



Conductance spectra of NbN-Bi₂Se₃-Au junctions: Tunneling gaps & ZBCPs

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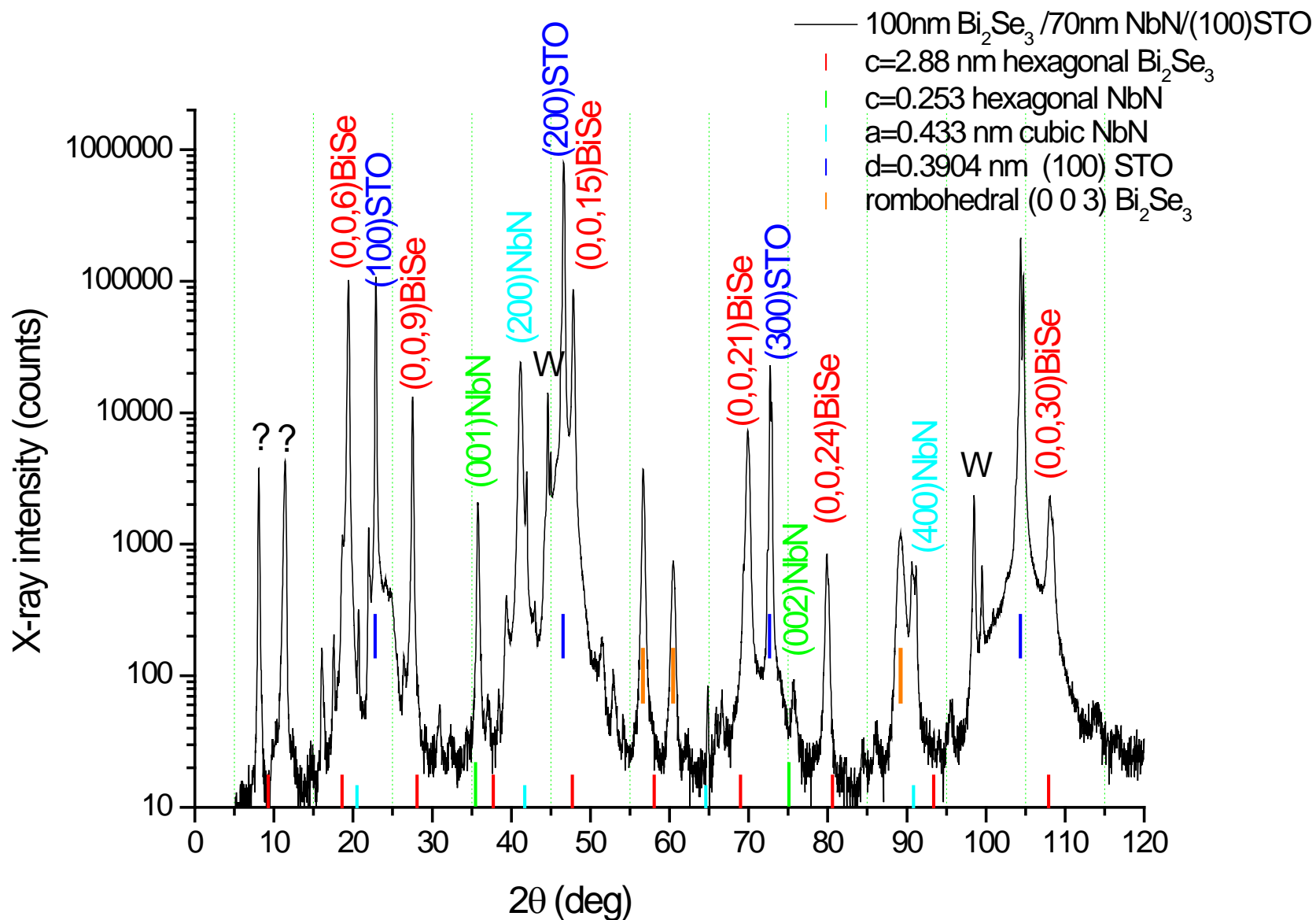
<http://lanl.arxiv.org/abs/1207.5352>

Some basics & what is done here

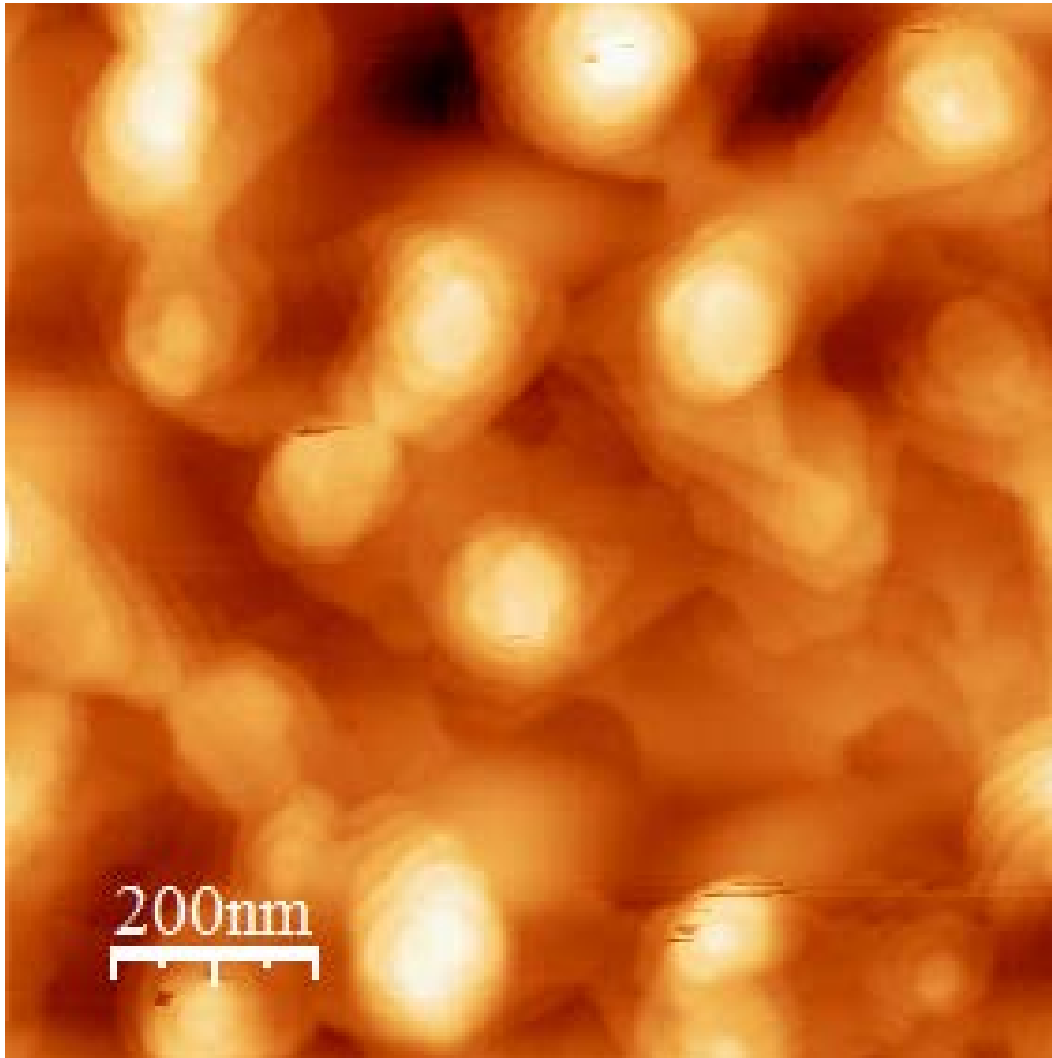
- Topological Insulators (TOI) are band-gap insulators with conduction by surface states due to a strong spin-orbit interaction
- The hallmark of Topological Superconductivity is the presence of Majorana fermions (MF), which are a special sub-group of surface Andreev bound states (SABS) at zero energy
In the experiments, they will appear as ZBCPs
- We use **proximity induced superconductivity in a TOI** in NbN-Bi₂Se₃-Au junctions, similar to Yang & Li Lu, PE in Sn/Bi₂Se₃, PRB **85**, 104508 (2012)
& G. Koren et al., PE in Bi/Bi₂Se₃, PRB **84**, 224521 (2011)
- We characterize the films structure & growth orientation, & measure transport (R vs T), and conductance spectra vs H & T
- The observed ZBCPs are ~0.15mV & ~0.9mV wide (FWHM). The narrower one is kT limited (~2K) & the broader one might be a result of thermally smeared split ZBCP (Yamakage & Tanaka, PRB **85**, 180509(R) (2012) or other non zero bound states

X-ray diffraction of Bi_2Se_3 / NbN bilayer on (100)STO

Shows hexagonal Bi_2Se_3 & mostly cubic but some hexagonal NbN



An AFM image of a
100nm thick Bi_2Se_3 on 70nm thick NbN bilayer on (100) SrTiO_3



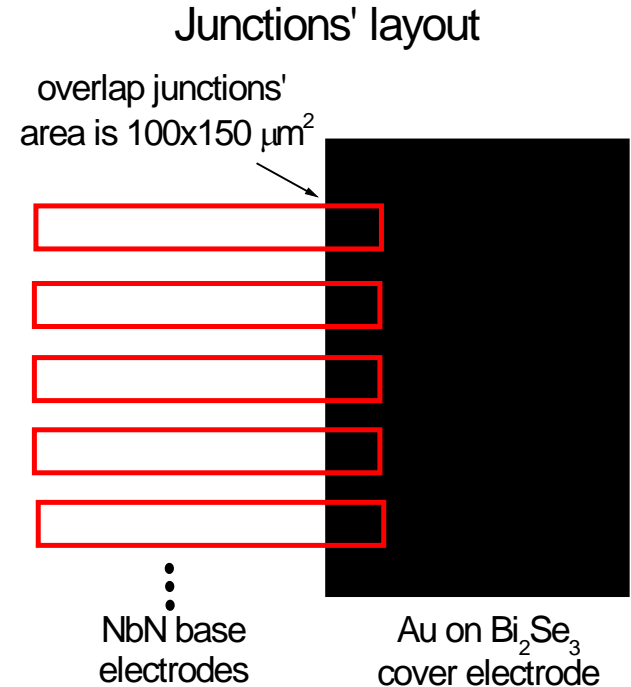
Crystallized well,
in laterally disordered
hexagonal form

$1 \times 1 \mu\text{m}^2$

Junctions layout

- 10 base electrodes were patterned first in a NbN film on half the wafer
- Then a cover bilayer of Au/Bi₂Se₃ was deposited on the other half of the wafer with the aid of a **shadow mask**
- No further patterning was done, only Ag paste was added on the NbN contact pads

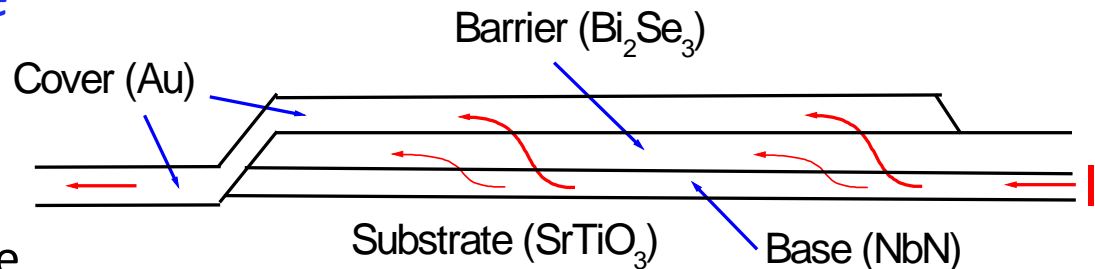
Junctions with *ex-situ* interface



Junction with *in-situ* interface

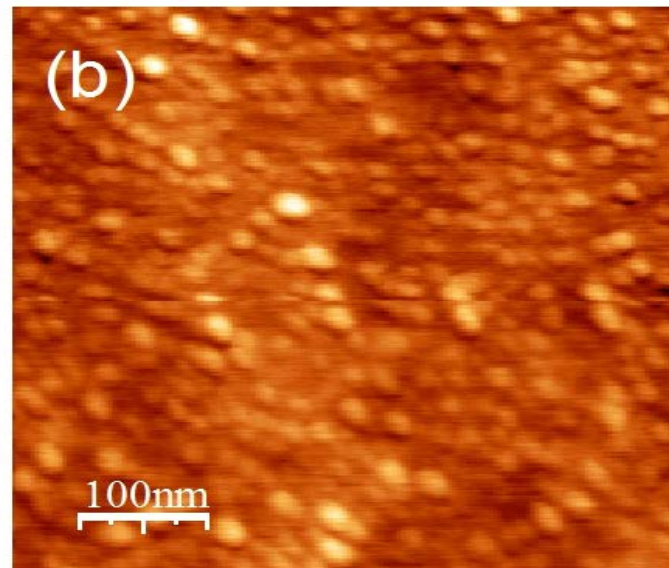
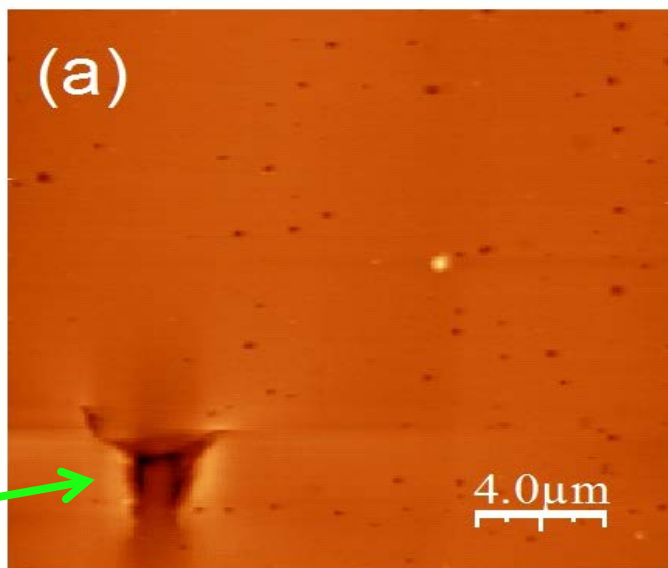
(Full patterning)

- Start with a Bi₂Se₃/NbN bilayer (base)
- Then add the Au cover electrode



AFM images

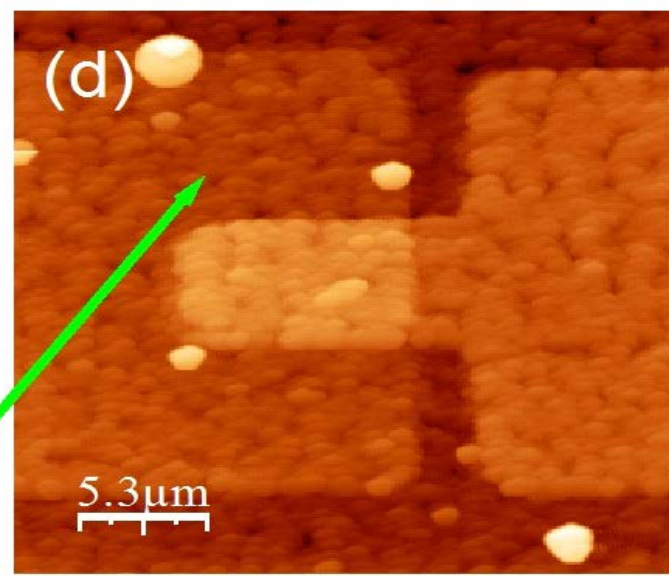
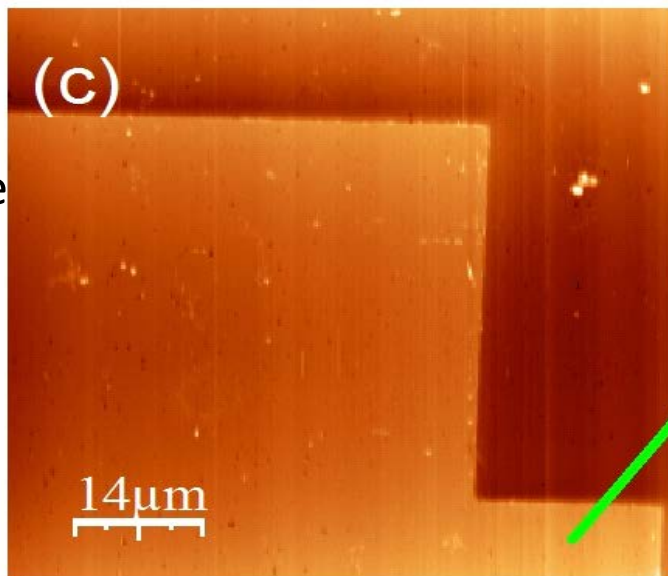
NbN films



As deposited
virgin films (a & b)

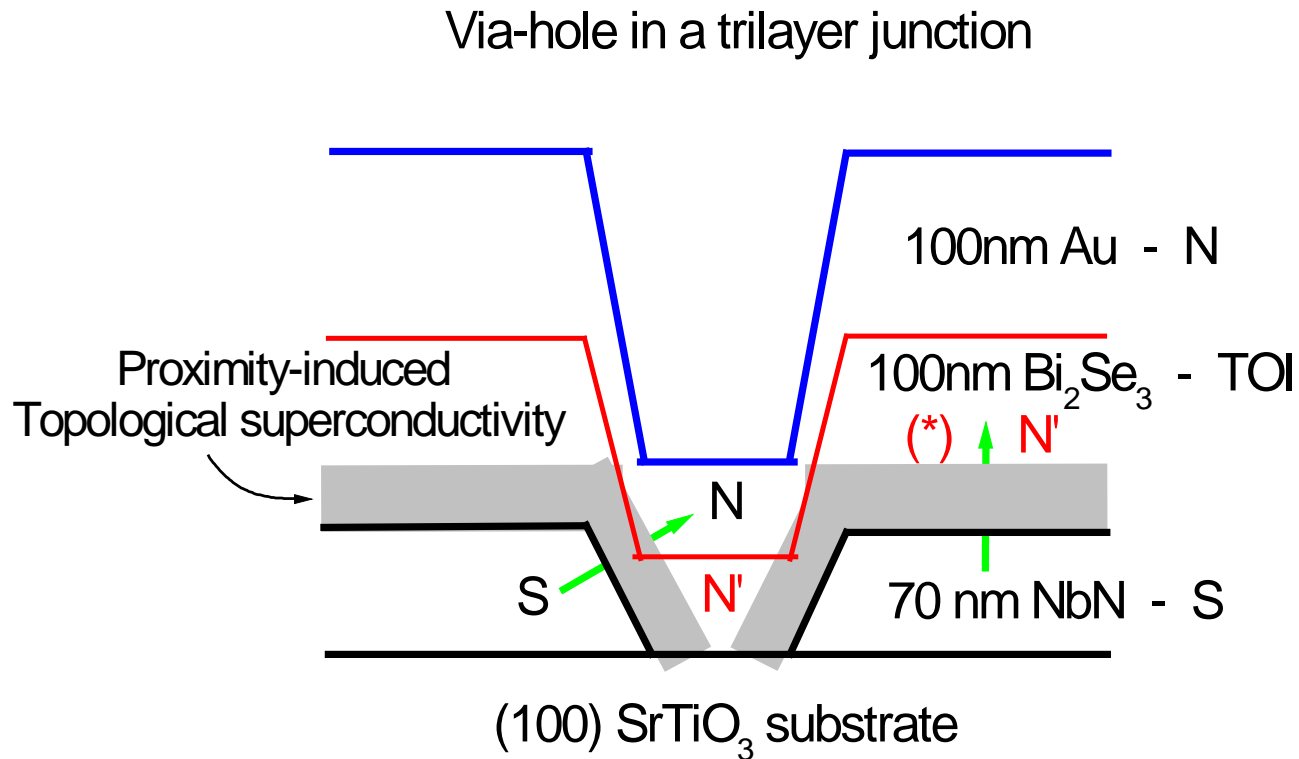
A large via-hole

After patterning (c)
of the base electrode



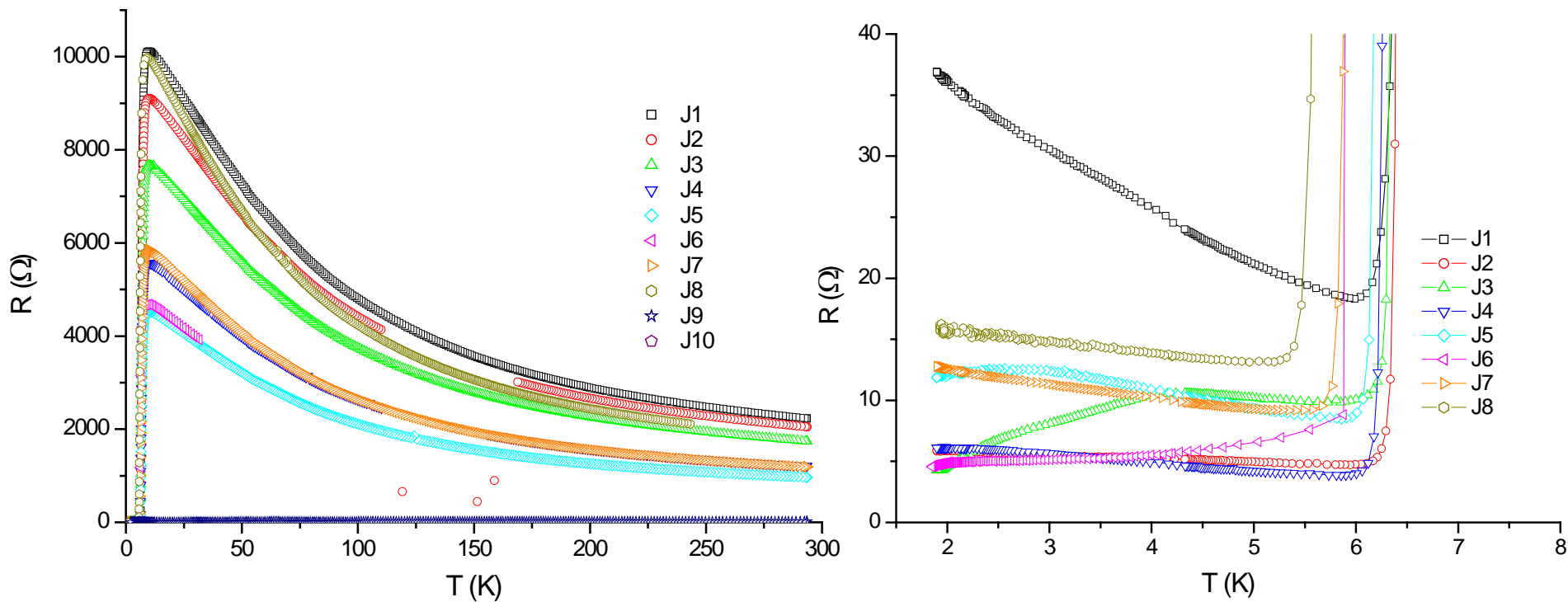
(d) A small junction
with full patterning

A schematic cross section near a via hole



(*) Since the Bi₂Se₃ film is heavily hole doped by Se vacancies, it can also be considered as N'

R vs T of a 100x150 μm^2 area *ex-situ* NbN-Bi₂Se₃-Au junctions



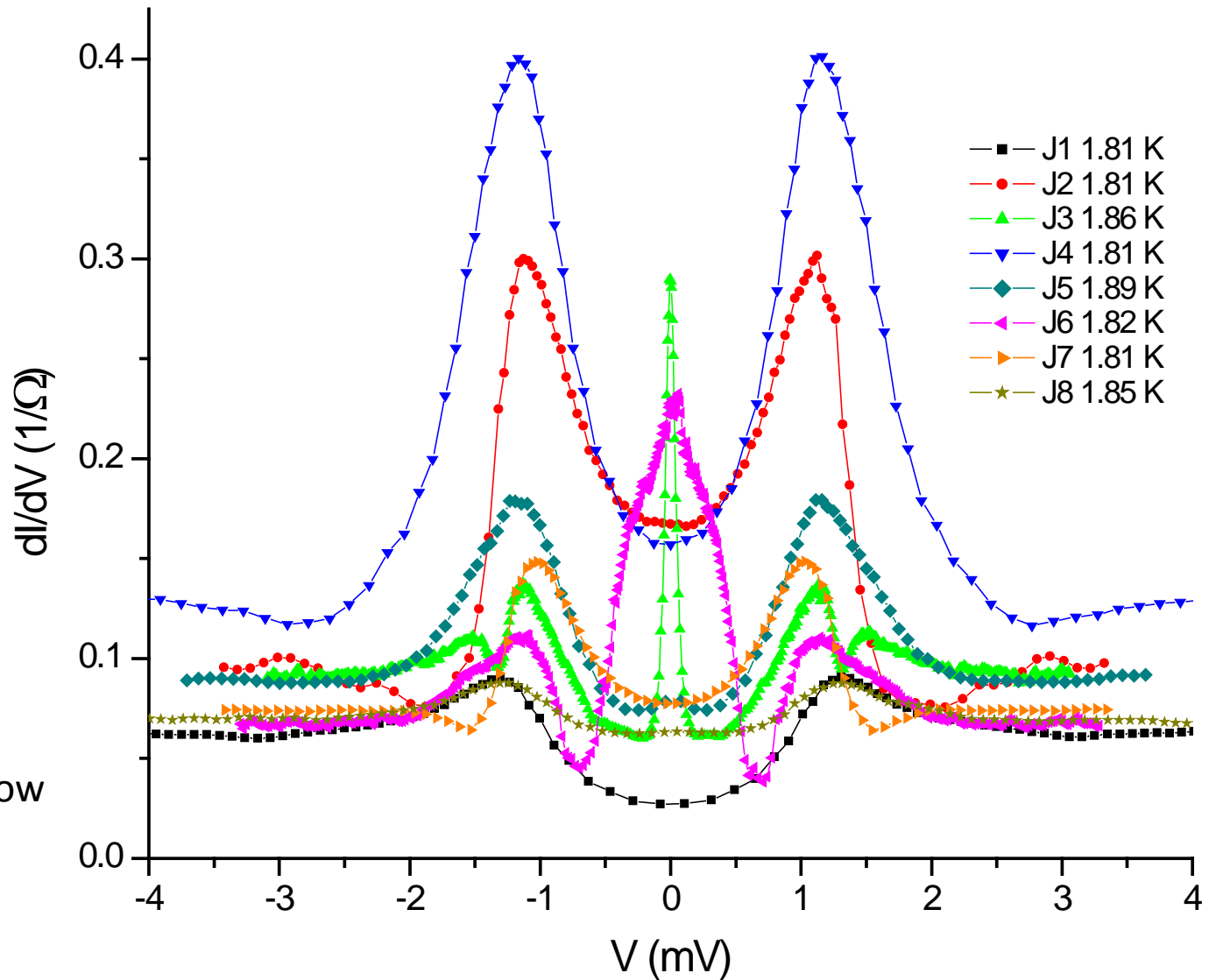
- Insulating vs T due to the low T_c NbN film (deposited in excess N_2)
- Superconductive transition only at $T_c \sim 8\text{K}$ (onset)
- Junctions' resistance at low T is of about 10 Ω
- Junctions with a metallic behavior at low T have ZBCPs (#3, #5 & #6)

Conductance spectra of all working *ex-situ* junctions

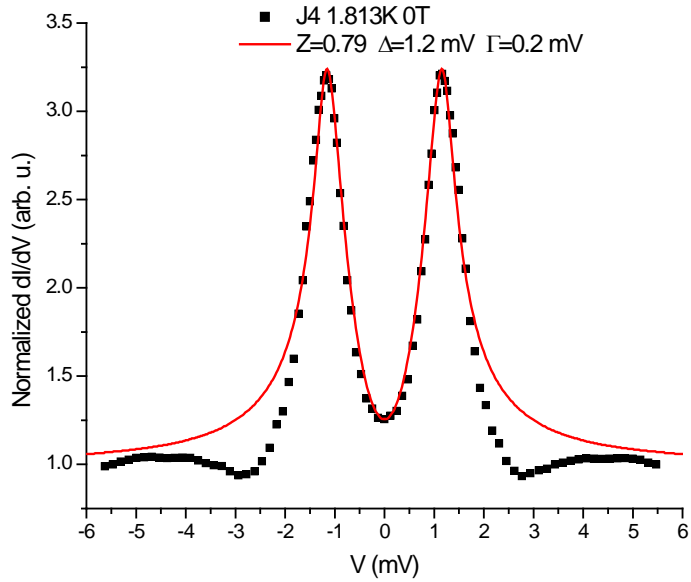
All junctions have coherence peaks of the NbN SC electrode at 1-1.2 mV, consistent with the BCS ratio: $2\Delta/kT_c \sim 3.5$

Narrow (kT -limited) and broad ZBCPs

The broad one might be a thermally smeared split one, or with a narrow ZBCP superimposed



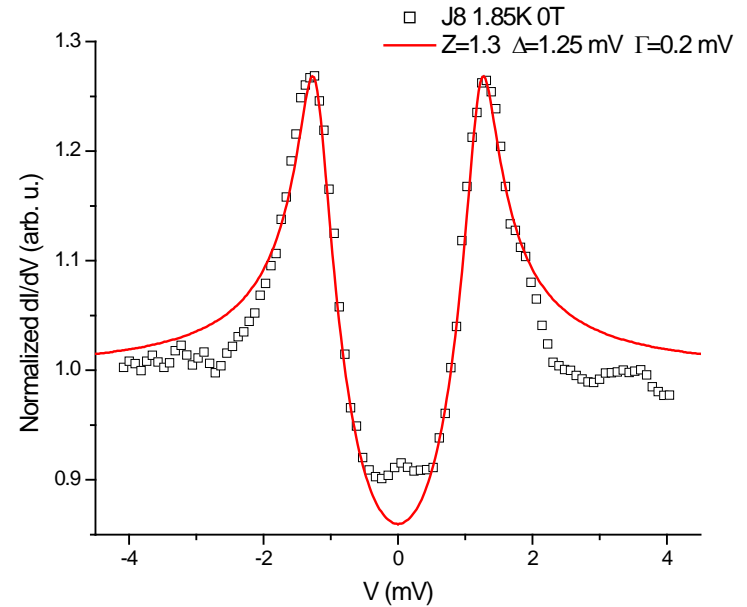
Conductance spectra with s-wave BTK fits of a few NbN-Au/Bi₂Se₃ junctions



• Andreev



• tunneling



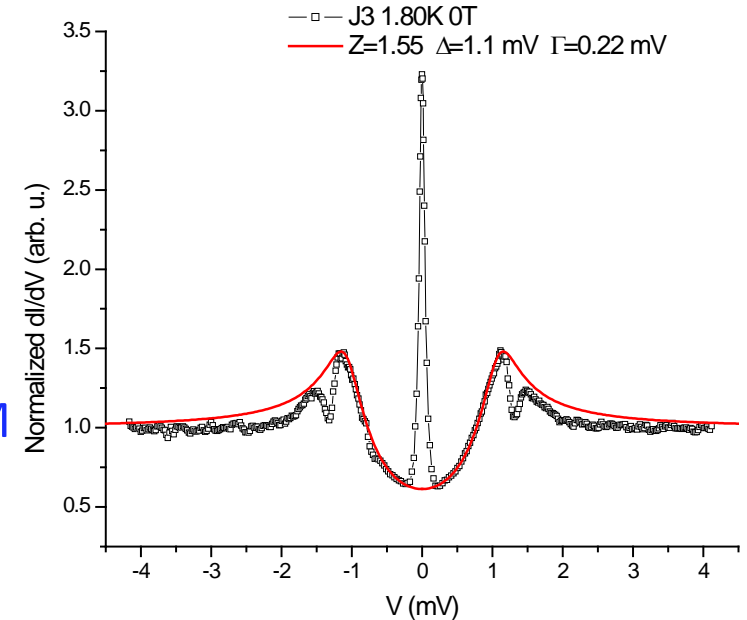
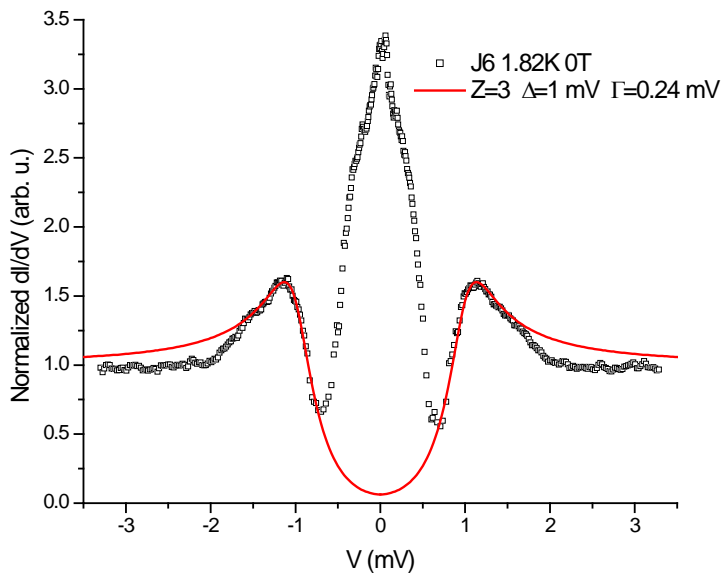
• Broad ZBCP

0.9mV FWHM



• Narrow ZBCP

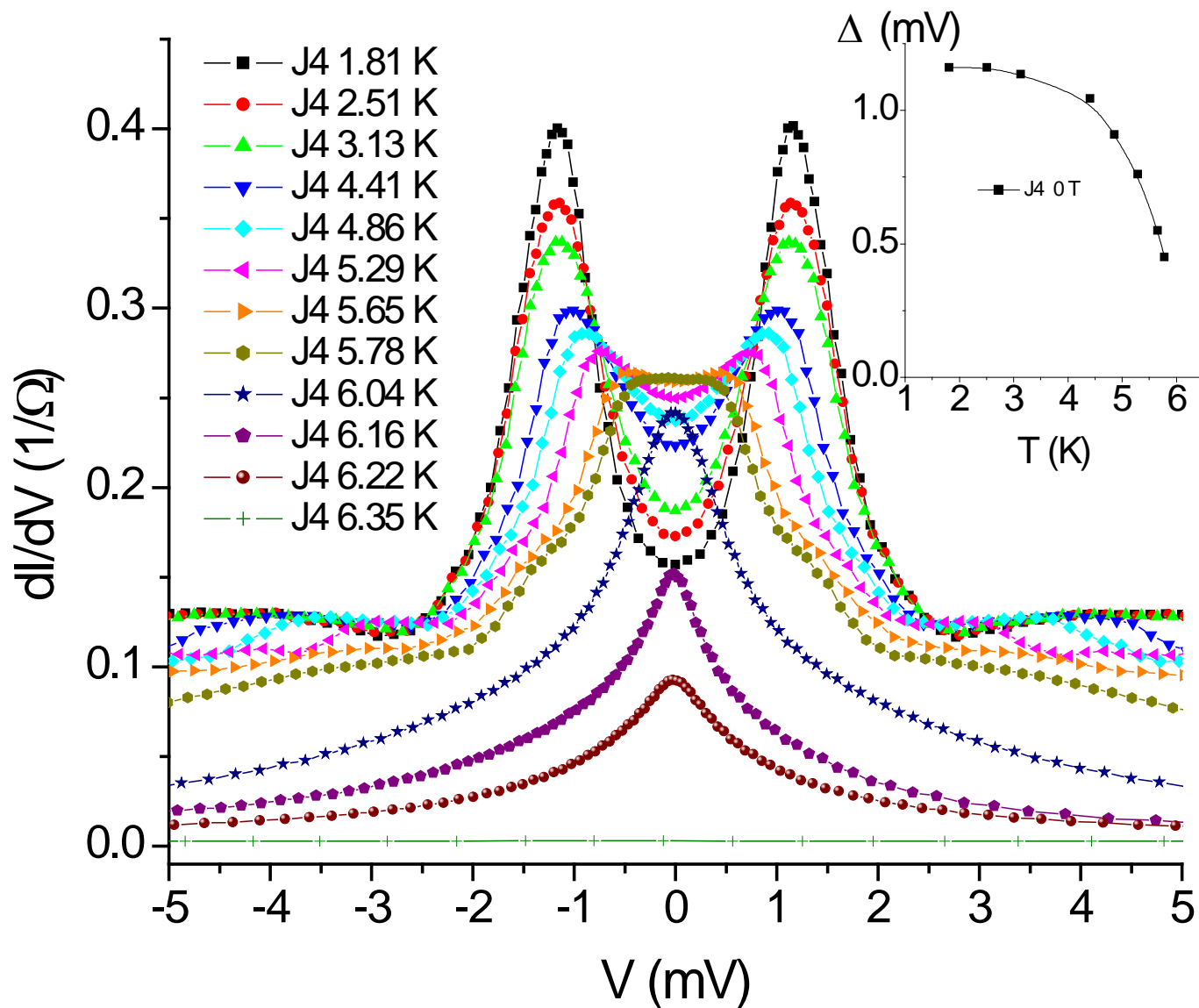
0.15mV FWHM



Conductance spectra at various temperatures

Coherence peaks
(CP) & SC gap
are suppressed
vs T (as usual)

BCS-like gap Δ

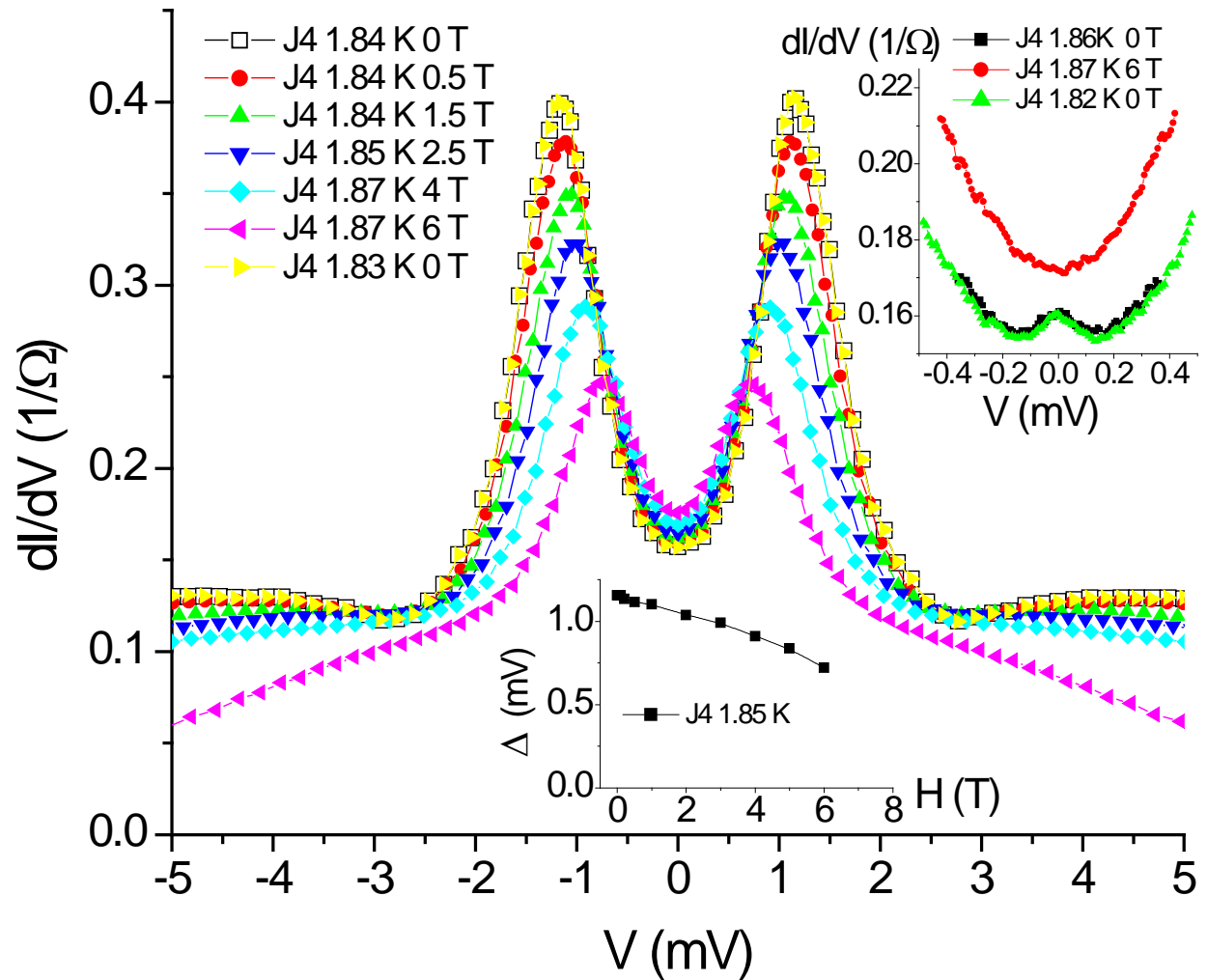


Conductance spectra under various magnetic fields

Coherence peaks
& SC gap are
suppressed
vs H (also as usual)

A small ZBCP at
Low T, suppressed
under 6 T, and
recovers when
field is removed

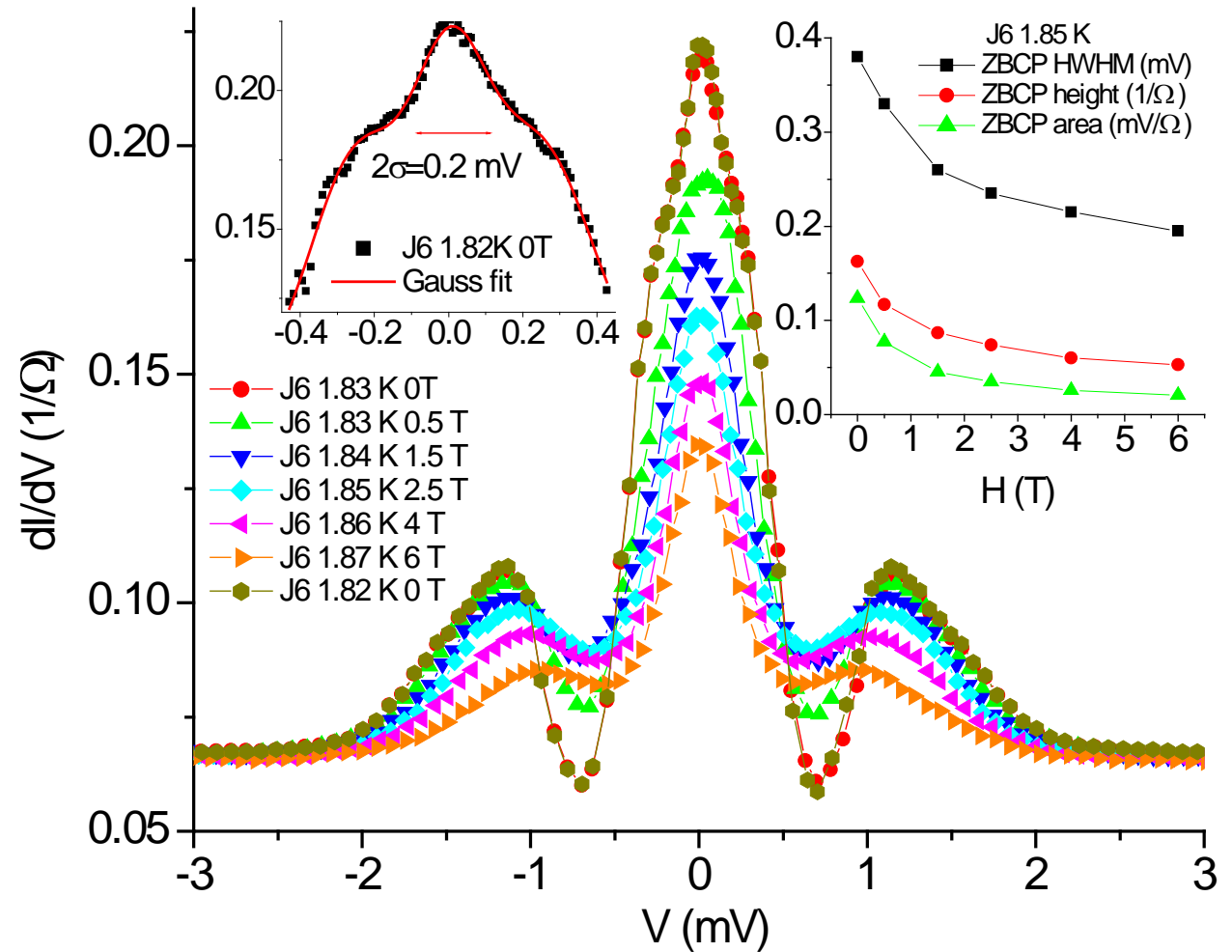
Need higher fields
to test BCS behavior
of the gap Δ



Conductance spectra at various T, CP & broad ZBCP

Coherence peaks,
SC Gap and ZBCP
height, width & area
are suppressed
vs H

Broad ZBCP can be
due to a split ZBCP,
with superimposed
narrow ZBCP
Yamakage & Tanaka,
PRB **85**, 180509(R) (2012)



Conductance spectra at various T, narrow ZBCP

Narrow ZBCP with kT limited width (thermal smearing at 1.8 K)

Its' ZBCP height is suppressed differently than the height of the broad ZBCP

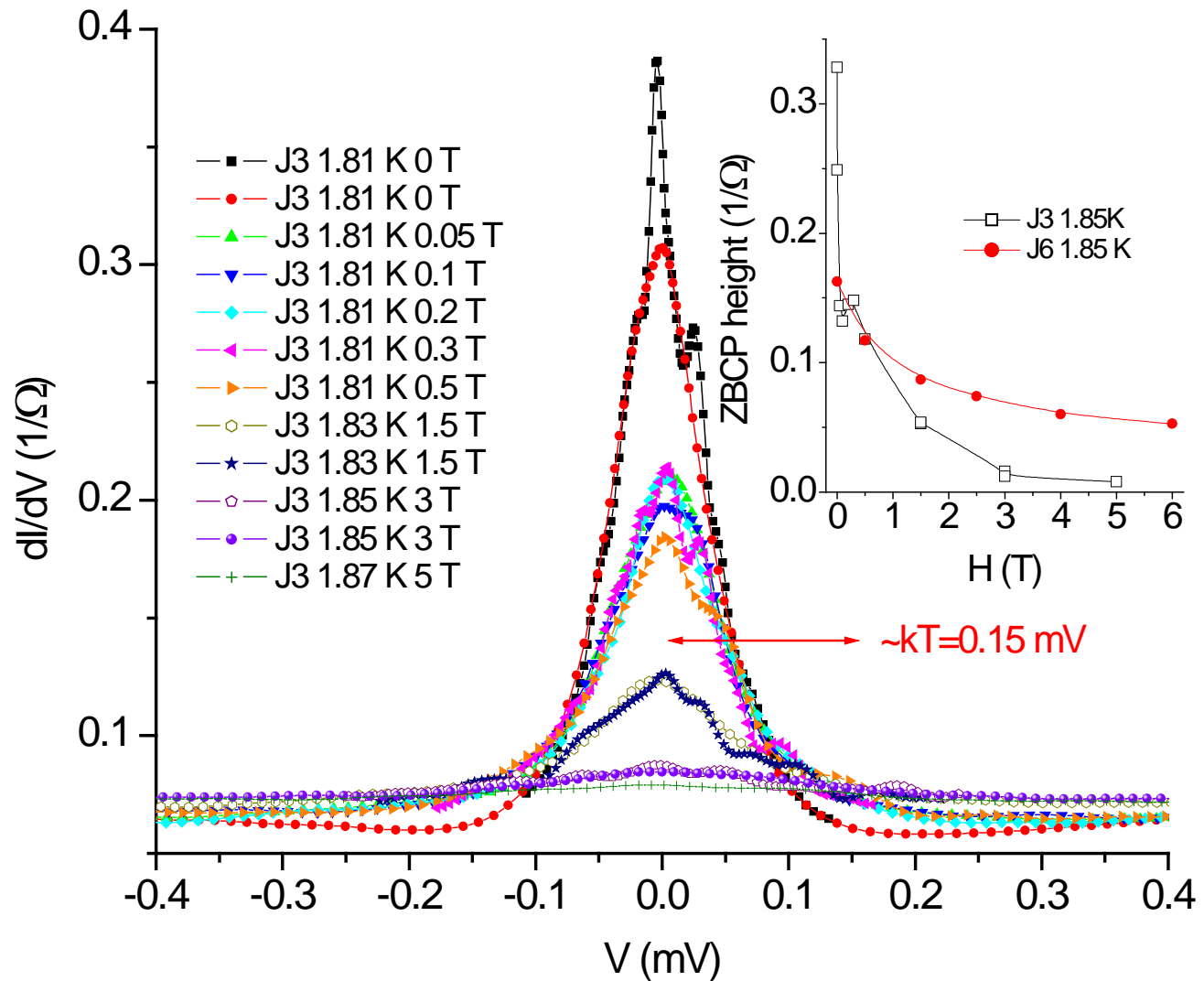
Can be due to Majorana fermions

Yang & Li Lu PRB 85, 104508 (2012)

Mourik et al., Science 336, 1003 (2012).

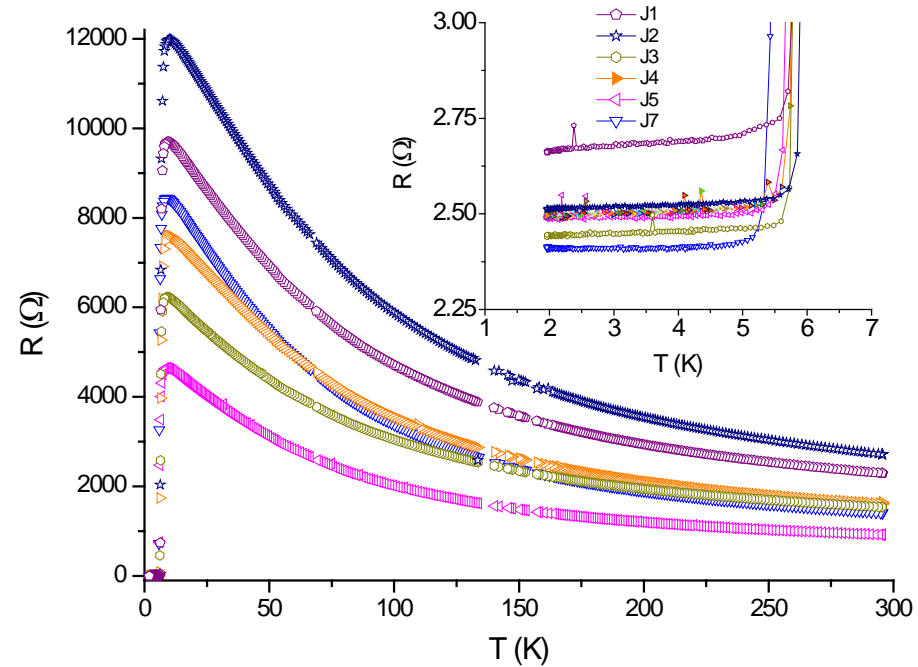
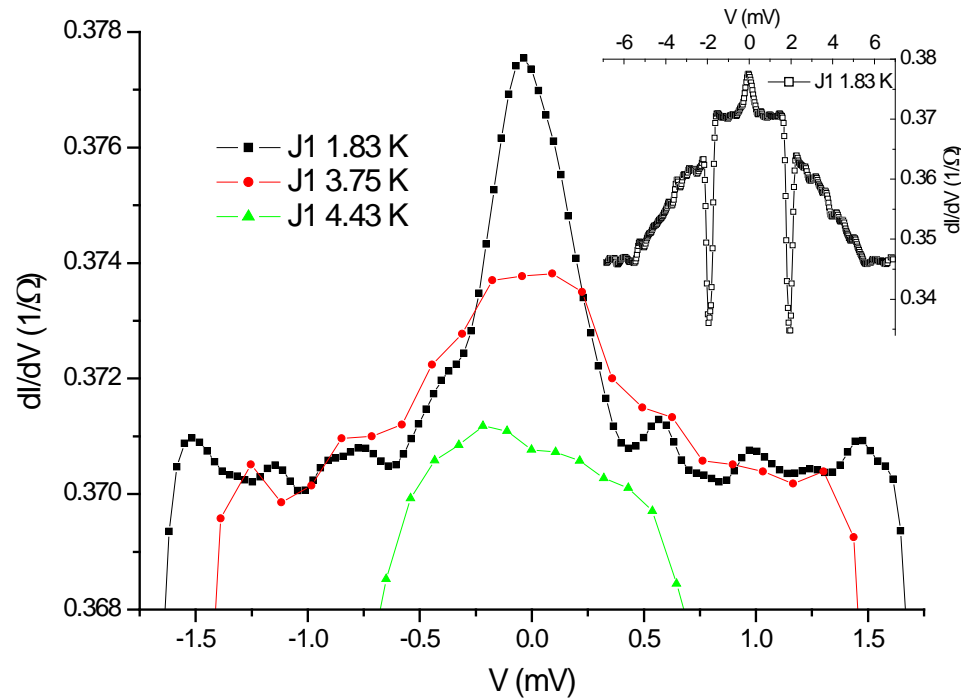
Das et al., arXiv:1205.7073

0.03-0.05 mV widths at 0.1K



Smaller $10 \times 5 \mu\text{m}^2$ area *ex-situ* NbN-Bi₂Se₃-Au junctions

R(junction) went down!
Se vacancies doping by milling.
⇒ Get a low R shunt resistance,
Strongly shunted junction.



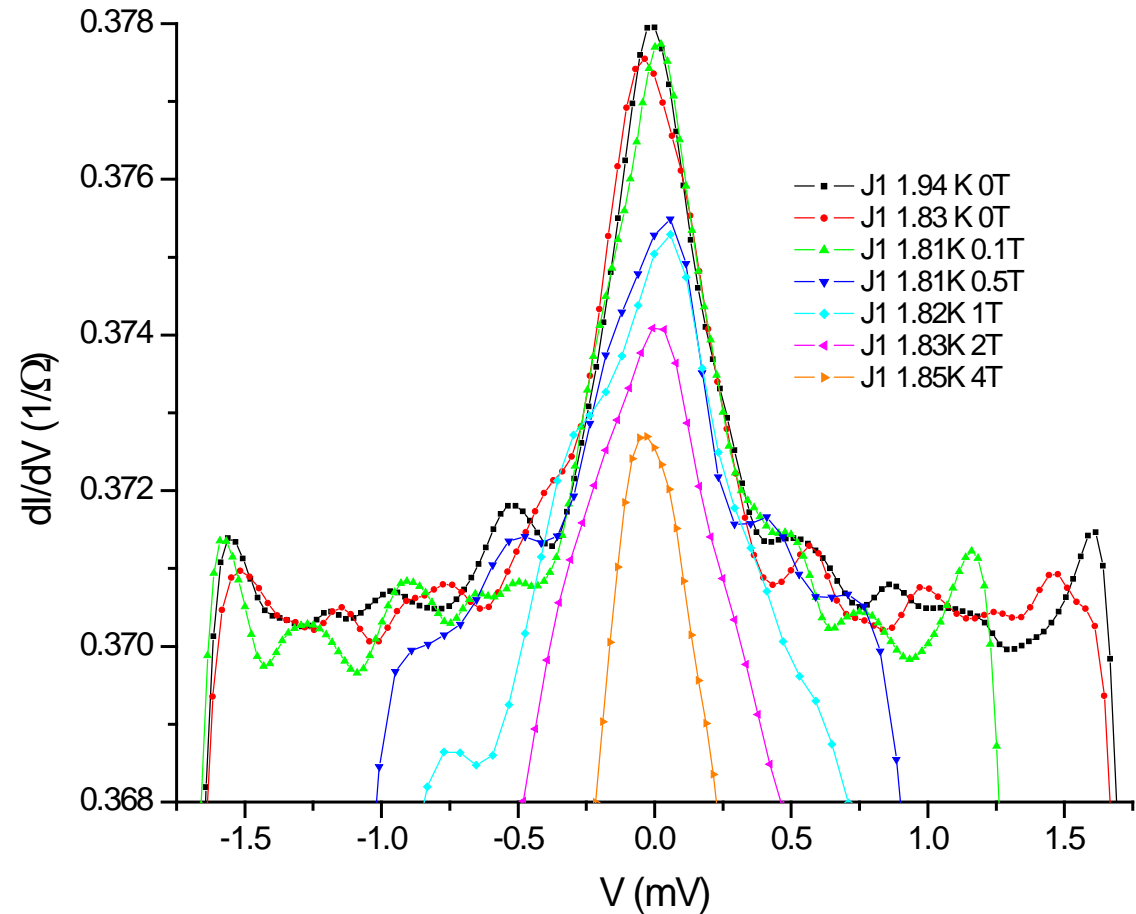
One weak ZBCP survived

No coherence peaks now
(must be a milling effect)

10x5 μm^2 area *ex-situ* NbN-Bi₂Se₃-Au junctions under different fields

The ZBCP is suppressed with increasing magnetic field

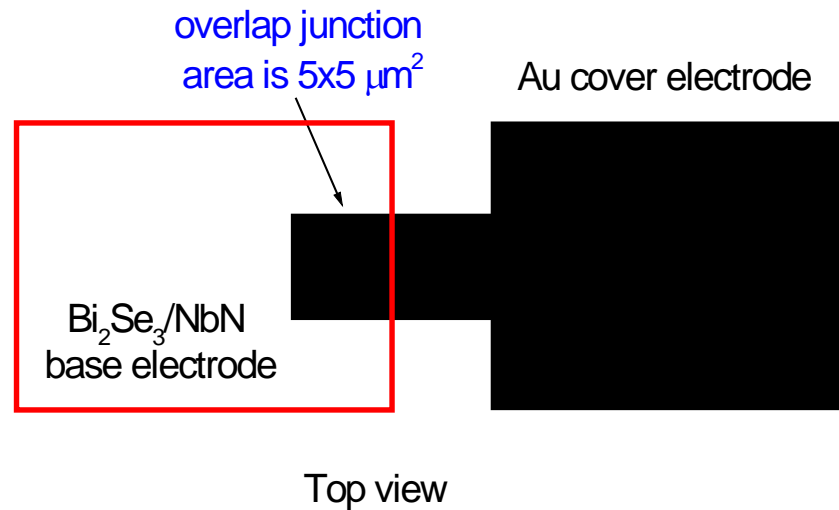
Also the critical current goes down with increasing field (the “dips”)



Layout of a fully lithographically-patterned, *in-situ* junction

- 10 base electrodes are patterned first in a $\text{Bi}_2\text{Se}_3/\text{NbN}$ bilayer on half the wafer
- Then a cover gold film is deposited on the whole wafer, followed by further patterning of this cover electrode
- The resulting junction area is $A \sim 5 \times 5 \mu\text{m}^2$

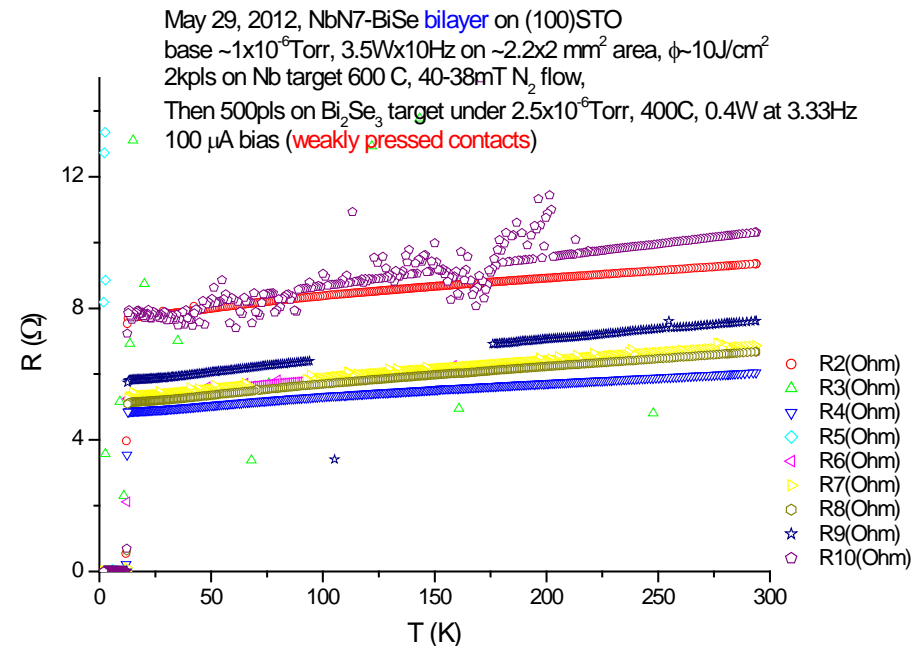
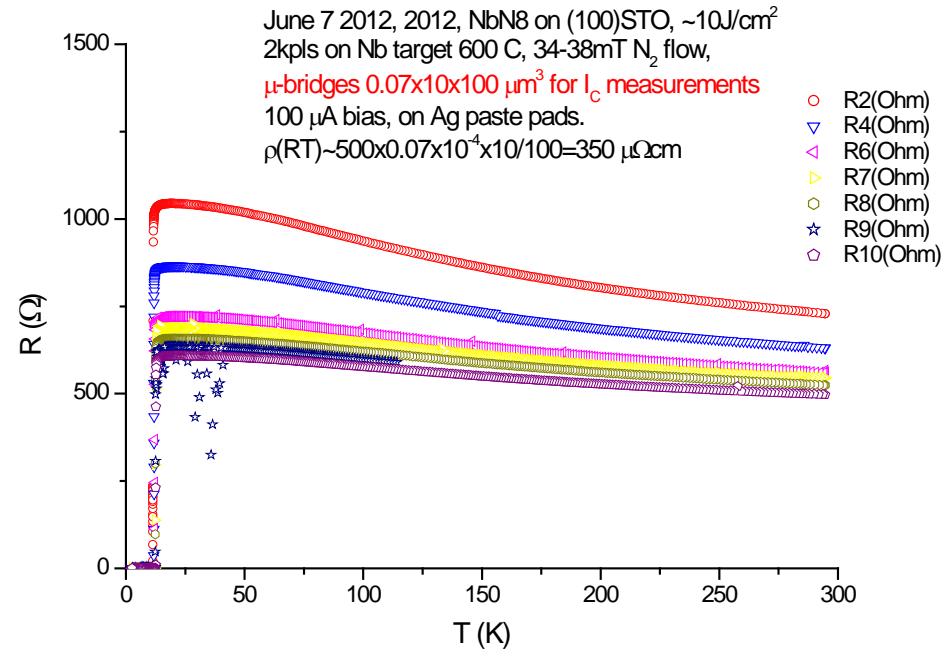
A fully lithographically-patterned junction layout



R vs T of bridges in a 70nm thick NbN & of 100nm Bi₂Se₃/70nm NbN bilayer

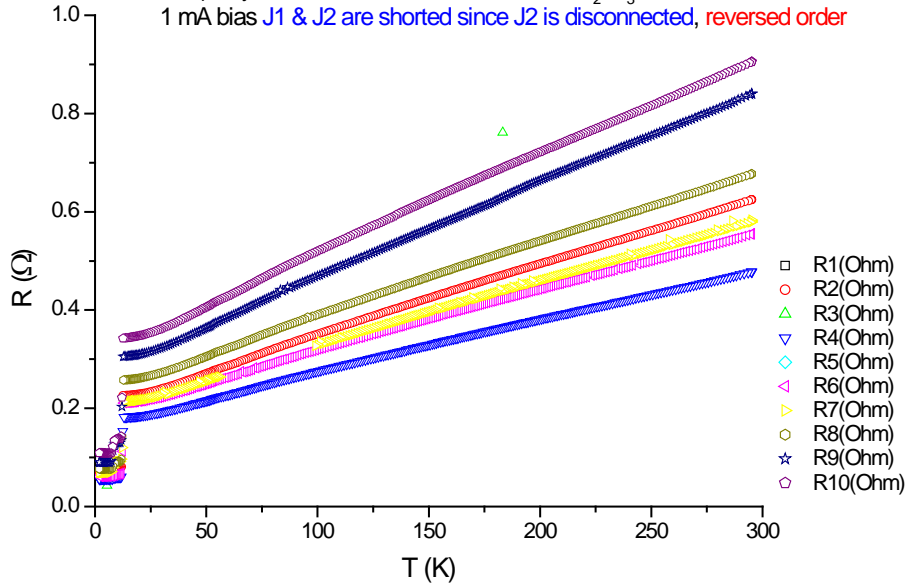
- ρ of bare NbN bridges at RT is similar to that of YBCO, but is not metallic vs T.
- T_c of the bare NbN film is $\sim 12.5K$ (onset)
- $J_c(2K)$ of NbN bridges with $0.07 \times 10 \mu m^2$ cross-section area is $\sim 1.4 MA/cm^2$, which is equivalent to $10 mA/10 \mu m$ width

- An un-patterned bilayer. When patterned, its' resistance R would probably be similar to that of the patterned NbN bridges.

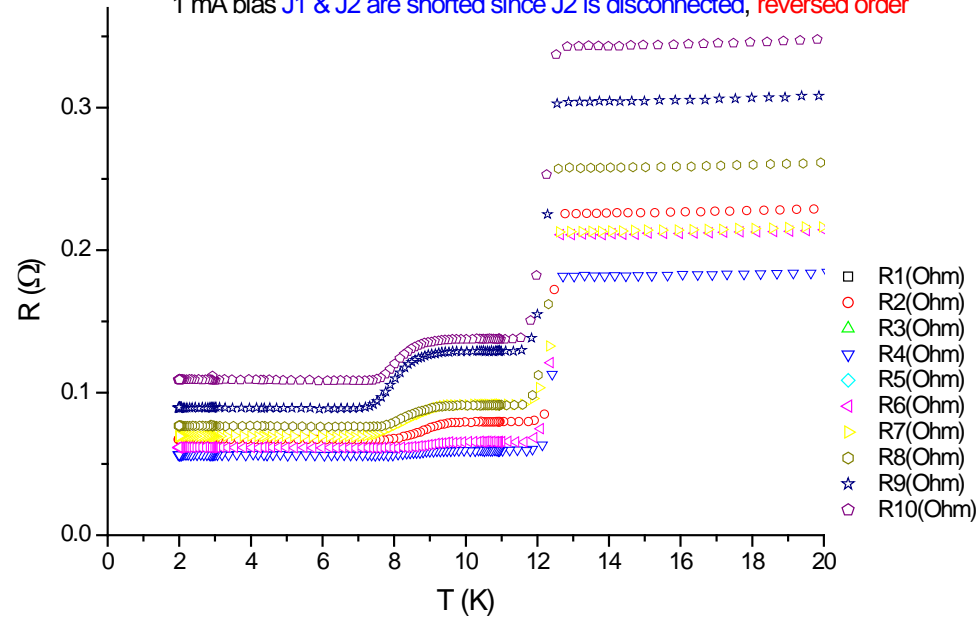


R versus T of 100nm Bi₂Se₃/70nm NbN-100nm Au *in-situ* junctions

June 4, 2012, NbN7-BiSe-Au junctions on (100)STO
 5x5 μm² junctions' area, 100nm Au-100nm Bi₂Se₃/70nm NbN
 1 mA bias J1 & J2 are shorted since J2 is disconnected, reversed order

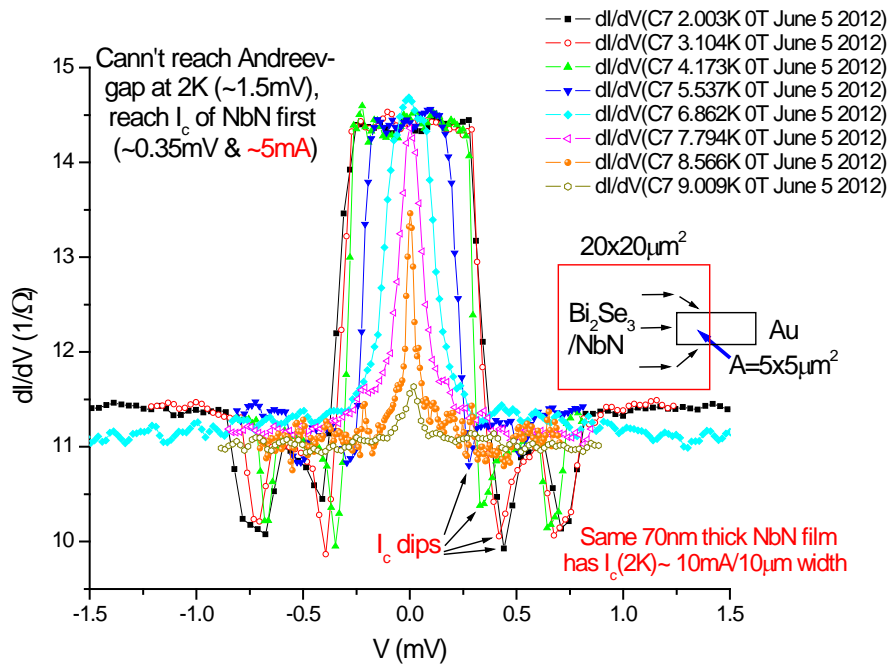


June 5, 2012, NbN7-BiSe-Au junctions on (100)STO
 5x5 μm² junctions' area, 100nm Au-100nm Bi₂Se₃/70nm NbN
 1 mA bias J1 & J2 are shorted since J2 is disconnected, reversed order



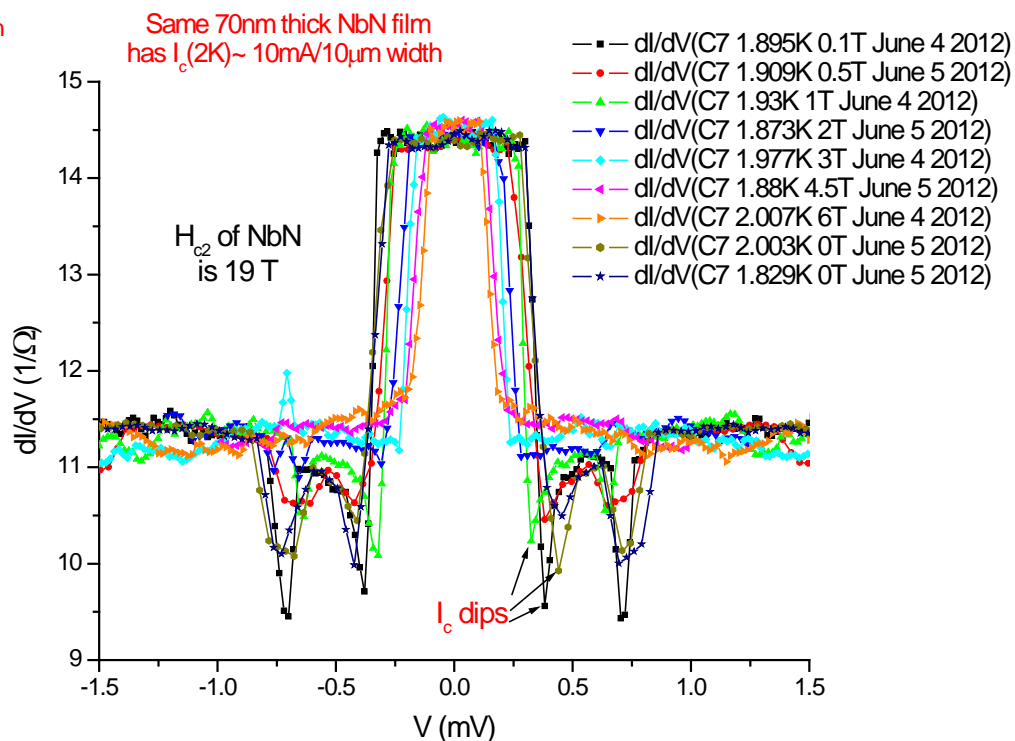
- **Metallic vs T.** Very low ρ even when compared to that of the unpattern bilayer (x10 lower). This seems to be due to additional Se vacancies in the Bi₂Se₃ layer (hole doping), possibly created by the direct Ar ion beam on the Bi₂Se₃ when milling the gold cover electrode. \longrightarrow Should use a shadow mask instead
- T_c of the NbN electrode is ~ 12.5 K (onset) & ~ 8 K of the junction
- Junctions' resistance at low T is very low, about 0.1Ω (cleanest interface due to *in-situ* deposition of the bilayer). In the previous *ex-situ* junctions with the x600 larger area, this resistance was x100 larger!!!

Conductance spectra of J7 at different T & H



- Looks like s-wave Andreev, but gap can't be at 0.35mV.
- Reach I_c first via the 5+10 μm length of the bilayer-Au overlap boundary. See I_c dips.

- 6T is not sufficiently high to suppress the low bias conductance



Conductance spectra of J1, an *in-situ* junction at different T

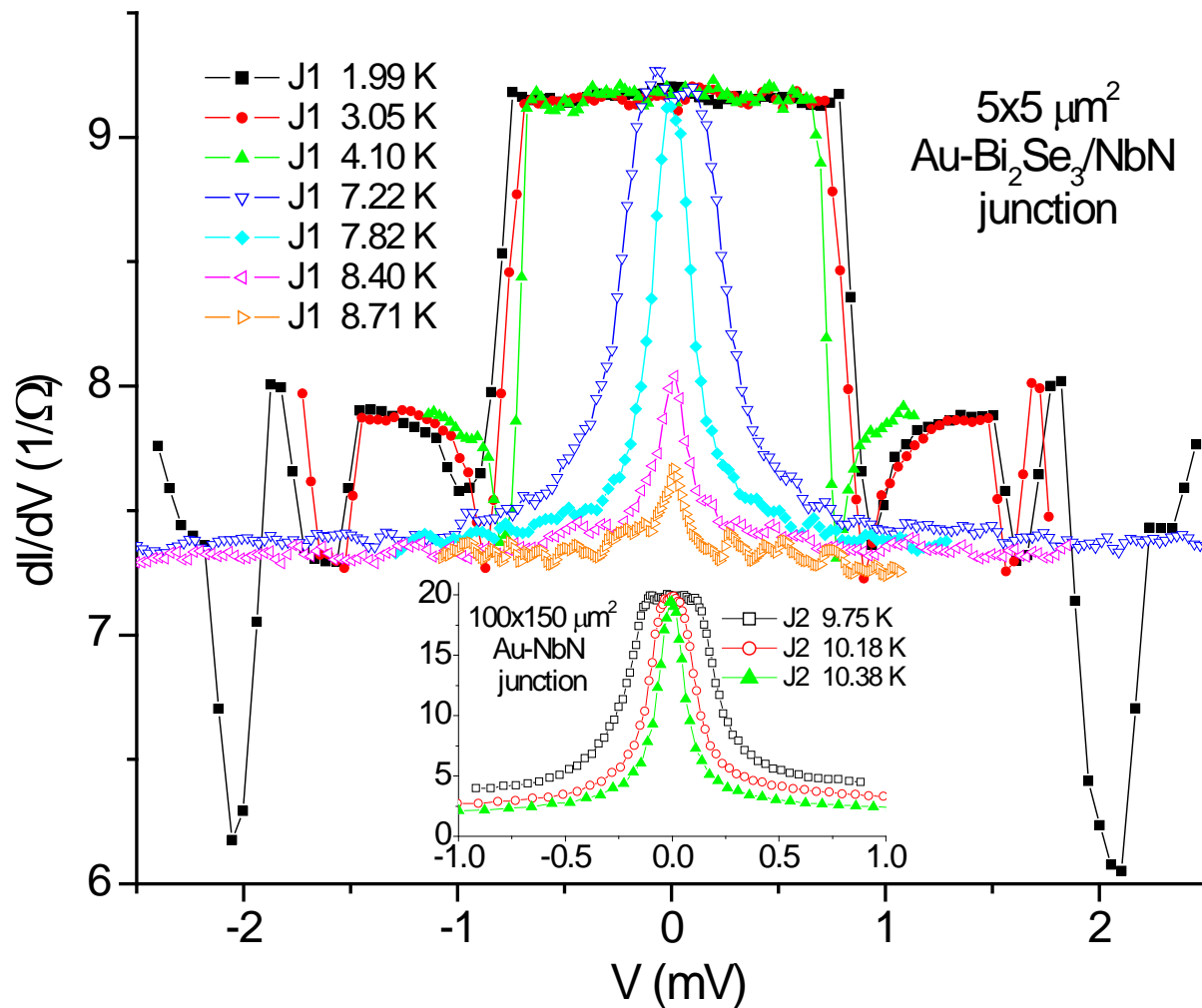
I_c dips at $\sim \pm 1, 1.5$ & 2 mV

No CP, I_c is reached first

No ZBCP, just decreasing critical current with increasing T (and H)

Same seemingly "emergent" ZBCP in reference NbN-Au junction with no Bi_2Se_3 in between

Need a strong tunneling Barrier to observe ZBCPs

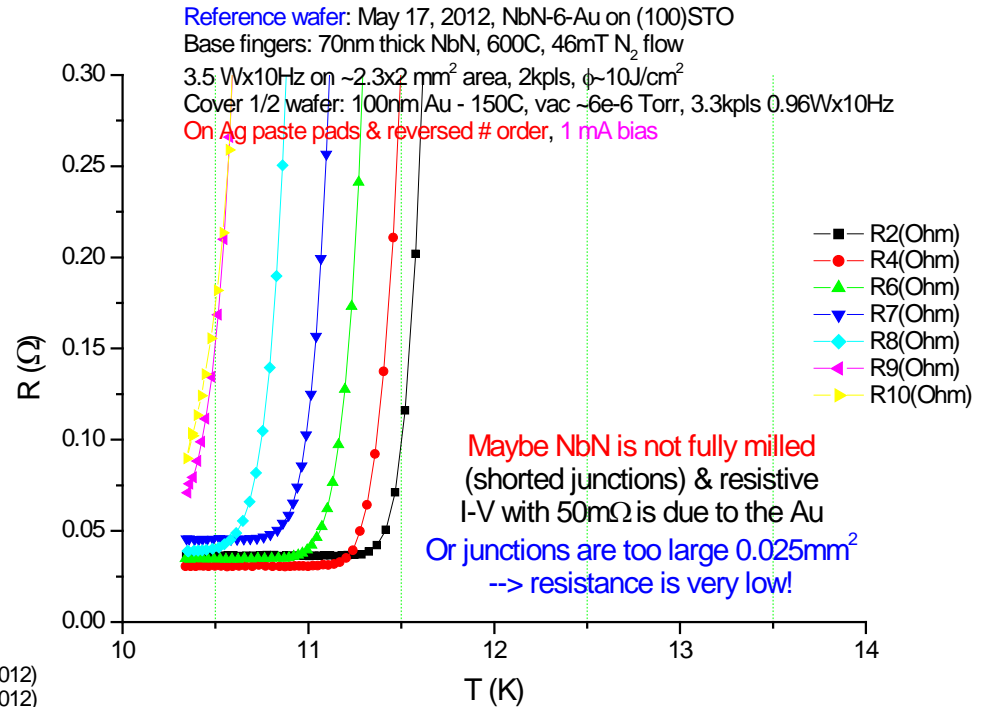
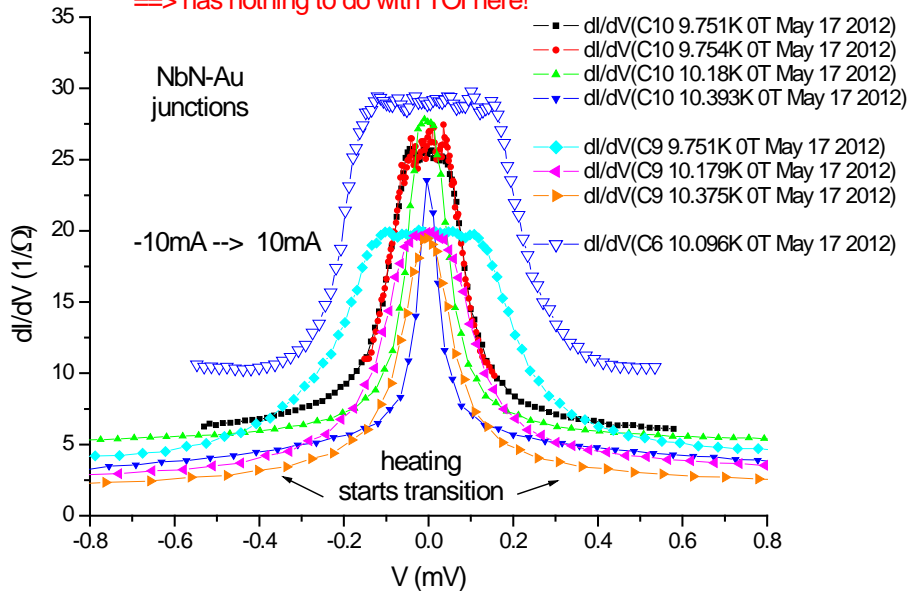


Reference NbN-Au junctions

$T_c^{\text{onset}} \sim 12.5\text{K}$

$R_{\text{junction}} \sim 50\text{ m}\Omega$

Similar to junctions with Bi_2Se_3 barrier
 \Rightarrow has nothing to do with TOI here!



- Same as in Au- Bi_2Se_3 /NbN junctions !!
- The “ZBCP” seems to be an I_c effect due to heating.

Conclusions

- Proximity superconductivity is induced by NbN in Topological Bi_2Se_3 films
- Tunneling and Andreev bound states (ZBCP) were observed
- The broad ZBCP ($\sim 0.9\text{mV}$) can't originate in simple Andreev since coherence peaks exist. It may be due thermally smeared, split ZBCP.
- The narrow ZBCP ($\sim 0.15\text{mV}$) width is kT limited ($\sim 1.8\text{K}$)
- Both types of ZBCP could originate in Majorana fermions of a Topological SC