D) = \frac{\lambda_F^{d-1} v_F \tau_D}{L^d} \equiv \frac{1}{g}$$

$$g = \frac{D}{v_F \lambda_F^{d-1}} L^{d-2}$$



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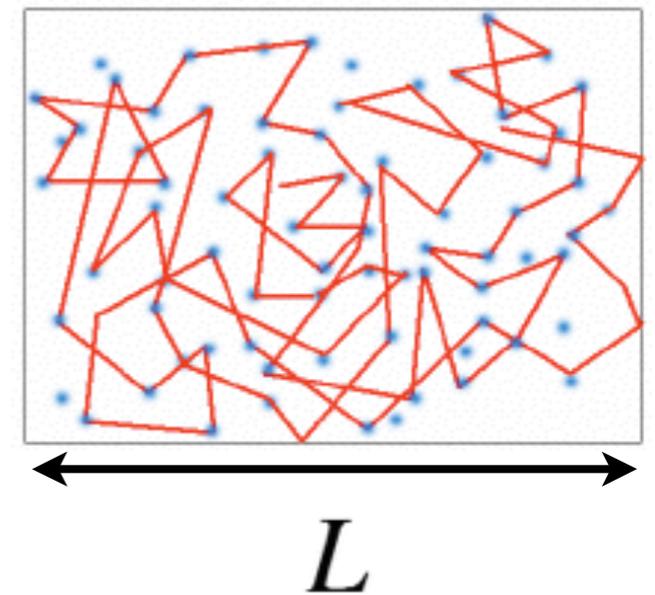
$$p_{\times}(t) = \frac{\lambda_F^d}{V}$$



mi velocity

The time spent by a diffusing electron is $t \sim L^2/D$ so that

$$g \sim \frac{V}{\lambda_F^{d-1} v_F \tau_D}$$



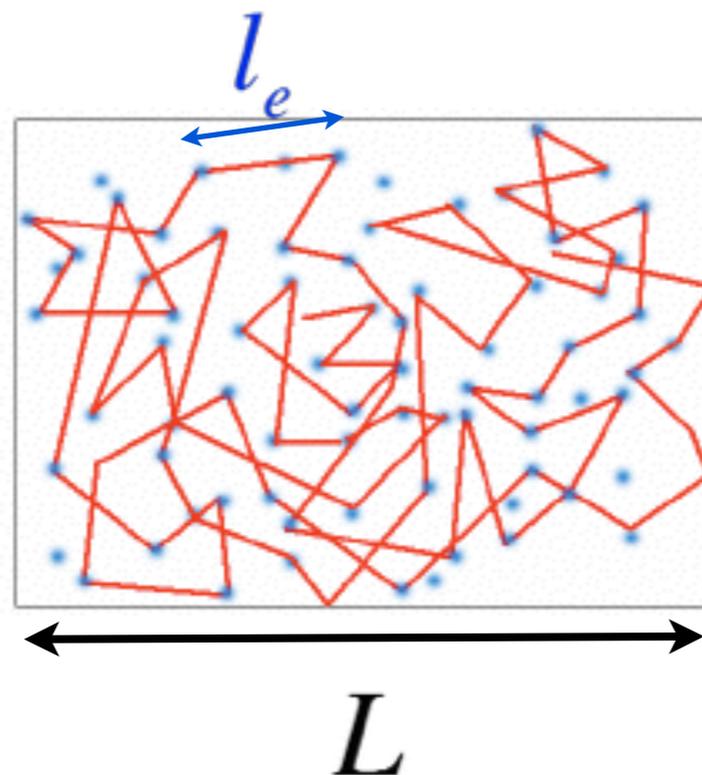
$$g = \frac{D}{v_F \lambda_F^{d-1}} L^{d-2}$$

Physical meaning of this parameter ?

Electrical conductance of a metal

A metal can be modeled as a quantum gas of electrons scattered by an elastic disorder l_e .

Classically, the conductance of a cubic sample of size L^d is given by Ohm's law: $G = \sigma L^{d-2}$ where σ is the conductivity.



Conductance = ratio of two volumes

Conductivity (Drude) $\sigma = \frac{ne^2\tau_e}{m} = e^2 D \rho$

$$D = \frac{v_F l_e}{3}$$

Conductance (Ohm's law) $G = \sigma \frac{S}{L}$

$$S = W^{d-1}$$

Dimensionless conductance $G = g \frac{e^2}{h}$

$$g \sim \frac{D}{v_F \lambda^{d-1}} \frac{W^{d-1}}{L} = \frac{D}{v_F \lambda^{d-1}} \frac{V}{L^2} \sim \frac{V}{\lambda^{d-1} v_F \tau_D}$$

A direct consequence: quantum corrections to electrical transport

Classical transport : $G_{cl} = g \times \frac{e^2}{h}$ with $g \gg 1$

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Quantum corrections: $\Delta G = G_{cl} \times \frac{1}{g}$

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Classical transport : $G_{cl} = g \times \frac{e^2}{h}$ with $g \gg 1$

Quantum corrections: $\Delta G = G_{cl} \times \frac{1}{g}$

so that $\Delta G \simeq \frac{e^2}{h}$ is universal

A direct consequence: quantum corrections to electrical transport

Classical transport : $G_{cl} = g \times \frac{e^2}{h}$ with $g \gg 1$

Quantum corrections: ΔG

What does it mean ?

so that $\Delta G \simeq \frac{e^2}{h}$ is universal

A direct consequence: quantum corrections to electrical transport

Classical transport : $G_{cl} = g \times \frac{e^2}{h}$ with $g \gg 1$

Quantum

Independent of the microscopic
(and often unknown) disorder -
Depends only on the geometry

so that $\Delta G \simeq \frac{e^2}{h}$ is universal

A direct consequence: quantum corrections to electrical transport

Classical $\Delta G_{cl} = g \times \frac{e^2}{h}$ with $g \gg 1$

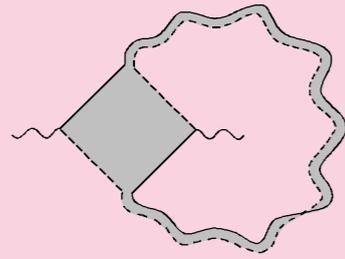
Not that simple!

numbers... Need to sum up Feynman diagrams.

Quantum corrections

so that $\Delta G = \# \frac{e^2}{h}$ is universal

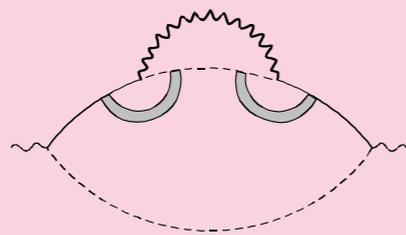
A direct consequence: quantum corrections to electrical



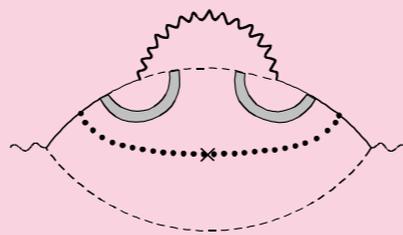
$$\text{Diamond} = \text{Diamond} + \text{Diamond} + \text{Diamond}$$

$$\Delta\sigma^* = \text{Diagram 1} + \text{Diagram 2} + \text{Diagram 3}$$

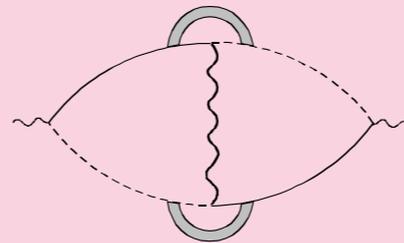
$$c) \quad \text{Vertex} = \text{Wavy Line}$$



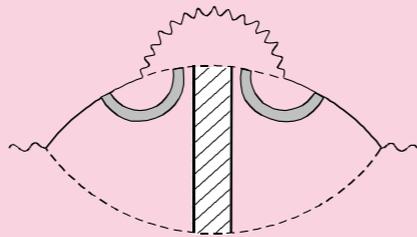
(a)



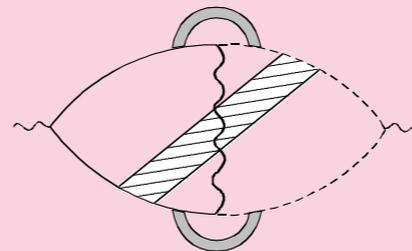
(b)



(c)



(d)

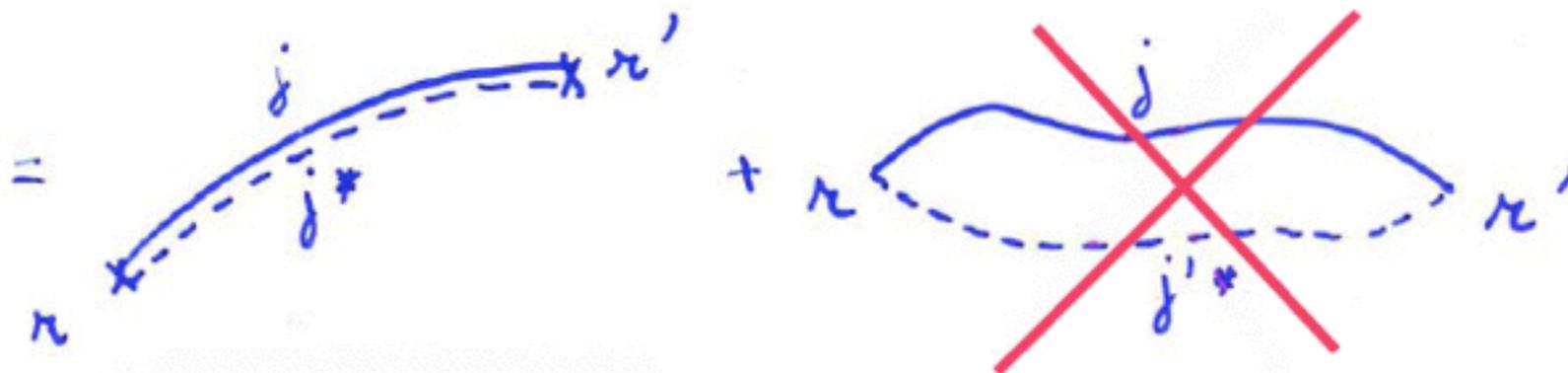


(e)

Summary

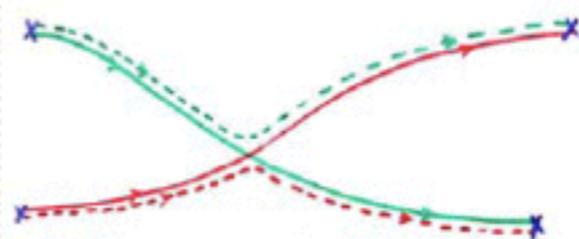
conductance \sim transmission \sim probability

$$P_{cl}(r, r')$$



classical diffusion

quantum corrections



quantum crossing \rightarrow $1/g$ correction

$$P_x(\tau_D) \sim \frac{\lambda_F^{d-1} v_F \tau_D}{V} \sim \frac{1}{g}$$

classical transport $\propto g \frac{e^2}{h}$

quantum effects $\propto \frac{e^2}{h}$

Complexity of a quantum mesoscopic system

Elastic disorder does not break phase coherence
and it does not introduce irreversibility

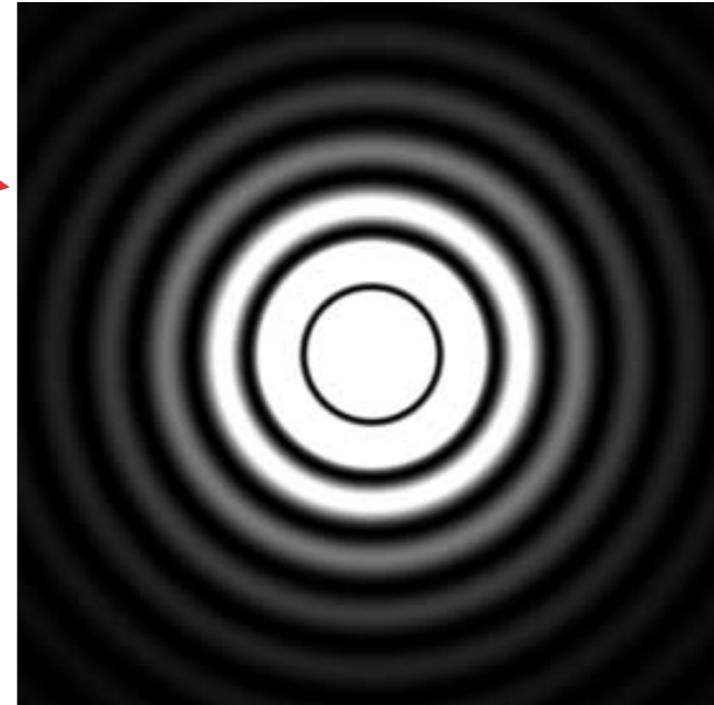
Elastic disorder does not break phase coherence
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Disorder introduces randomness and
complexity:

All symmetries are lost, there are no good
quantum numbers.

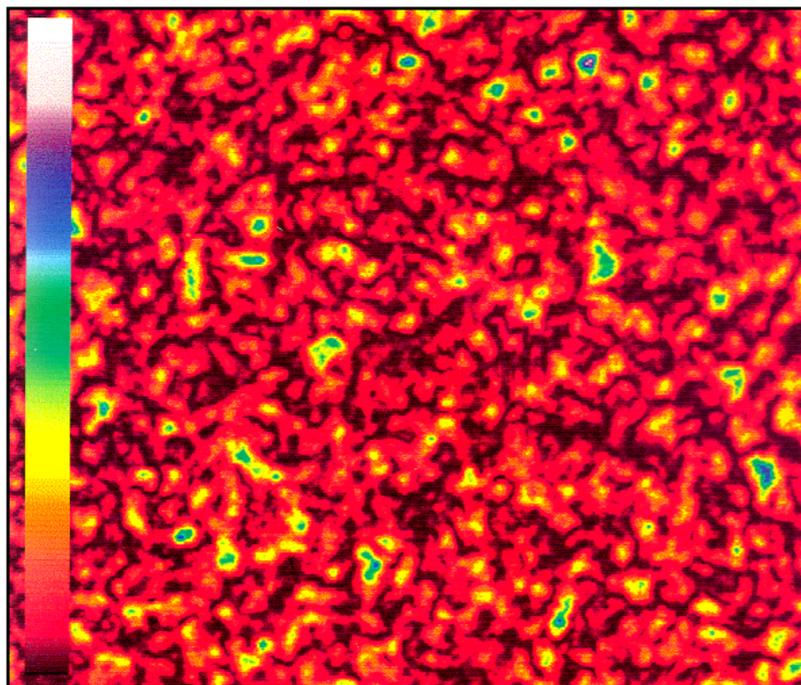
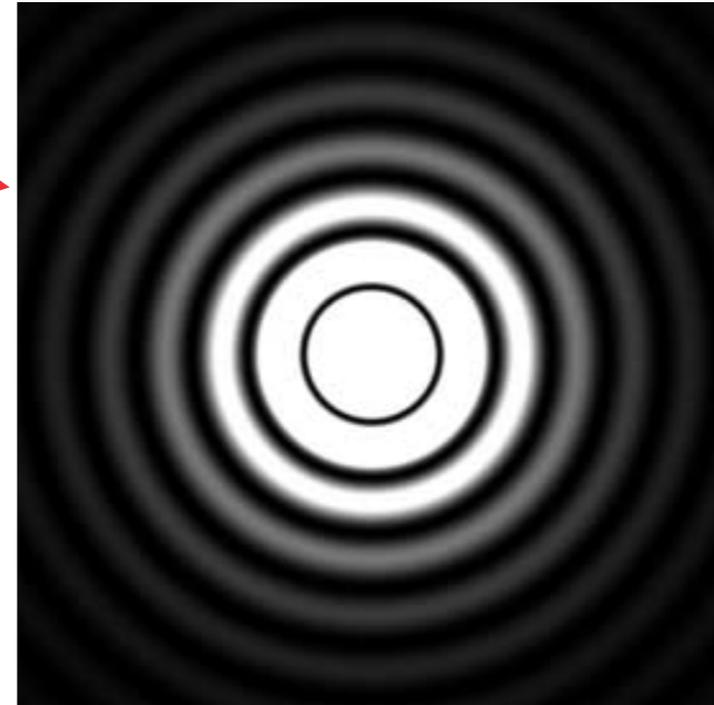
Example: speckle patterns in optics

Diffraction
through a circular
aperture: order in
interference



Example: speckle patterns in optics

Diffraction
through a circular
aperture: order in
interference



Transmission of
light through a
disordered
suspension:
complex system

Mesoscopic quantum systems

- Most (all ?) quantum systems are complex
- Complexity (randomness) and decoherence are separate and independent notions.
- Complexity
quantum numbers)
- Decoherence
coherence $L \gg L_\varphi$

A mesoscopic quantum system is a coherent complex quantum system with $L \leq L_\varphi$

Phase coherence and self-averaging: universal fluctuations.

Classical limit : $L \gg L_\varphi$

The system is a collection of $N = (L/L_\varphi)^d \gg 1$
statistically independent subsystems.

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A macroscopic observable defined in each subsystem
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Classical limit : $L \gg L_\varphi$

The system is a collection of $N = (L/L_\varphi)^d \gg 1$ statistically independent subsystems.

A macroscopic observable defined in each subsystem takes independent random values in each of the N pieces.

Law of large numbers: any macroscopic observable is equal with probability one to its average value.

The system performs an
average over realizations of
the disorder.

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For $L \ll L_\varphi$, we expect deviations from self-averaging which reflect the underlying quantum coherence.

Quantum conductance fluctuations

Classical self-averaging limit : $\frac{\delta G}{\overline{G}} = \frac{1}{N} = \left(\frac{L_\varphi}{L}\right)^{d/2}$

where $\delta G = \sqrt{\overline{G^2} - \overline{G}^2}$ and $\overline{G} = \sigma L^{d-2}$

$\overline{\quad}$ is the average over disorder.

$$\delta G^2 \propto L^{d-4}$$

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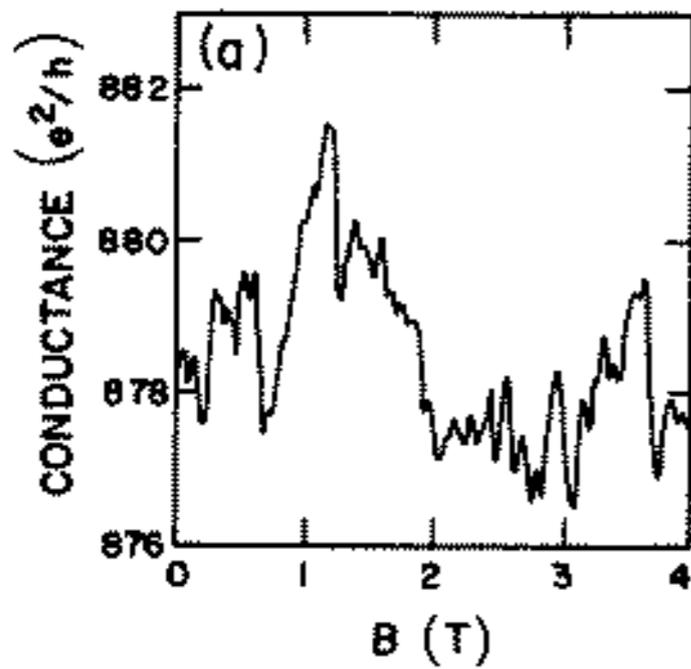
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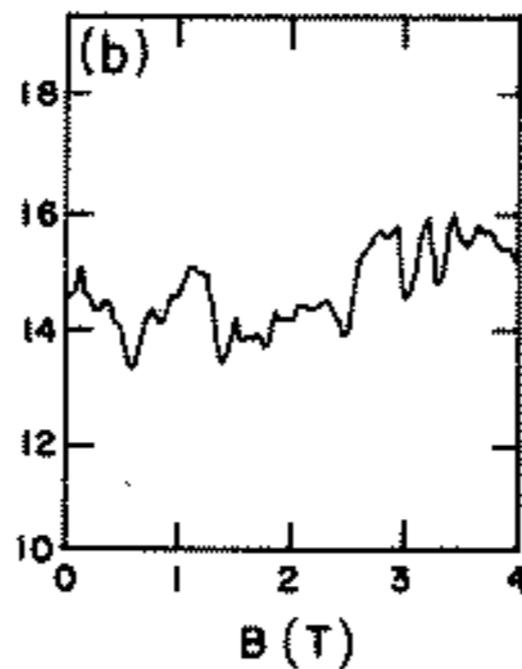
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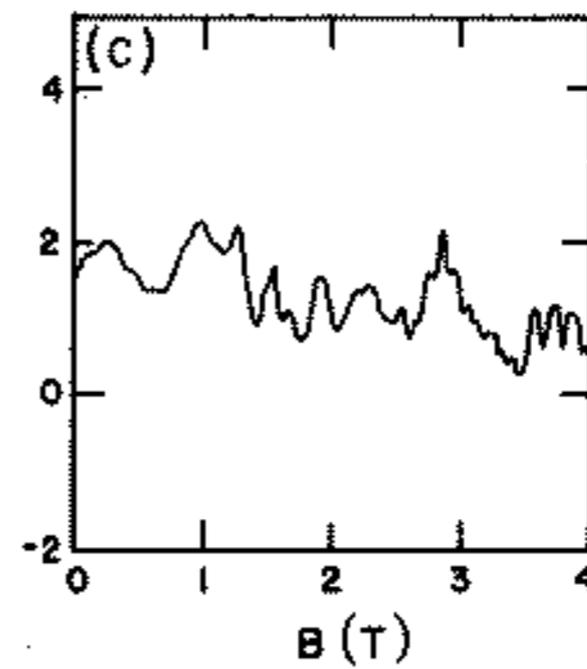
In the mesoscopic limit, the electrical conductance is not self-averaging.



Gold ring



Si-MOSFET



NUMERICS ON
THE ANDERSON MODEL

Summary : key ideas and concepts in quantum mesoscopic physics

1. Classical diffusion

2.

3.

crossings) are propagated over long distances by means of classical diffusion.

3. Complexity

from decoherence.