בית ספר הלן דילר לחישוב קוואנטי, אוקטובר 2020

הרצאה 6

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TECHNION | Helen Diller Quantum Center

Outline:

1. No signaling
2. No cloning
3. Quantum teleportation

No signaling

More about measurements...

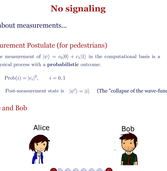
Measurement Postulate (for pedestrians)

The measurement of $|\psi\rangle = \alpha_0|0\rangle + \alpha_1|1\rangle$ in the computational basis is a physical process with a **probabilistic** outcome:

★ $\text{Prob}(i) = |\alpha_i|^2, \quad i = 0, 1$

★ Post-measurement state is $|\psi'\rangle = |i\rangle$ (The "collapse of the wave-function")

Alice and Bob



$|\psi\rangle = \alpha_{00}|00\rangle + \alpha_{01}|01\rangle + \alpha_{10}|10\rangle + \alpha_{11}|11\rangle$

We can also talk about single-side measurement:

Alice: $\text{Prob}(0) = |\alpha_{00}|^2 + |\alpha_{01}|^2 \qquad |\psi'\rangle = \frac{1}{\sqrt{|\alpha_0|^2 + |\alpha_1|^2}}(\alpha_{00}|00\rangle + \alpha_{01}|01\rangle)$

$\text{Prob}(1) = |\alpha_{10}|^2 + |\alpha_{11}|^2 \qquad |\psi'\rangle = \frac{1}{\sqrt{|\alpha_0|^2 + |\alpha_1|^2}}(\alpha_{10}|10\rangle + \alpha_{11}|11\rangle)$

For example:

$$|\psi\rangle = \frac{1}{\sqrt{14}}(|00\rangle - 2|10\rangle + 3|11\rangle)$$

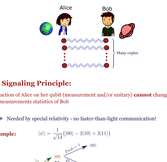
? What is the probability that Alice measures 1 in her qubit, and what is the post measurement state?

$$\alpha_{00}|00\rangle + \alpha_{11}|11\rangle = \frac{1}{\sqrt{14}}(-2|10\rangle + 3|11\rangle)$$

$$\text{Prob}(1) = |\alpha_{10}|^2 + |\alpha_{11}|^2 = \frac{2^2 + 3^2}{14} = \frac{13}{14}$$

$$|\psi'\rangle = \frac{1}{\sqrt{13}}(-2|10\rangle + 3|11\rangle)$$

? Can a measurement or a unitary gate Alice applies "change the state of Bob's qubit"? Can it influence its measurement results?



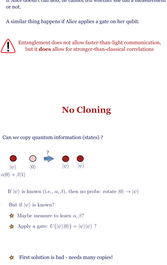
No Signaling Principle:

Any action of Alice on her qubit (measurement and/or unitary) **cannot** change the measurements statistics of Bob

Needed by special relativity - no faster-than-light communication!

Example:

$$|\psi\rangle = \frac{1}{\sqrt{14}}(|00\rangle - 2|10\rangle + 3|11\rangle)$$



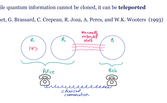
If Alice doesn't call Bob, he cannot tell whether she did a measurement or not.

A similar thing happens if Alice applies a gate on her qubit.

Entanglement does not allow faster-than-light communication, but it **does** allow for stronger-than-classical correlations

No Cloning

? Can we copy quantum information (states) ?



★ If $|\psi\rangle$ is known (i.e., α, β), then no probs: rotate $|0\rangle \to |\psi\rangle$

★ But if $|\psi\rangle$ is known?

- ✯ Maybe measure to learn α, β?
- ✯ Apply a gate: $U|\psi\rangle|0\rangle = |\psi\rangle|\psi\rangle$?

- ✯ **First solution is bad - needs many copies!**
- ✯ **Second solution is impossible:**

Assume by contradiction that such U exists.

Then:

$$U|0\rangle|0\rangle = |0\rangle|0\rangle \ , \qquad U|1\rangle|0\rangle = |1\rangle|1\rangle$$

$$U\Big[|\psi\rangle|0\rangle = U\big[\alpha|0\rangle + \beta|1\rangle\big]|0\big] = \alpha U|0\rangle|0\rangle + \beta U|1\rangle|0\rangle$$
$$= \alpha|0\rangle|0\rangle + \beta|1\rangle|1\rangle \neq \underbrace{(\alpha|0\rangle + \beta|1\rangle)(\alpha|0\rangle + \beta|1\rangle)}_{=|\psi\rangle|\psi\rangle}$$

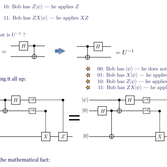
However, cloning (copying) is possible if we restrict ourselves to a subset of states.

? Show that cloning in the standard basis is possible: $U|z\rangle|0\rangle = |z\rangle|z\rangle$.

Teleportation

While quantum information cannot be cloned, it can be **teleported**

C. Bennet, G. Brassard, C. Crepeau, R. Joza, A. Peres, and W.K. Wooters (1993)



The 1-qubit case:

Mathematical Preliminary: Bell's basis

$$U = \boxed{H}$$

$$|\beta_{00}\rangle = U|00\rangle = \boxed{H}\boxed{\ }\ \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$$

$$|\beta_{01}\rangle = U|01\rangle = \boxed{H}\boxed{X}\ \frac{1}{\sqrt{2}}(|01\rangle + |10\rangle)$$

$$|\beta_{10}\rangle = U|10\rangle = \boxed{\ }\boxed{H}\ \frac{1}{\sqrt{2}}(|00\rangle - |11\rangle)$$

$$|\beta_{11}\rangle = U|11\rangle = \boxed{X}\boxed{H}\ \frac{1}{\sqrt{2}}(|01\rangle - |10\rangle)$$

Mathematical fact:

For every single qubit state $|\psi\rangle$

$$|\psi\rangle \otimes |\beta_{00}\rangle = \frac{1}{2}\Big[|\beta_{00}\rangle \otimes |\psi\rangle + |\beta_{01}\rangle \otimes (X|\psi\rangle) + |\beta_{10}\rangle \otimes (Z|\psi\rangle) + |\beta_{11}\rangle \otimes (ZX|\psi\rangle)\Big]$$

1. Let Alice apply $U_{AA'}^\dagg$ on her qubits:

$$(U_{AA'}^\dagg \otimes I_B)|\psi\rangle \otimes |\beta_{00}\rangle = \frac{1}{2}\Big[|00\rangle \otimes |\psi\rangle + |01\rangle \otimes (X|\psi\rangle) + |10\rangle \otimes (Z|\psi\rangle) + |11\rangle \otimes (ZX|\psi\rangle)\Big]$$

2. Now she measures her qubits and calls Bob to tell him the results:

- ✯ 00: Bob has $|\psi\rangle$ — he does nothing
- ✯ 01: Bob has $X|\psi\rangle$ — he applies X
- ✯ 10: Bob has $Z|\psi\rangle$ — he applies Z
- ✯ 11: Bob has $ZX|\psi\rangle$ — he applies XZ

What is U^{-1} ?

$$U = \boxed{H}\boxed{\ } \quad \Rightarrow \quad \boxed{\ }\boxed{H}\boxed{X}\boxed{Z} = U^{-1}$$

- ✯ 00: Bob has $|\psi\rangle$ — he does nothing
- ✯ 01: Bob has $X|\psi\rangle$ — he applies X
- ✯ 10: Bob has $Z|\psi\rangle$ — he applies Z
- ✯ 11: Bob has $ZX|\psi\rangle$ — he applies XZ

Wrapping it all up:

Proving the mathematical fact:

(*) $|\psi\rangle \otimes |\beta_{00}\rangle = \frac{1}{2}\Big[|\beta_{00}\rangle \otimes |\psi\rangle + |\beta_{01}\rangle \otimes (X|\psi\rangle) + |\beta_{10}\rangle \otimes (Z|\psi\rangle) + |\beta_{11}\rangle \otimes (ZX|\psi\rangle)\Big]$

We will use the identities:

$$|\beta_{00}\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$$
$$|\beta_{01}\rangle = \frac{1}{\sqrt{2}}(|01\rangle + |10\rangle) = X \otimes I|\beta_{00}\rangle = I \otimes X|\beta_{00}\rangle$$
$$|\beta_{10}\rangle = \frac{1}{\sqrt{2}}(|00\rangle - |11\rangle) = Z \otimes I|\beta_{00}\rangle = I \otimes Z|\beta_{00}\rangle$$
$$|\beta_{11}\rangle = \frac{1}{\sqrt{2}}(|01\rangle - |10\rangle) = ZX \otimes I|\beta_{00}\rangle = I \otimes XZ|\beta_{00}\rangle$$

$$RHS = (I \otimes I_B \otimes I) \otimes X \otimes I \otimes Z \otimes I + I_B \otimes X \otimes Z \otimes (XZ)|\beta_{00}\rangle_{AA'}|\psi\rangle_B$$

To prove (*) we will multiply the RHS by $_{AA'}|\beta_{00}\rangle$ and see that we get $|\psi\rangle_B$.

The identities also imply:

$$\langle\beta_{00}| = \langle\beta_{00}|X \otimes X = \langle\beta_{00}|Z \otimes Z = \langle\beta_{00}|XZ \otimes ZX$$

$$\Rightarrow \langle\beta_{00}| \cdot RHS = {}_{AA'}\langle\beta_{00}| \cdot |\beta_{00}\rangle_{AA'}|\psi\rangle_B$$

Finally, we claim $_{AA'}\langle\beta_{00}| \cdot |\beta_{00}\rangle_{AA'}|\psi\rangle_B = |\psi\rangle_B$.

This is a linear identity, and therefore it is sufficient to verify for $|\psi\rangle = |0\rangle$ and $|\psi\rangle = |1\rangle$ — which can easily be done explicitly.